

# Analytic heliospheric magnetic field modeling

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## Talk outline:

### ① The Base Model

Introduction and Idea  
Solution Properties

### ② Improvements

Extension I: Non-circular Cross Sections  
Extension II: Compressible Flow  
(Some ideas for) Extension III: The Inside



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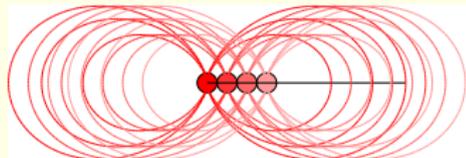
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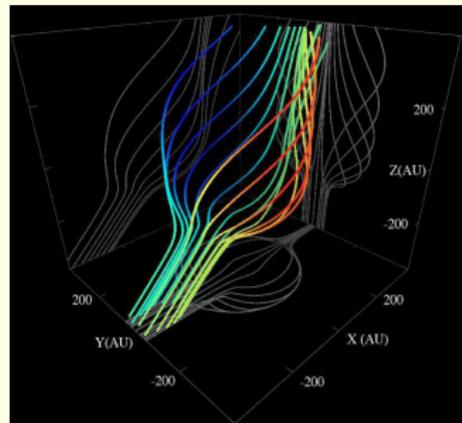
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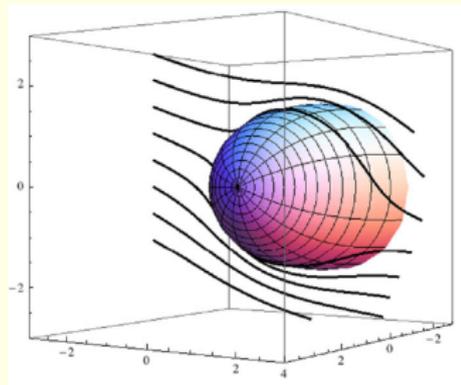
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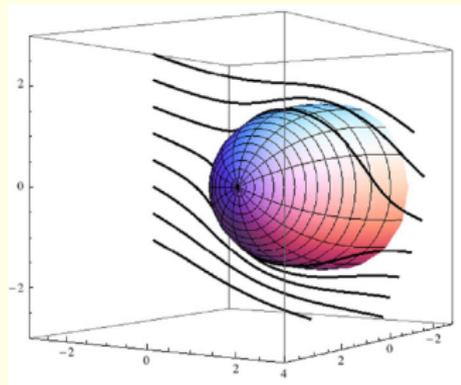
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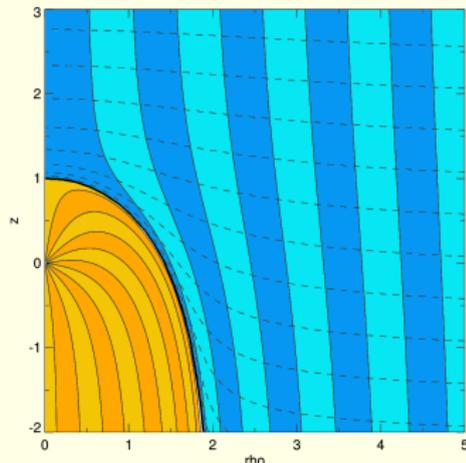
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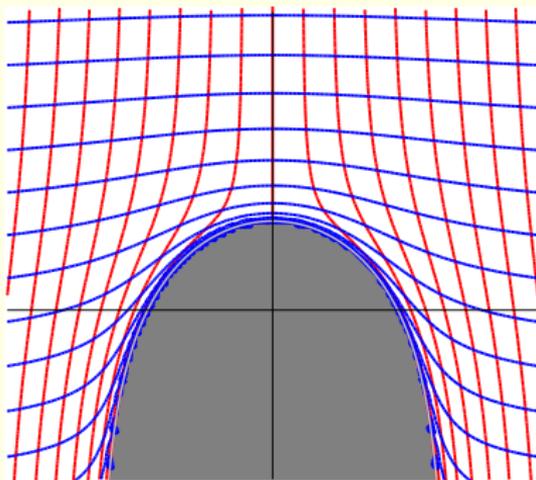
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## Key idea

- **Stream lines** and **isochrones** (lines of constant travel time  $T$ ) form a non-orthog. coordinate system  $\mathcal{K}$  exterior to HP.
- Advected **B** field components are constant w. r. t.  $\mathcal{K}$  due to **line conservation** [Cauchy 1816].



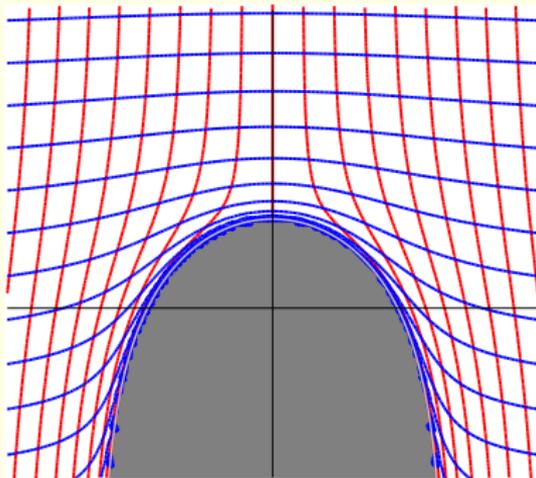
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Analytic evaluation of  
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$$T(r, \vartheta) = \int_{\infty}^r \frac{dr'}{u_r(r')} = \int_0^{\vartheta} \frac{r(\vartheta') d\vartheta'}{u_{\vartheta}(\vartheta')}$$

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## The exact solution [Röken+, ApJ 2015]

$$B_\rho(\rho, \varphi, z) = -\frac{q\rho}{r^3} B_{z0} + \left[ \frac{q^{3/2}\rho}{r^3 a} \mathcal{T} + \frac{a}{\rho} \left( 1 + \frac{qz}{r^3} \right) \right] B_{\rho0}$$

$$B_\varphi(\rho, \varphi, z) = \frac{\rho}{a} B_{\varphi0}$$

$$B_z(\rho, \varphi, z) = \left( 1 - \frac{qz}{r^3} \right) B_{z0} + \left[ \left( \frac{qz}{r^3} - 1 \right) \frac{\sqrt{q}}{a} \mathcal{T} + \frac{qz^2 a}{r^3 \rho^2} \right] B_{\rho0}$$

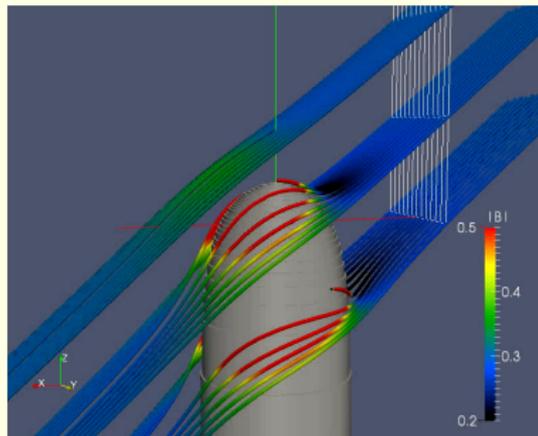
with  $\mathbf{B}_0 = \mathbf{B}|_\infty$ , incomplete elliptic integrals  $\{F, E\}$ , and

$$\mathcal{T} := (2 - \kappa^{-2}) E(\lambda, \kappa) - (1 - \kappa^{-2}) F(\lambda, \kappa) \quad \left| \quad \begin{array}{l} \lambda := \sqrt{1 - (a/\rho)^2} \\ \kappa := \sqrt{1 + a^2/(4q)} \end{array} \right.$$

$$a := \sqrt{\rho^2 + 2q(z/r - 1)}$$

Visualization of field line structure shows expected behavior:

- undisturbed ISM field at large distances from HP
- field lines do not penetrate HP, but drape around it, eventually becoming tangential to HP
- draping increases field strength (reaching  $\infty$  on HP!)



# Model extension I: Non-circular cross sections

- $\mathbf{u}$  axially symmetric  $\Rightarrow$  heliotail has **circular cross section** (despite 'squeezing' due to  $\mathbf{B}_{\text{ism}}$ ). This is somewhat unrealistic...
- ...but can be fixed (i.e., adjusted to tailor-made aspect ratio  $\eta(z) := a/b$  and area ( $\propto a b$ ) using a **distortion flow  $\mathbf{w}$  ( $\neq \mathbf{u}$ )** and still remain exact if  $\nabla(\nabla \cdot \mathbf{w}) = \mathbf{0}$ . [Kleimann+, ApJ 2016]

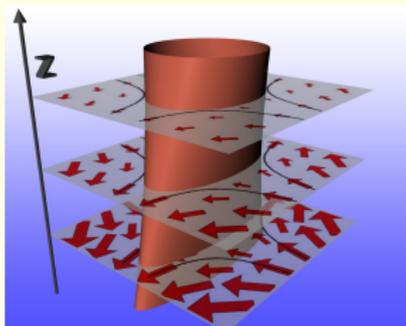
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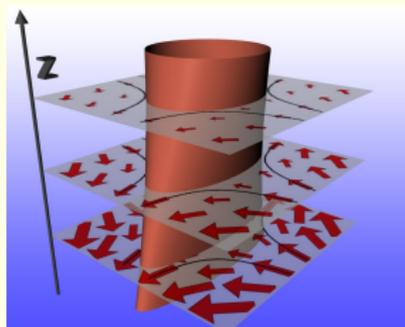


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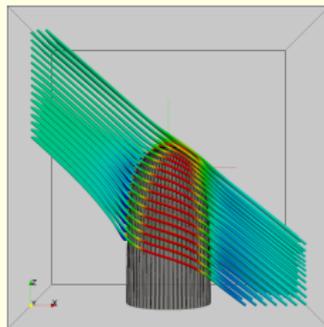
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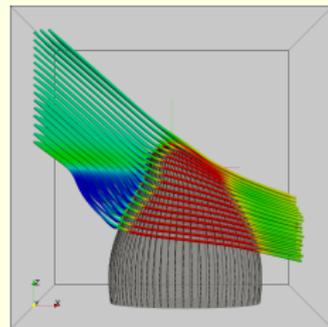
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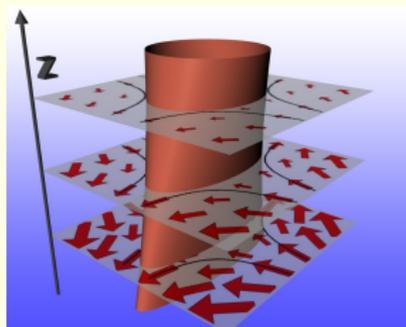
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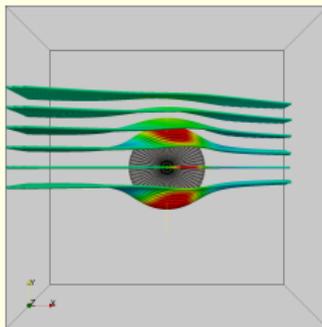
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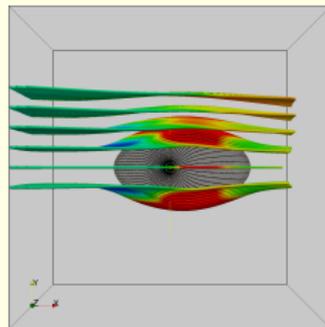
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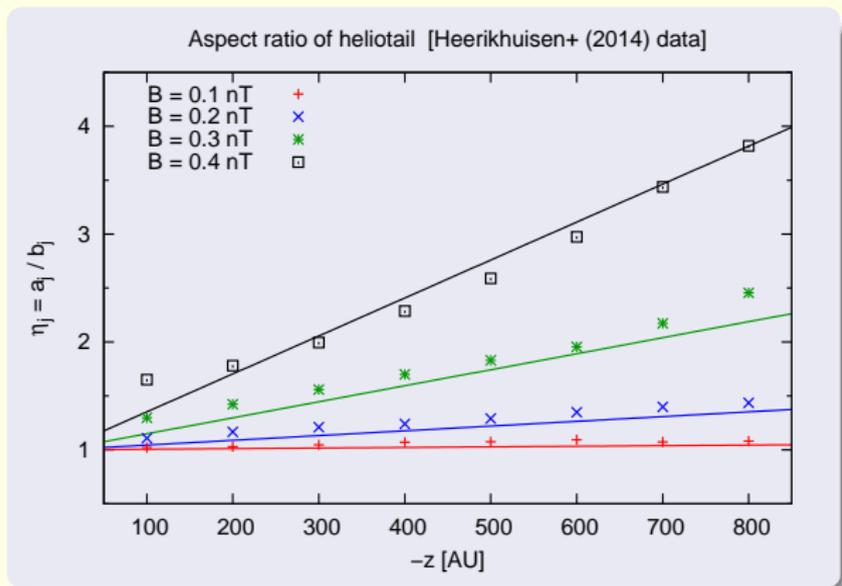
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Realistic parameters for  $\eta(z, B_{\text{ism}})$  from simulations

$$\eta_{\text{fit}}(z, B_{\text{ism}}) = 1 + 3.27 \left( \frac{-z}{100 \text{ AU}} \right) \left( \frac{B_{\text{ism}}}{\text{nT}} \right)^{2.5}$$

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- Issues:**
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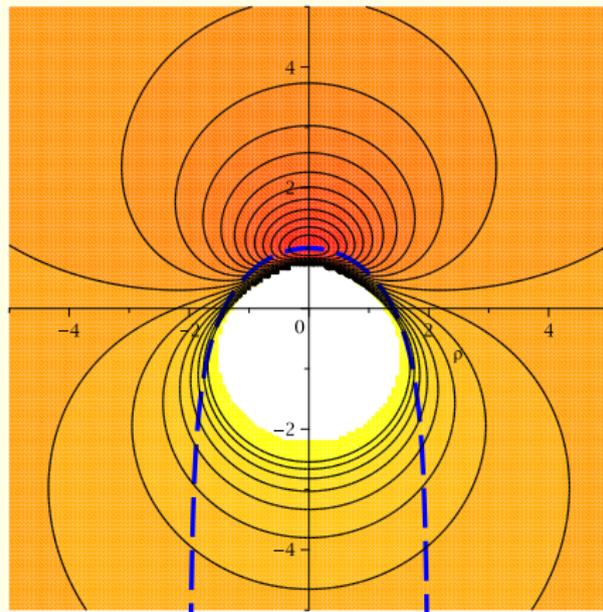
**Requires:** System closure through **Bernoulli's law**  
 $(\mathbf{u} \cdot \nabla)\mathbf{u} = -(\nabla P)/n$  along flow lines,  
plus equation of state  $P \propto n^\gamma$ ,  $\gamma \in \{1, 5/3\}$

## Key findings (density / flow field)

- 1 New parameter: **upstream Mach number**  $m \in [0, 1[$ .  
( $m = 0 \Leftrightarrow n = \text{const.}$ )
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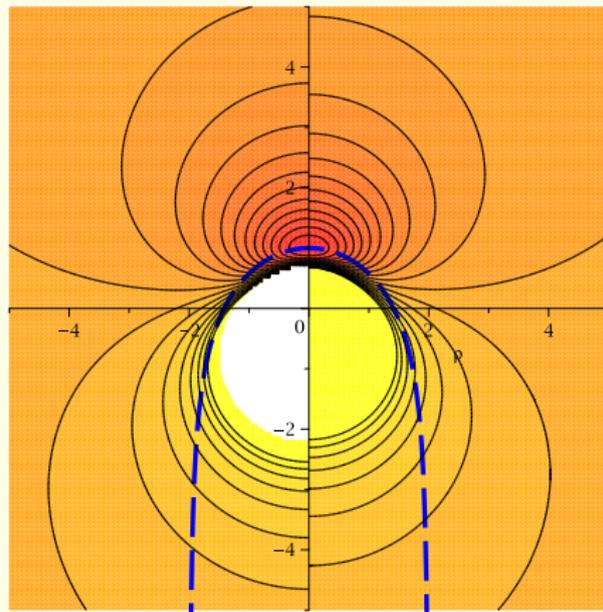
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$n$  contours ( $\gamma = 1, m = 0.6$ )

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- 3 Density looks plausible; reasonable polynomial **approximation** for  $n$



$n_{\text{exact}}$  vs.  $n_{\text{approx}}$  ( $\gamma = 1, m = 0.6$ )

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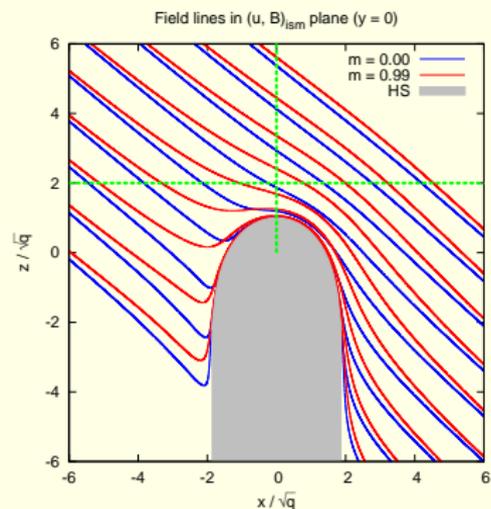
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 $\mathbf{B}$  thus derived from  $\mathbf{u}_{\text{approx}}$  is again an exact MHD solution!

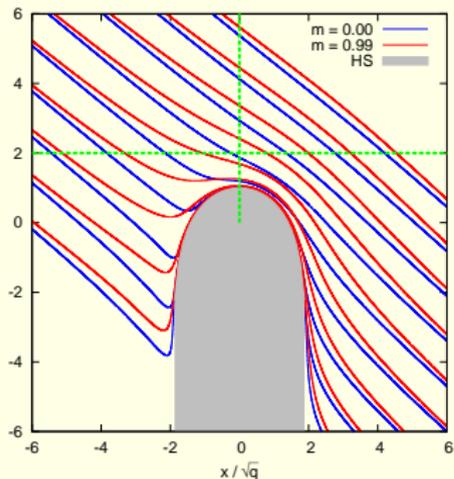
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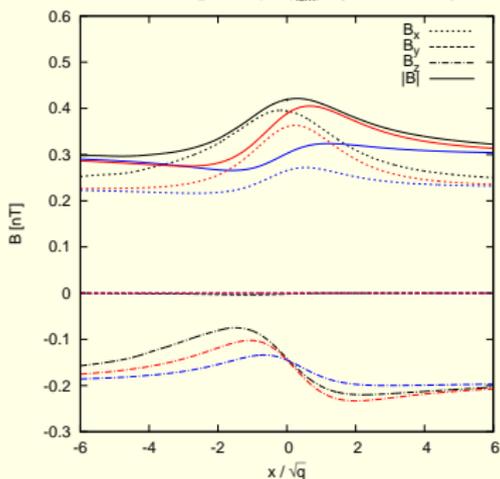
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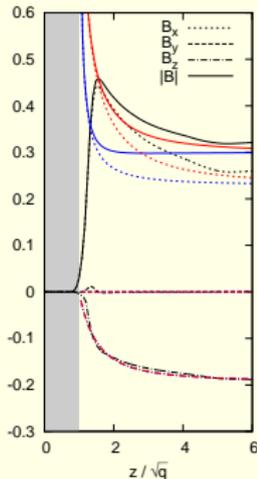
Field lines (**compress.**  
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Comparison  $\mathbf{B}_{\text{old}} \leftrightarrow \mathbf{B}_{\text{new}} \leftrightarrow$  self-consistent MHD sim.Field lines in  $(u, B)_{\text{ISM}}$  plane ( $y = 0$ )

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Crosswind [parallel  $(u, B)_{\text{ISM}}$ ] at  $y = 0, z = 2.0 \sqrt{q}$ 

Cut along  $x$  and  $z$  (numerics in black)  
Note **blue**  $\rightarrow$  **red**  $\rightarrow$  **black** ordering!

Upwind at  $x = 0, y = 0$ 

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**Goal:** Extend method to the actual heliosphere (= HP interior). In principle, **same PDEs** with different boundary conditions.

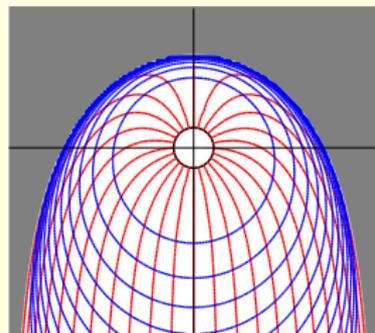
## Difficulties

- 1 stream lines kink at termination shock (TS)
- 2 Inside TS: diverging  $\|\nabla\Phi\| \propto 1/r^2$
- 3 realistic time-dependent solar cycle BCs that leave  $\nabla \cdot \mathbf{B} = 0$  intact

*Minimal working example:* magnetic rings

Parker spiral has  $B_\varphi \gg B_r$  almost everywhere.

$$\mathbf{B}_0(\mathbf{r}_0, t) = B_0(r, \vartheta', t) \mathbf{e}_{\varphi'}$$
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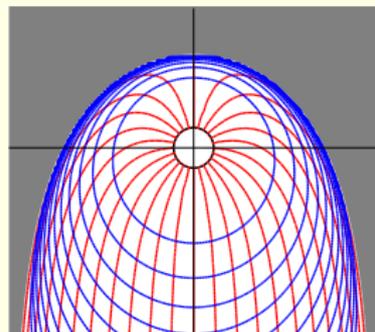


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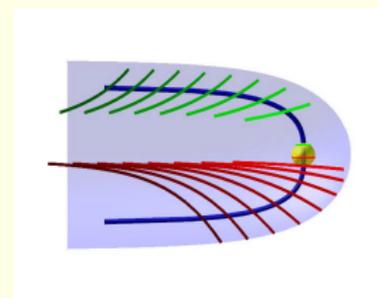
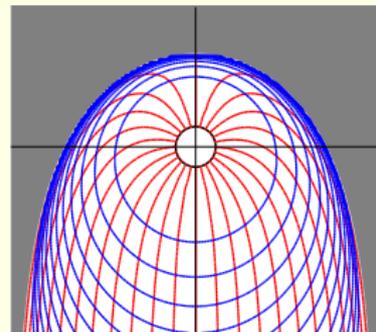
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- Exact(!) compressible magnetic field solution to now even closer to self-consistent MHD numerics.
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**BACKUP SLIDES**

# A note on the approximation of non-existent solutions

## Extension into “white” region

- lacks a measure of departure from “real” flow solution, but
- + follows stream lines,
- + conserves mass,
- + and looks as expected.

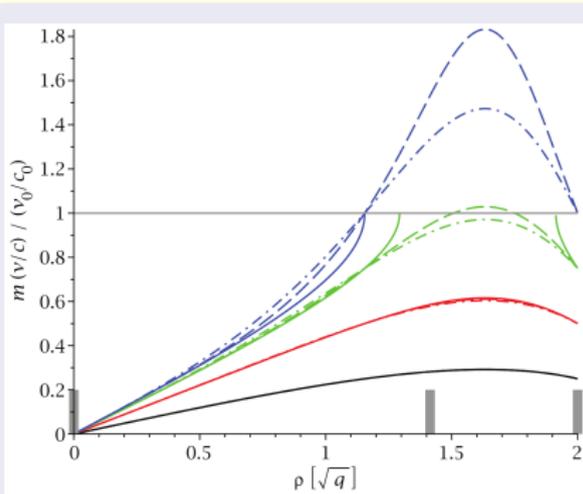
Local(!) Mach number for  $m \in \{0.25, 0.5, 0.75, 1.0\}$  along HP  
(Marks for up/cross/downwind)

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NB:  $\|\mathbf{u}_{\text{appr.}}\|/c_s \leq 1.83$   
(supersonic, despite  $m < 1$ )



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