

Radar detection of high-energy neutrino induced particle cascades in ice

Krijn de Vries ¹

Kael Hanson ^{2,3}

Thomas Meures ²

Aongus O'Murchadha ²

IIHE

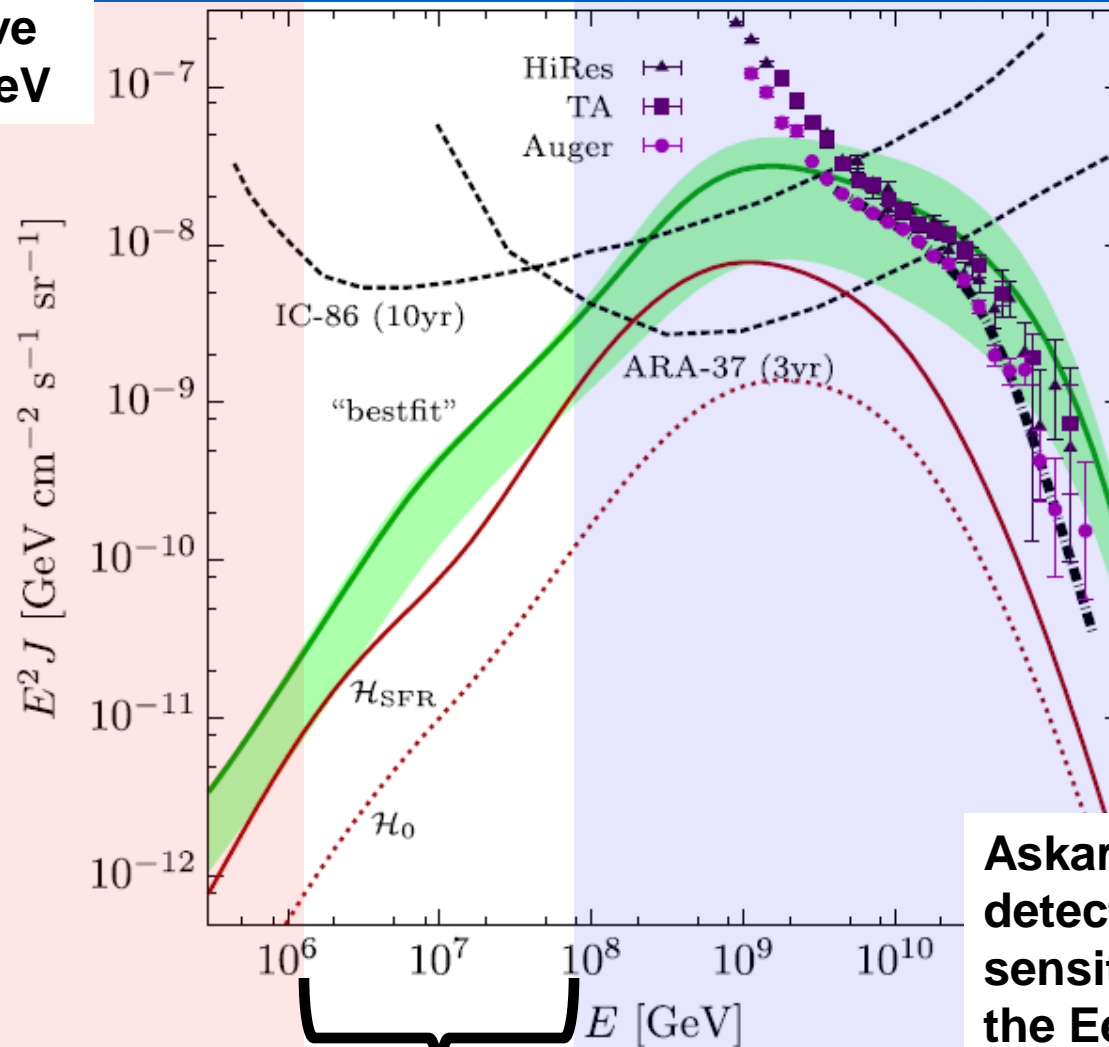
Vrije Universiteit Brussel ¹

UW -Madison ²

Université Libre de Bruxelles ³

Motivation

IceCube sensitive
below several PeV

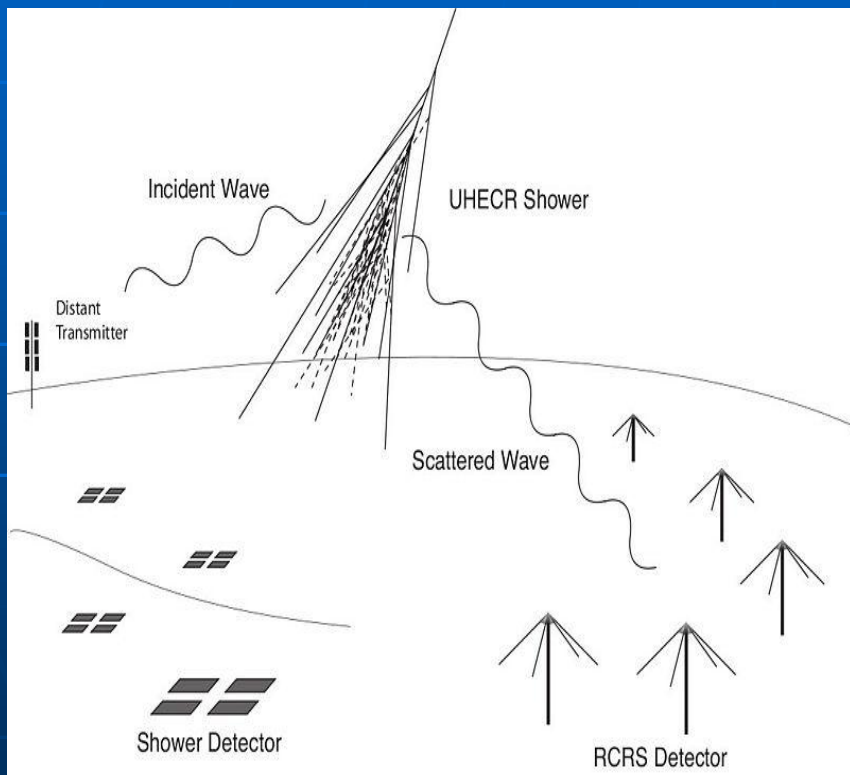


Sensitivity Gap in
PeV – EeV region

Askaryan Radio
detectors become
sensitive close to
the EeV region

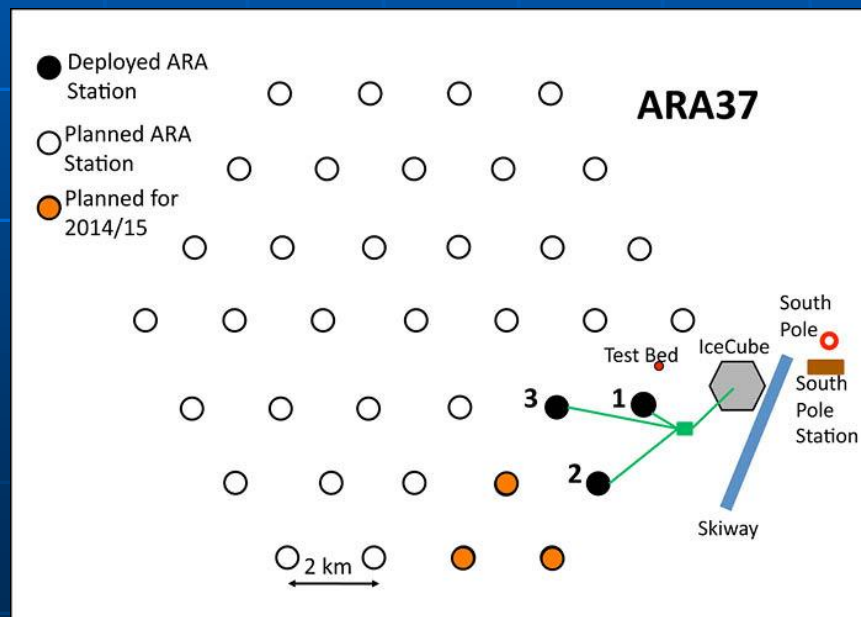
New detection method

If a RADAR signal can be bounced off of a neutrino induced cascade in ice, we have **control over the signal strength!**



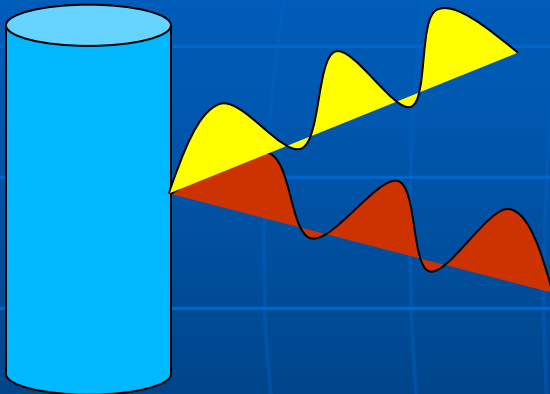
M. Abou Bakr Othman et al,
Proceedings 32nd ICRC, Beijing 2011

Infrastructure already available!



RADAR scattering

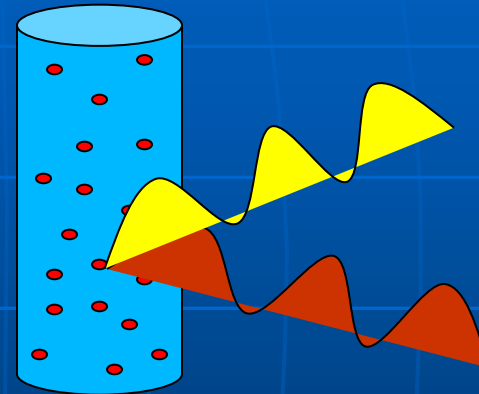
- Over-dense scattering:



Radar frequency < Plasma Frequency

Reflection from the surface of the plasma tube

- Under-dense scattering:



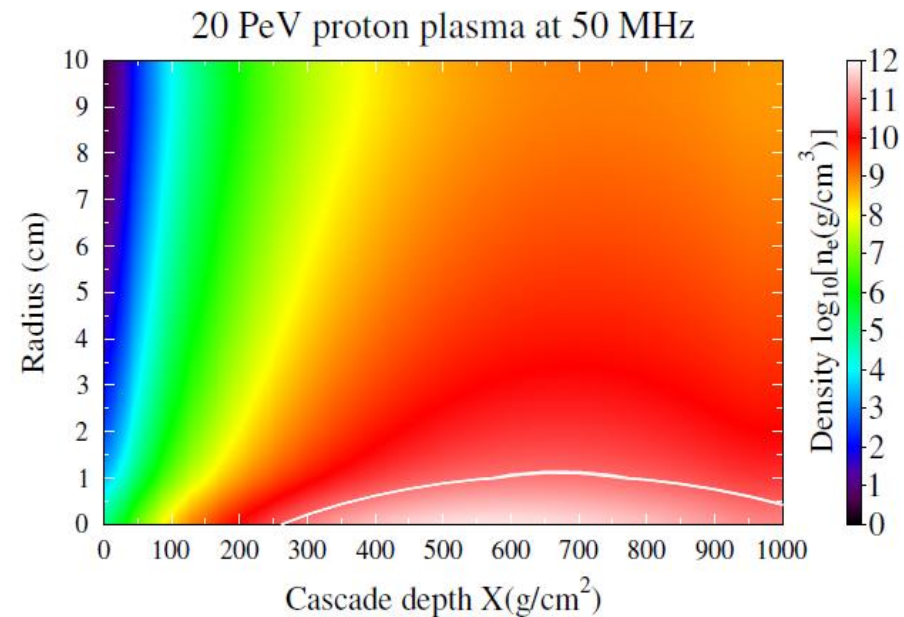
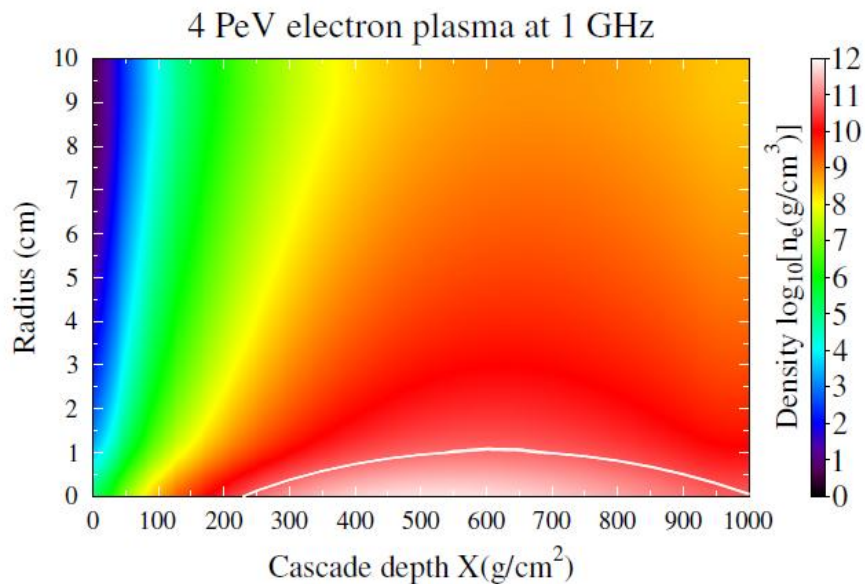
Radar frequency > Plasma Frequency

Scattering off of the individual charges in the plasma

Over-dense scattering

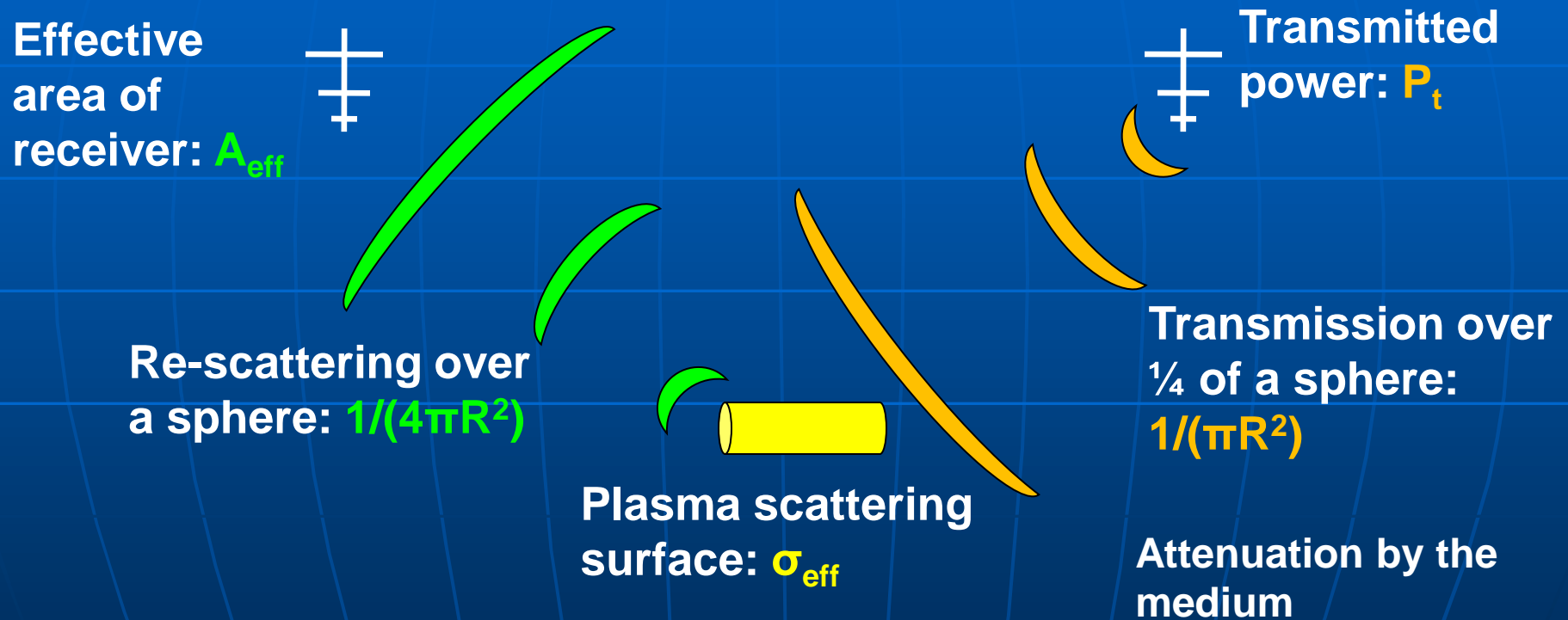
$$v_{Plasma} > v_{Radar} > \begin{cases} 1/\tau_{Plasma} & c_{med}\tau_e < l_c \\ c_{med}/l_c & c_{med}\tau_e > l_c \end{cases}$$

$$v_{Plasma} \propto \sqrt{n_{Plasma}} \propto \sqrt{E_{primary}}$$



RADAR return power estimation

Bi-static RADAR configuration



$$P_r = P_t \eta \frac{\sigma_{eff}}{\pi R^2} \frac{A_{eff}}{4\pi R^2} e^{-4R/L_\alpha}$$

RADAR return power estimation (single antenna)

$$P_r = P_t \eta \frac{\sigma_{eff}(\lambda)}{\pi R^2} \frac{A_{eff}(\lambda)}{4\pi R^2} e^{-4R/L_\alpha}$$

$$\lambda = 0.18 \text{ m}$$

$$\sigma_{eff}^{max} = 0.11 \text{ m}^2$$

$$\sigma_{eff}(\theta = 60^\circ, \phi = 60^\circ) = 1.6 \cdot 10^{-4} \text{ m}^2$$

$$L_\alpha = 1 \text{ km}$$

$$P_{noise} = k_b T_{sys} \Delta \nu$$

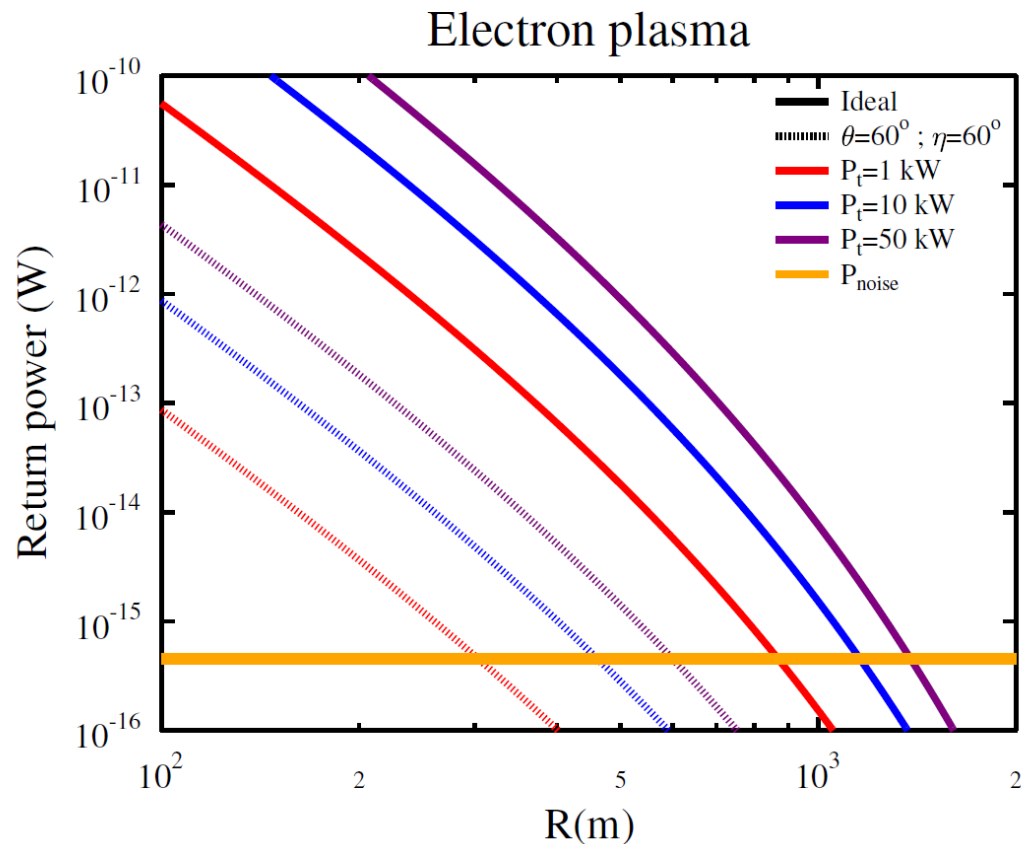
$$T_{sys} = 325 \text{ K}$$

$$\Delta \nu = 100 \text{ kHz}$$

N antennas :

$$P_{Noise}(N) = N \cdot P(N = 1)$$

$$P_{Signal}(N) = N^2 \cdot P(N = 1)$$



RADAR return power estimation (single antenna)

$$P_r = P_t \eta \frac{\sigma_{eff}(\lambda)}{\pi R^2} \frac{A_{eff}(\lambda)}{4\pi R^2} e^{-4R/L_\alpha}$$

$$\lambda = 3.6 \text{ m}$$

$$\sigma_{eff}^{\max} = 5.5 \text{ m}^2$$

$$\sigma_{eff}(\theta = 60^\circ, \phi = 60^\circ) = 1.2 \cdot 10^{-2} \text{ m}^2$$

$$L_\alpha = 1.4 \text{ km}$$

$$P_{\text{noise}} = k_b T_{\text{sys}} \Delta \nu$$

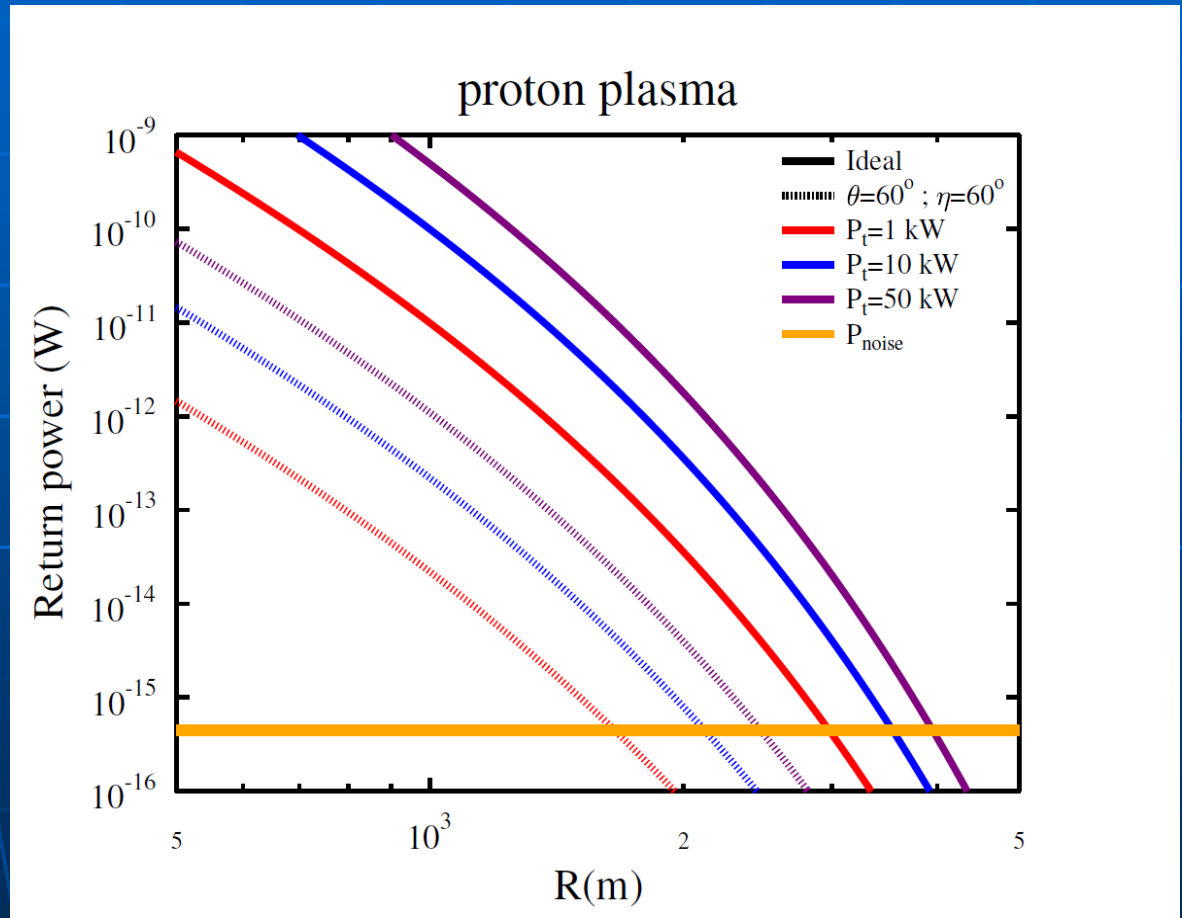
$$T_{\text{sys}} = 325 \text{ K}$$

$$\Delta \nu = 100 \text{ kHz}$$

N antennas :

$$P_{\text{Noise}}(N) = N \cdot P(N = 1)$$

$$P_{\text{Signal}}(N) = N^2 \cdot P(N = 1)$$



RADAR return power estimation (single antenna)

$$P_r = P_t \eta \frac{\sigma_{eff}(\lambda)}{\pi R^2} \frac{A_{eff}(\lambda)}{4\pi R^2} e^{-4R/L_\alpha}$$

$$\lambda = 2.6 \text{ m}$$

$$\sigma_{eff}^{max} = 5.5 \text{ m}^2$$

$$\sigma_{eff}(\theta = 60^\circ, \phi = 60^\circ) = 1.2 \cdot 10^{-2} \text{ m}^2$$

$$L_\alpha = 1.4 \text{ km}$$

$$P_{noise} = k_b T_{sys} \Delta \nu$$

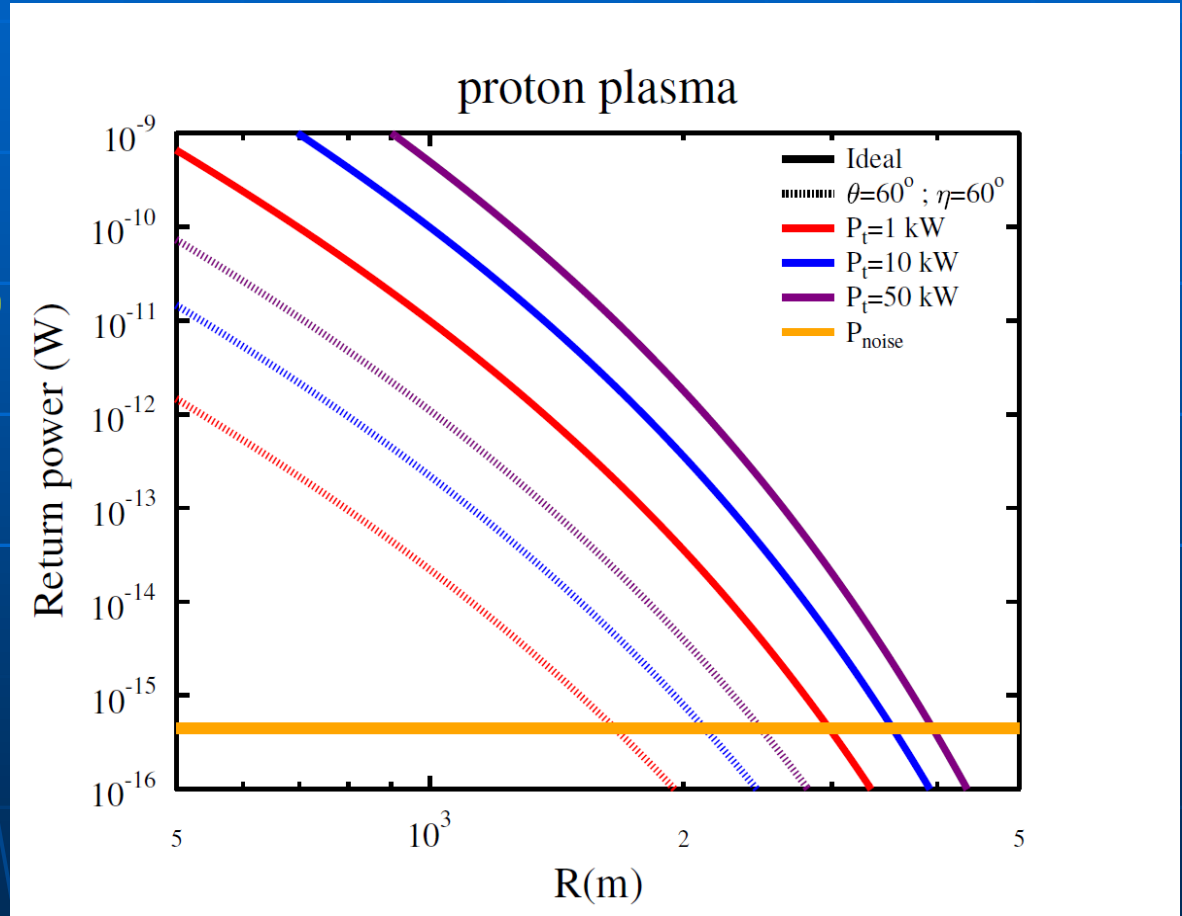
$$T_{sys} = 325 \text{ K}$$

$$\Delta \nu = 100 \text{ kHz}$$

N antennas :

$$P_{Noise}(N) = N \cdot P(N = 1)$$

$$P_{Signal}(N) = N^2 \cdot P(N = 1)$$



Open questions: The Plasma

- How large is the over-dense plasma?
- What is the influence of skin-effects?
- What is the lifetime of the plasma?
- Is the plasma collision frequency low enough?

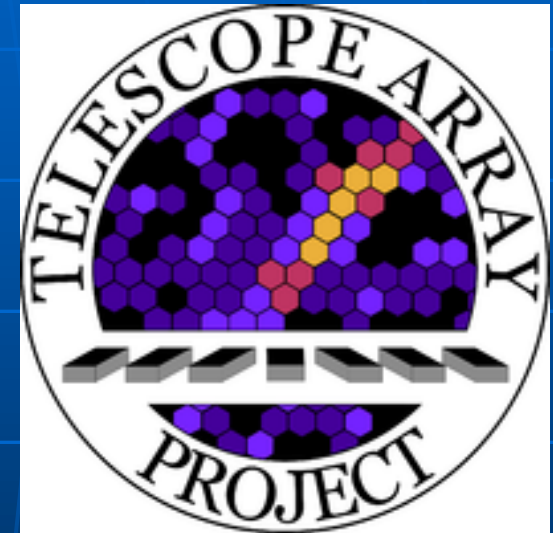


**Experimental verification
needed!**

Radar scattering experiment at TA-ELS

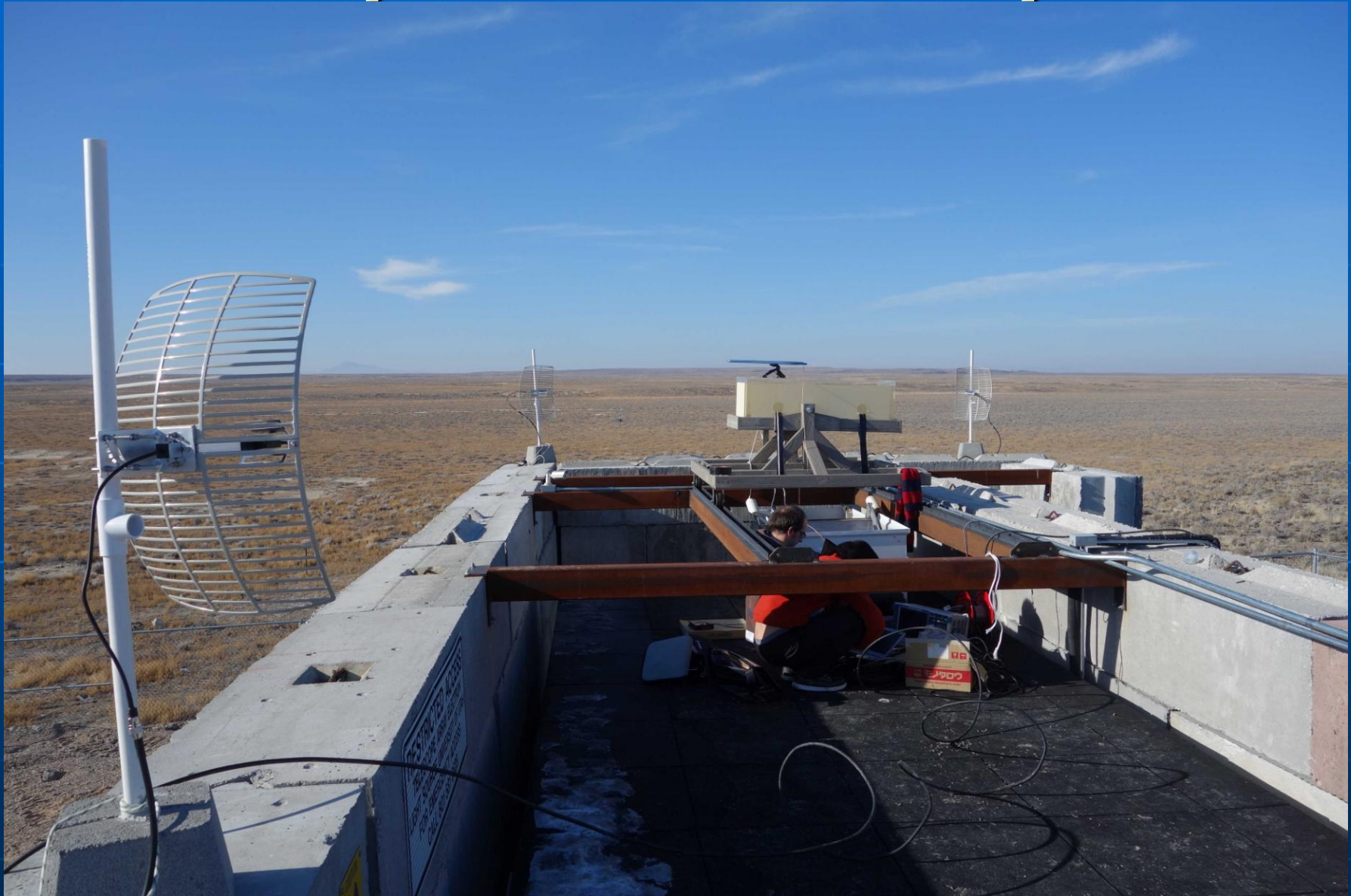


Aya
Matt
Kael

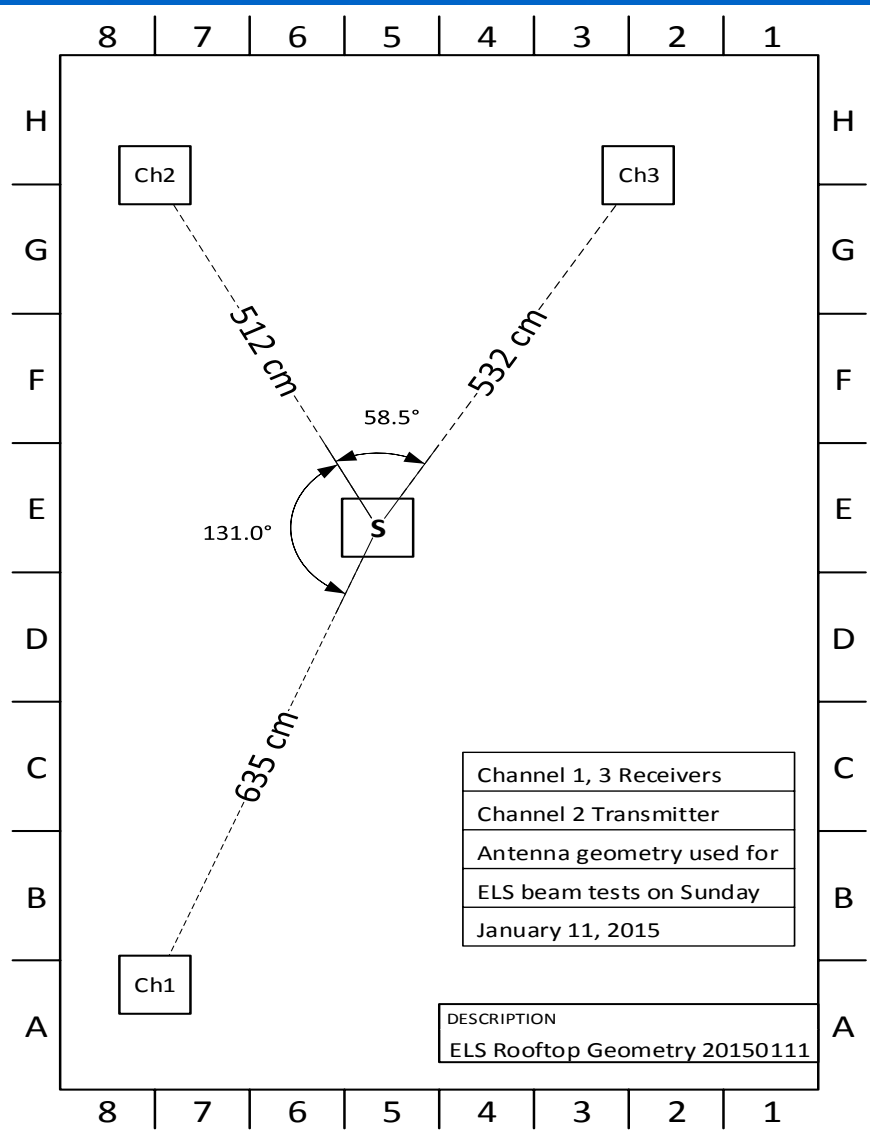


Many thanks to the **Chiba group** and
the **Telescope Array Collaboration** !

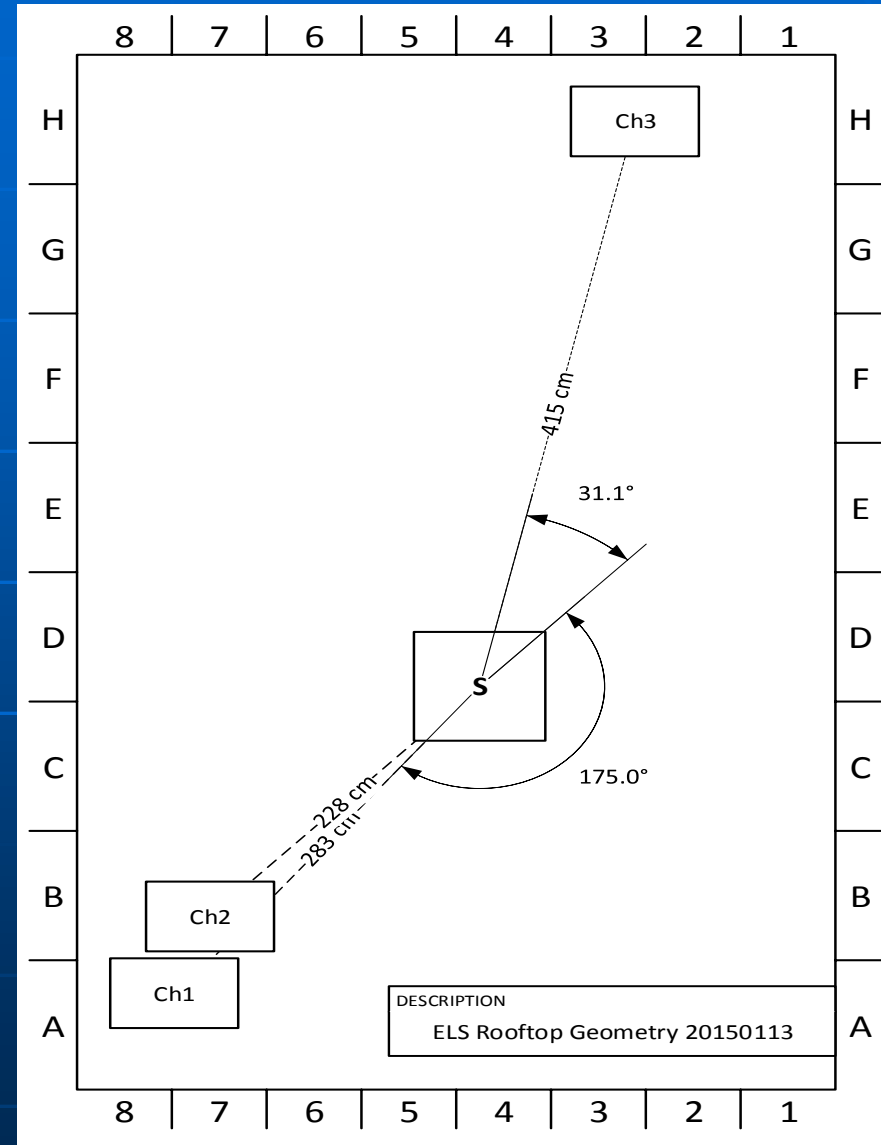
Experimental setup



Experimental setup



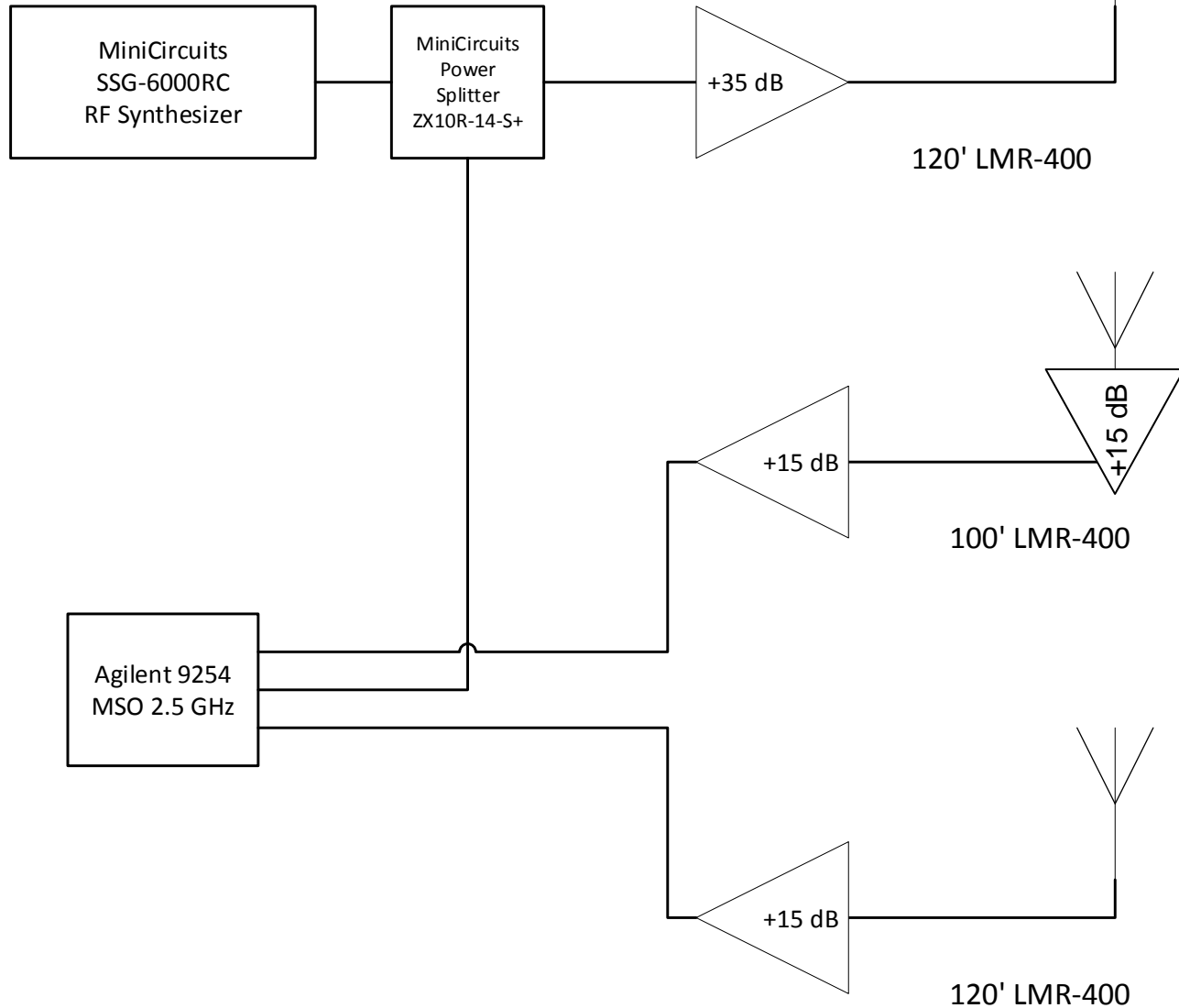
Early Configuration



Later Configuration

Signal chain

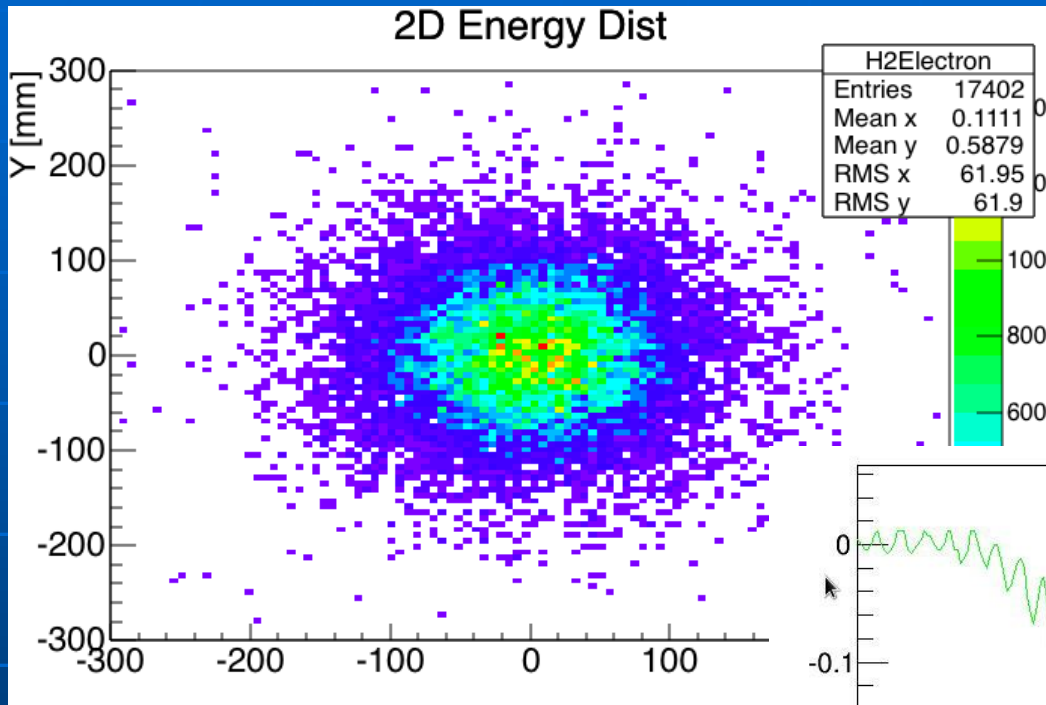
Tx Chain



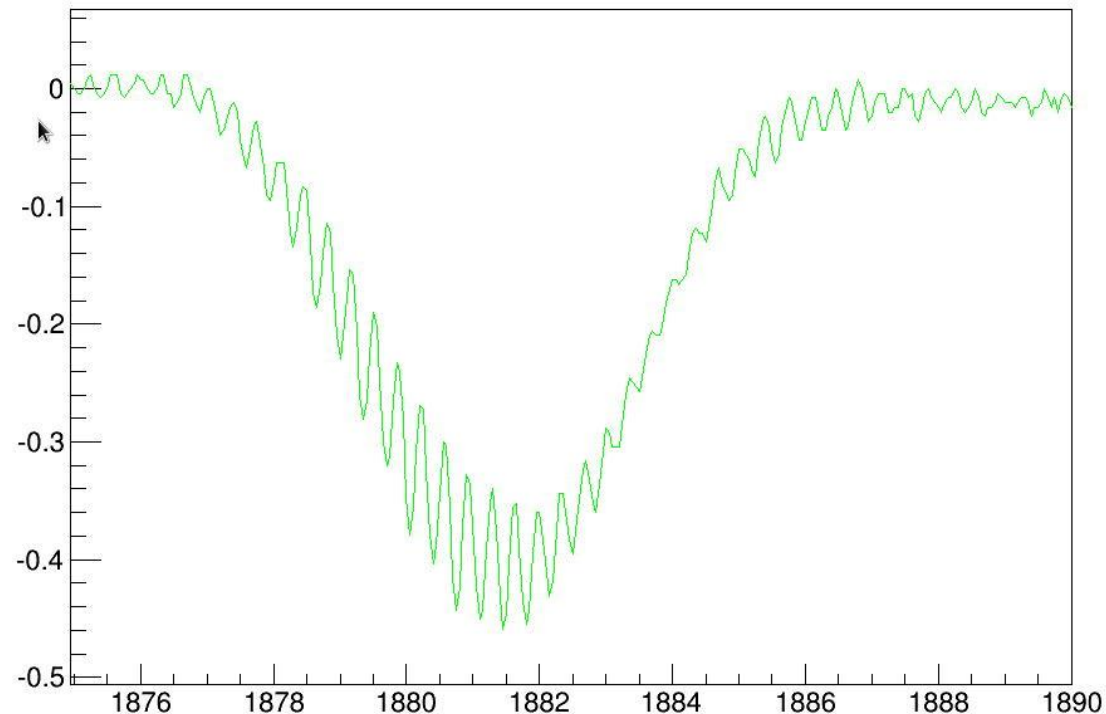
Rx Chain

Radar scattering

Beam characteristics



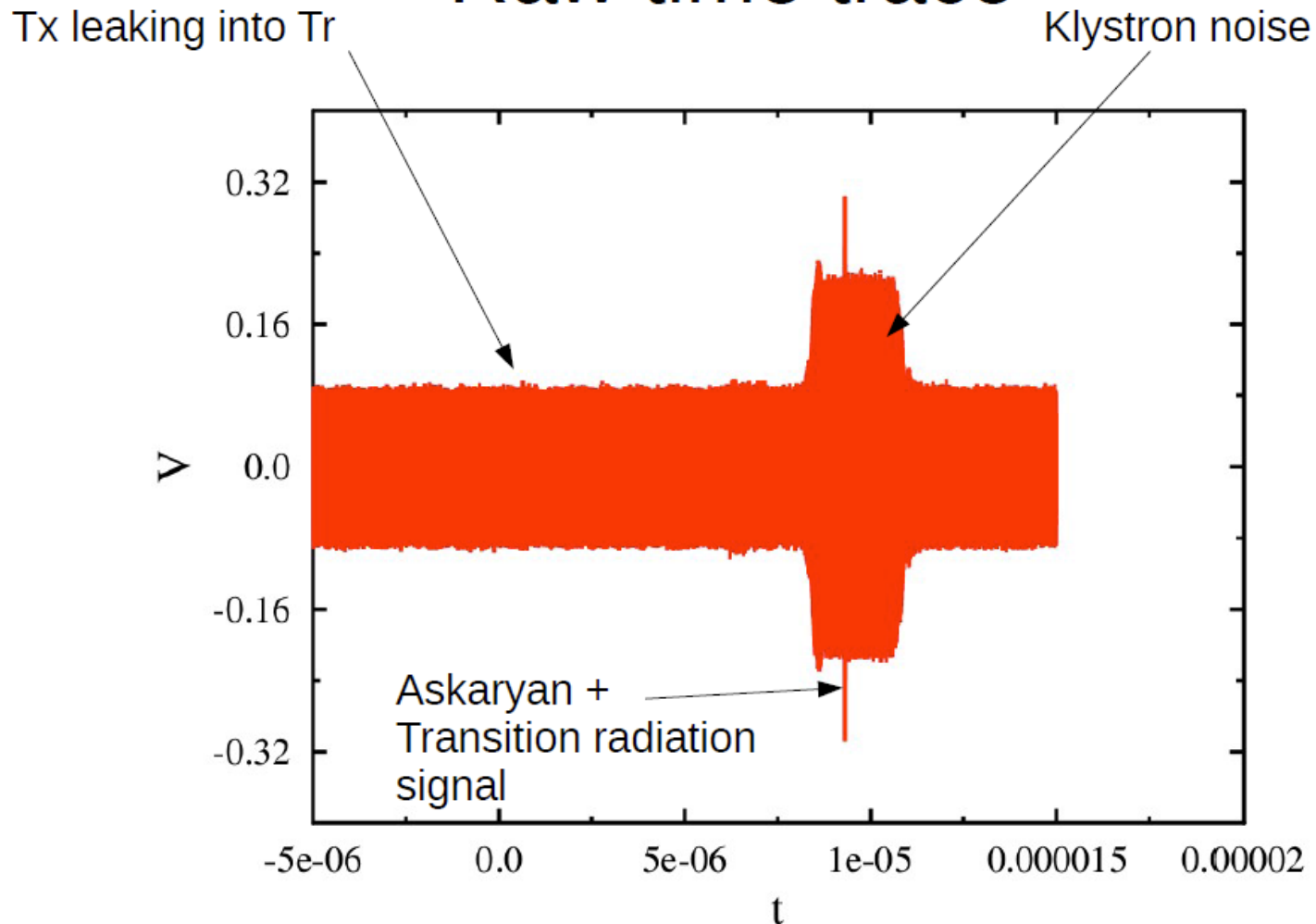
$\sim 10^9$ (40 MeV) electrons
 ~ 40 PeV



Radar scattering

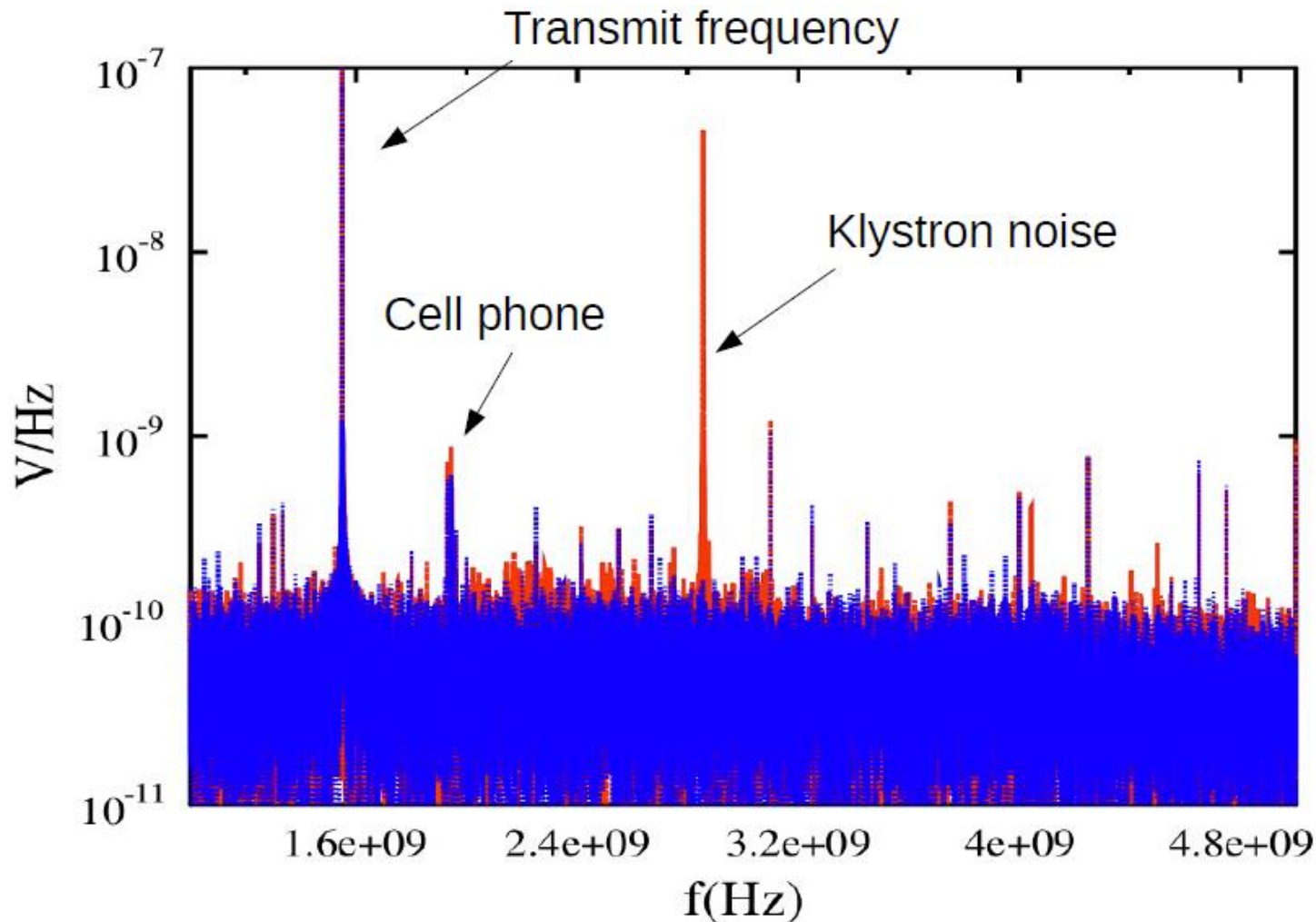
What do we see?

Raw time trace



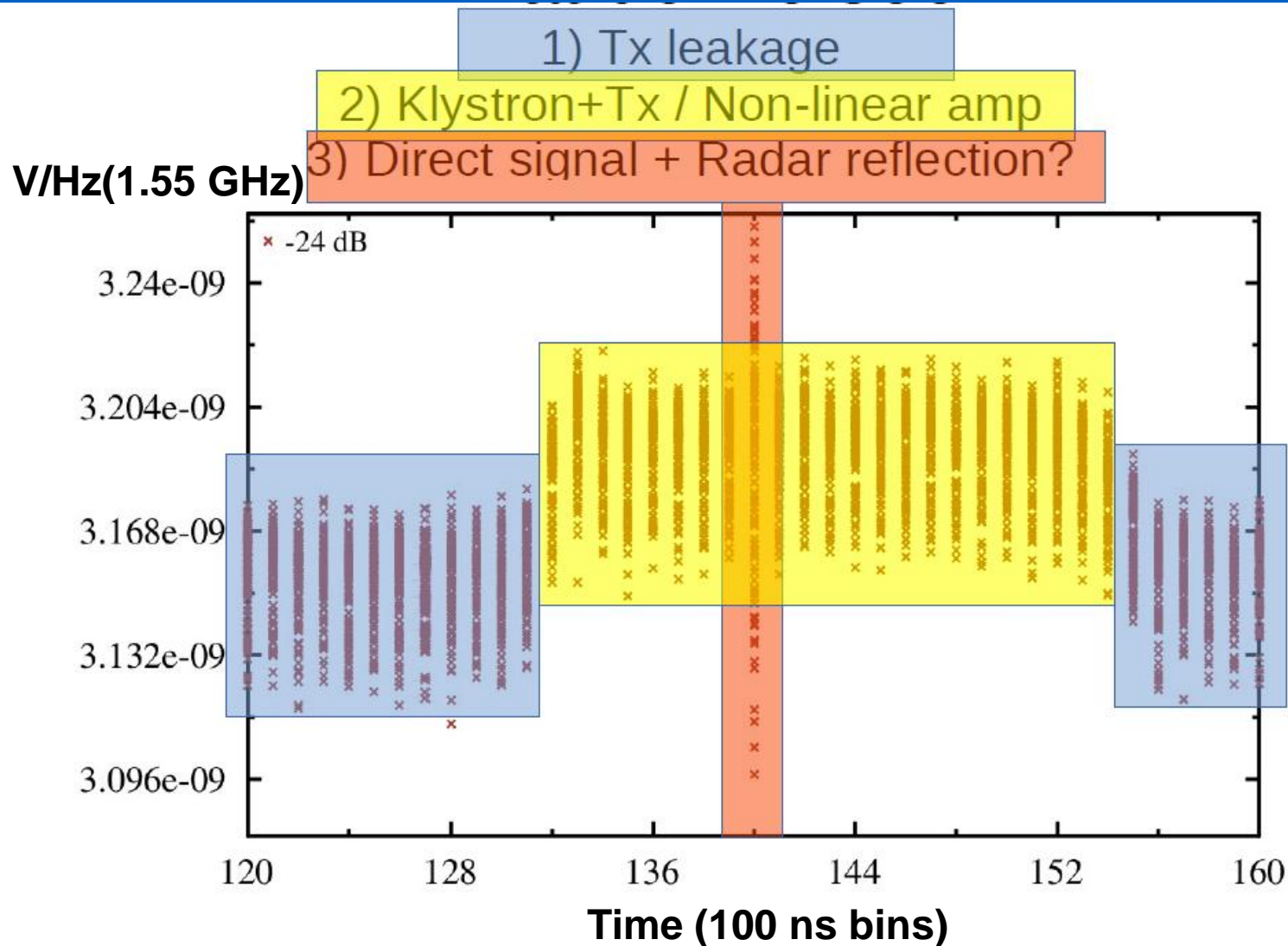
Radar scattering

What do we see?



Radar scattering

What do we see?

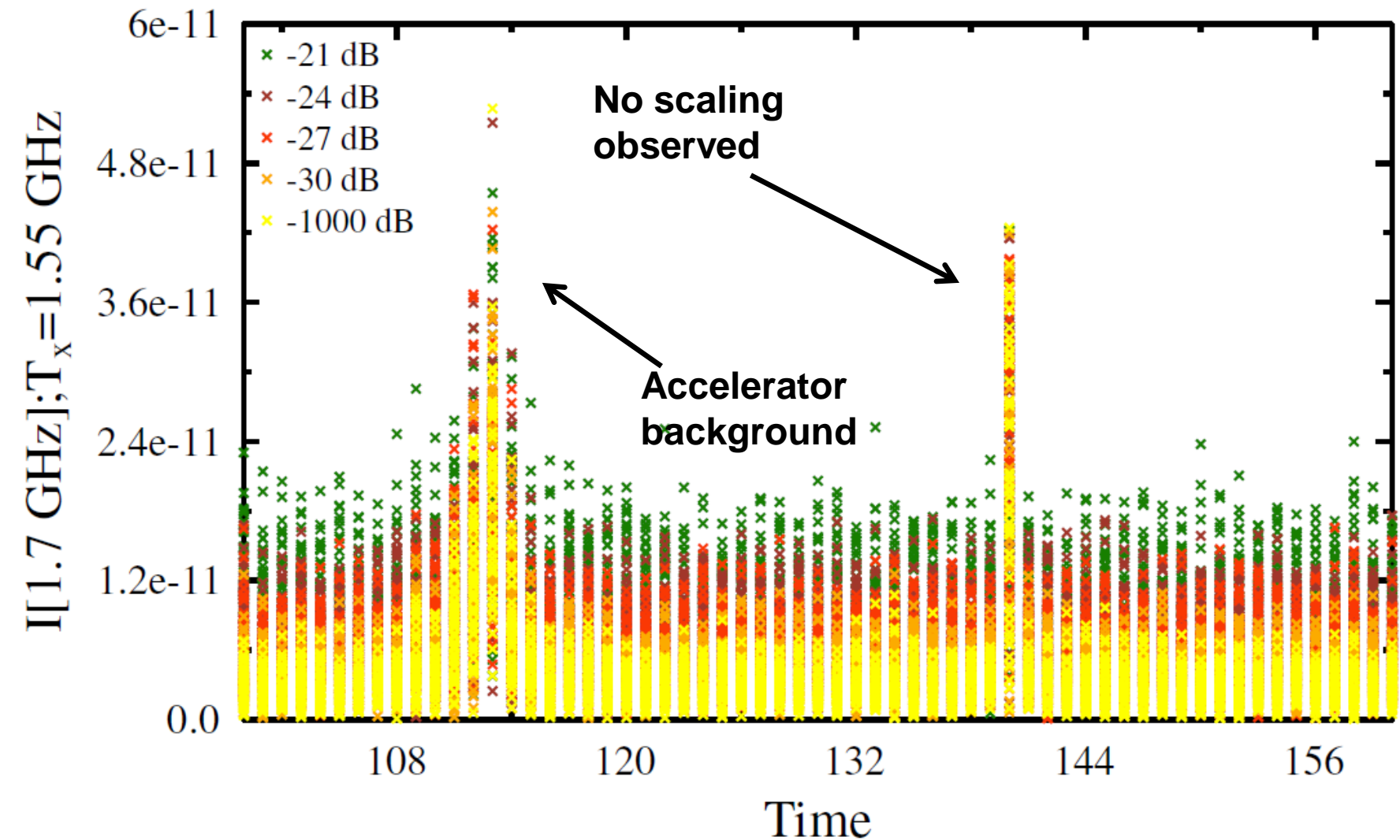


Radar scattering

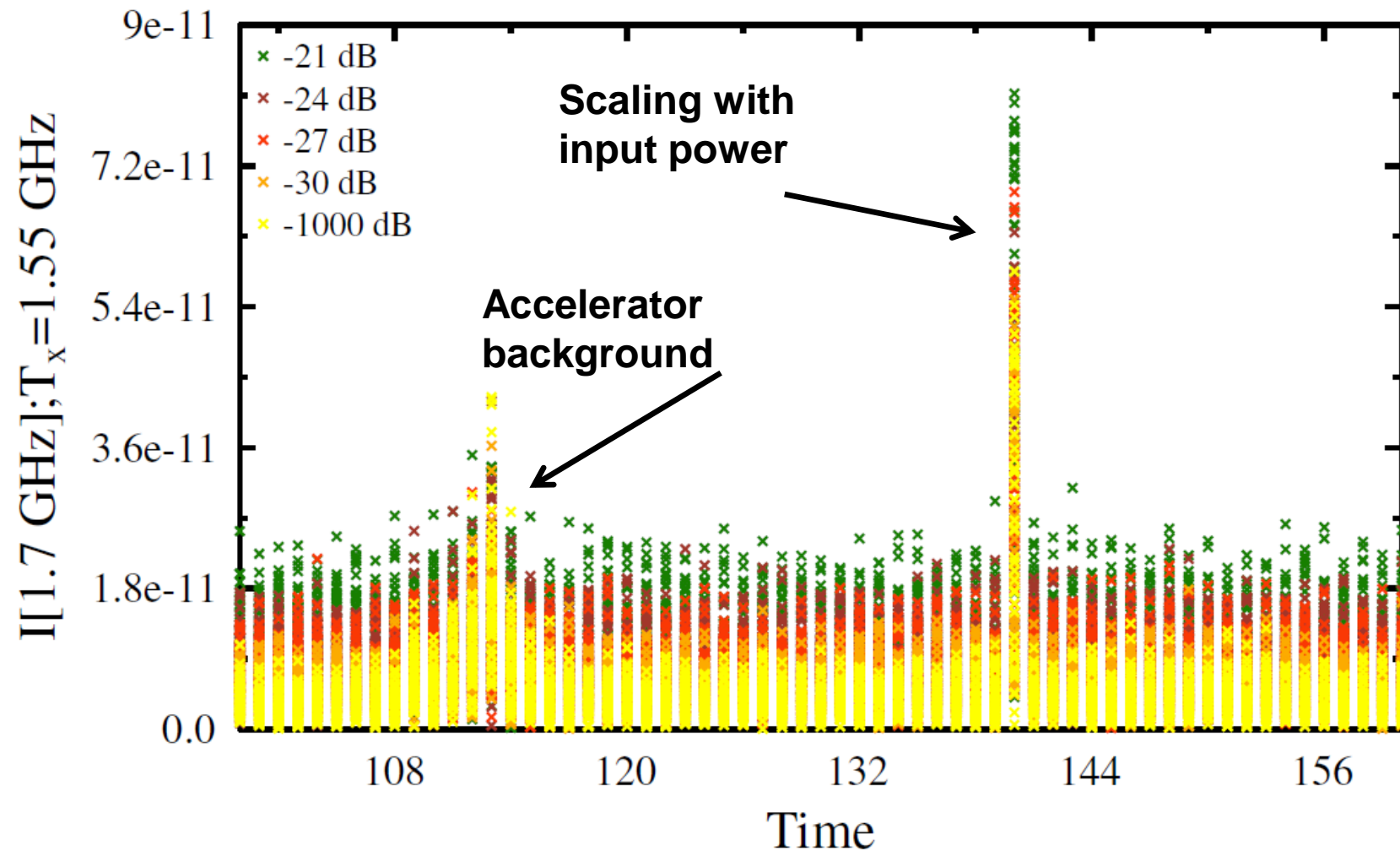
Interference and instrumental effects

- Accelerator noise interferes with our transmit signal
- Non-linear amplifier response
- **Signal can be mimicked by these effects!**
- What if we look at a different frequency than our transmit frequency?

Radar scattering Air



Radar scattering Ice



Conclusions

- Modeling the RADAR scattering of high-energy neutrino induced cascades gives an energy threshold of **several PeV**.
- We performed a measurement to determine the feasibility of this method.
- Obtained data **hints toward a scattered signal, analysis is ongoing.**

Three different types of plasma are considered

Leftover electrons from ionization:
Extension: $O(30 \text{ cm})$
Lifetime: $O(1-20 \text{ ns})$

Shower front electrons:
Extension: $R_L = O(10 \text{ cm})$
Lifetime: $O(100 \text{ ns})$
Moving!

Leftover protons from
ionization:
Wide extension: $O(5 \text{ m})$
Lifetime: $O(10-1000 \text{ ns})$

Ionization numbers come
from Physical Chemistry
research!

Figure from arXiv:1210.5140v2

6. Laws, J. O. & Parsons, D. A. *EOS* 24, 452-460 (1943)

Proton mobility in ice

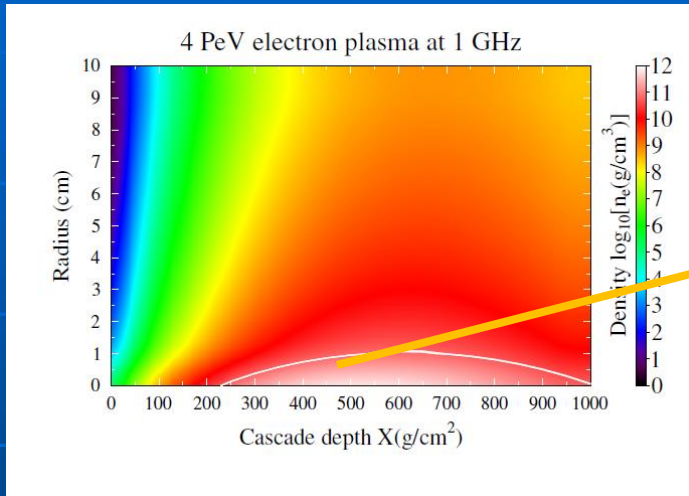
Marinus Kunst & John M. Warman

Interuniversitair Reactor Instituut, Mekelweg 15, 2629 JB Delft,
The Netherlands

Ice is frequently taken as a model when factors controlling proton transport in hydrogen-bonded molecular networks are discussed. Such discussions have increased with the acknowledgement that proton transfer across cell membranes may play a significant part in energy conversion and storage in biological systems¹⁻⁴ and that this transfer may involve hydrogen-bonded chains spanning the membrane^{5,6}. However, there is still much

Skin Effects

Model: Consider over-dense cylinders of equal density

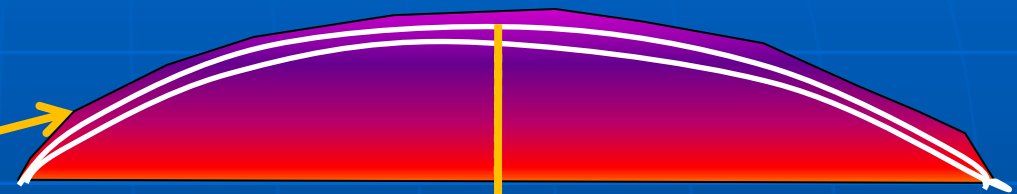


**Calculate skin depth
for a collision less plasma:**

$$\delta = \frac{c}{2\omega_p}$$

**Within 1 skin depth the
amount of power absorbed
and re-scattered equals:**

$$f_{skin}^{i+1} = (1 - f_{skin}^i) \left(1 - e^{-\frac{x}{\delta_i}}\right)$$



$$A_{Plasma}^i \approx L_i r_i$$



\approx



The over-dense radar cross-section

This approach:

1. Include skin-effects directly into the radar cross-section.
2. Consider projected area and polarization angles for in/out-going wave

$$\sigma_{od} = A_{plasma} \times f_{skin} \times f_{geom}$$

$$A_{Plasma}^i \approx L_i r_i$$

$$f_{skin}^{i+1} = (1 - f_{skin}^i)(1 - e^{-x/\delta_i})$$

$$f_{geom} = (\vec{e}_t \cdot \vec{e}_c)(\vec{e}_c \cdot \vec{e}_r)$$

$$\sigma_{od} = \sum_i L_i r_i (1 - f_{skin}^i)(1 - e^{-x/\delta_i})(\vec{e}_t \cdot \vec{e}_c)(\vec{e}_c \cdot \vec{e}_r)$$

The under-dense radar cross-section

The wave will scatter off of the individual electron given by the Thompson cross-section

$$\sigma_T = \left(\frac{m_e}{m_p} \right)^2 0.665 \cdot 10^{-28} \text{ m}^2$$

We have to take into account for the phase lag of the individual electrons w.r.t. each other:

$$\sigma_{ud} = \sum_{i=1}^N \sigma_T \cos(kx)$$

$$k = \frac{2\pi}{\lambda_d} \quad x = |\vec{x}_1 - \vec{x}_i| + |\vec{x}_2 - \vec{x}_i|$$