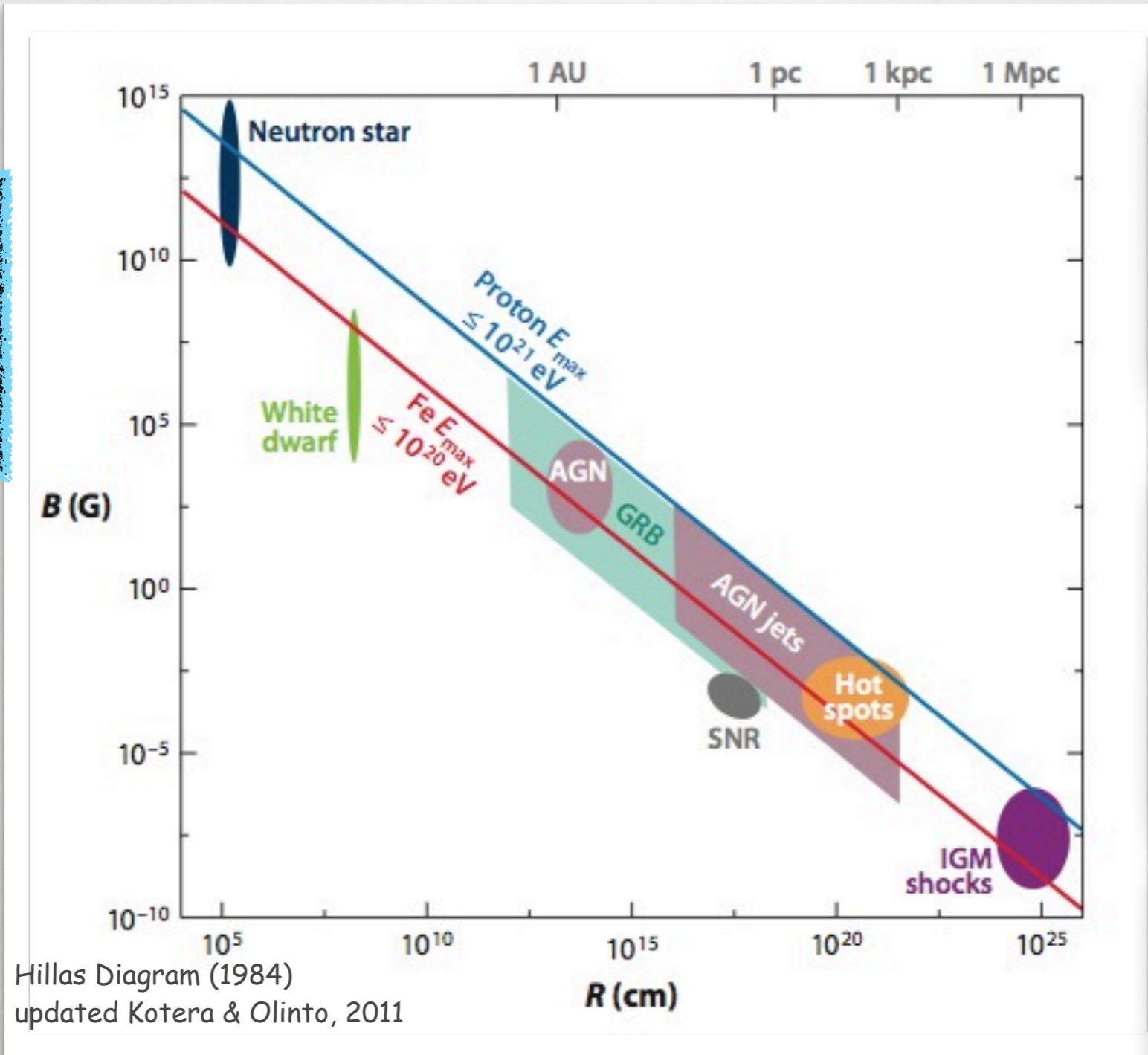
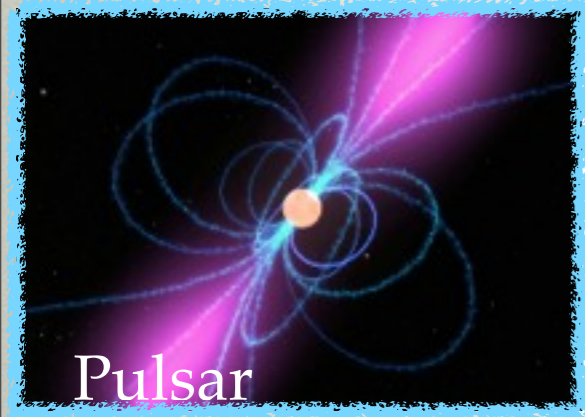


# Newborn Pulsars as Sources of UHECRs

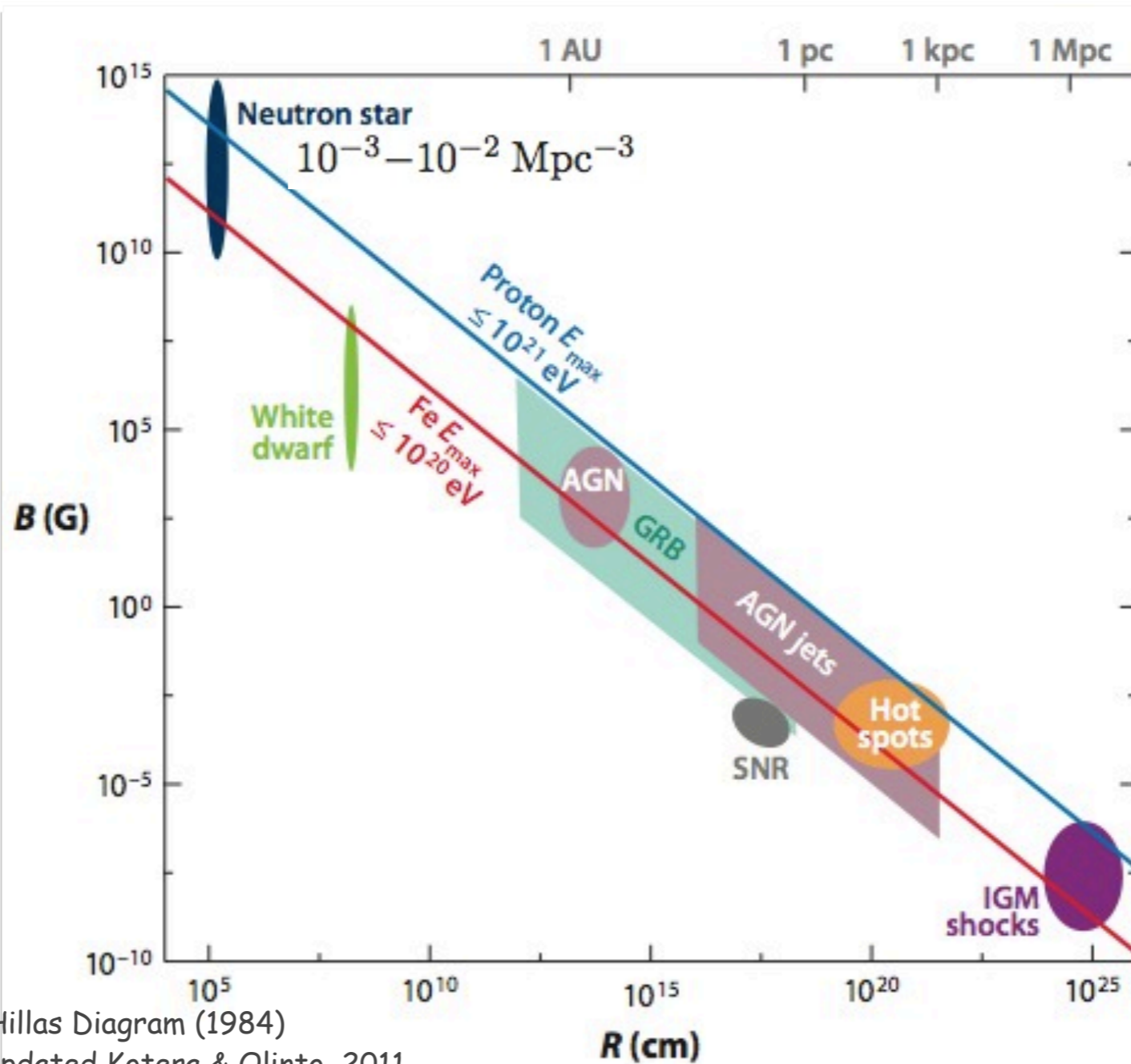
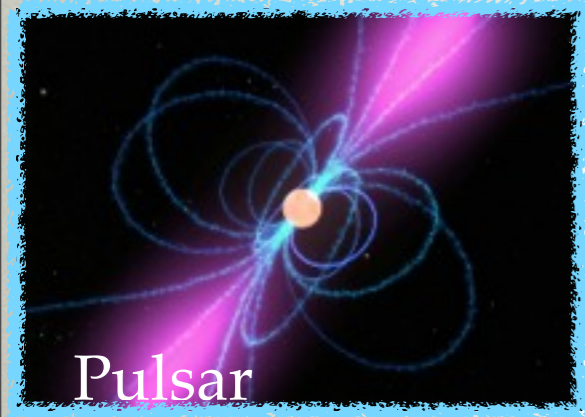
Ke Fang  
University of Chicago

Cosmic Ray Anisotropy Workshop  
Sep 26 2013

# Possible Candidates of UHECR Sources



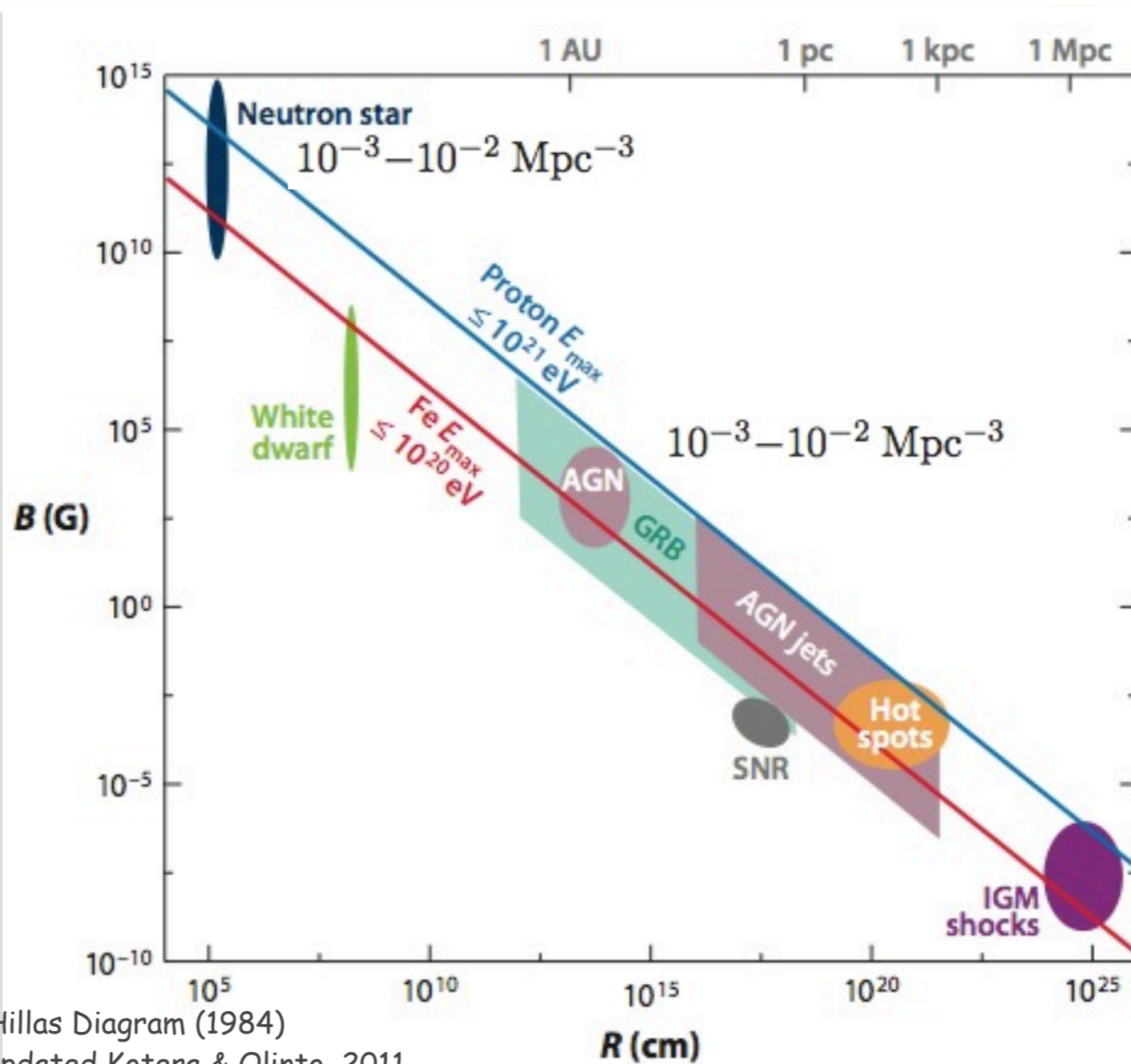
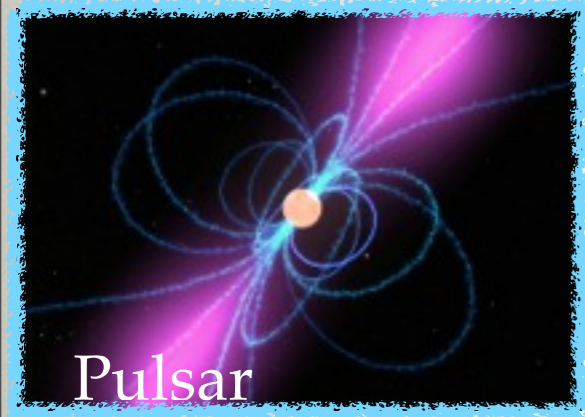
# Possible Candidates of UHECR Sources



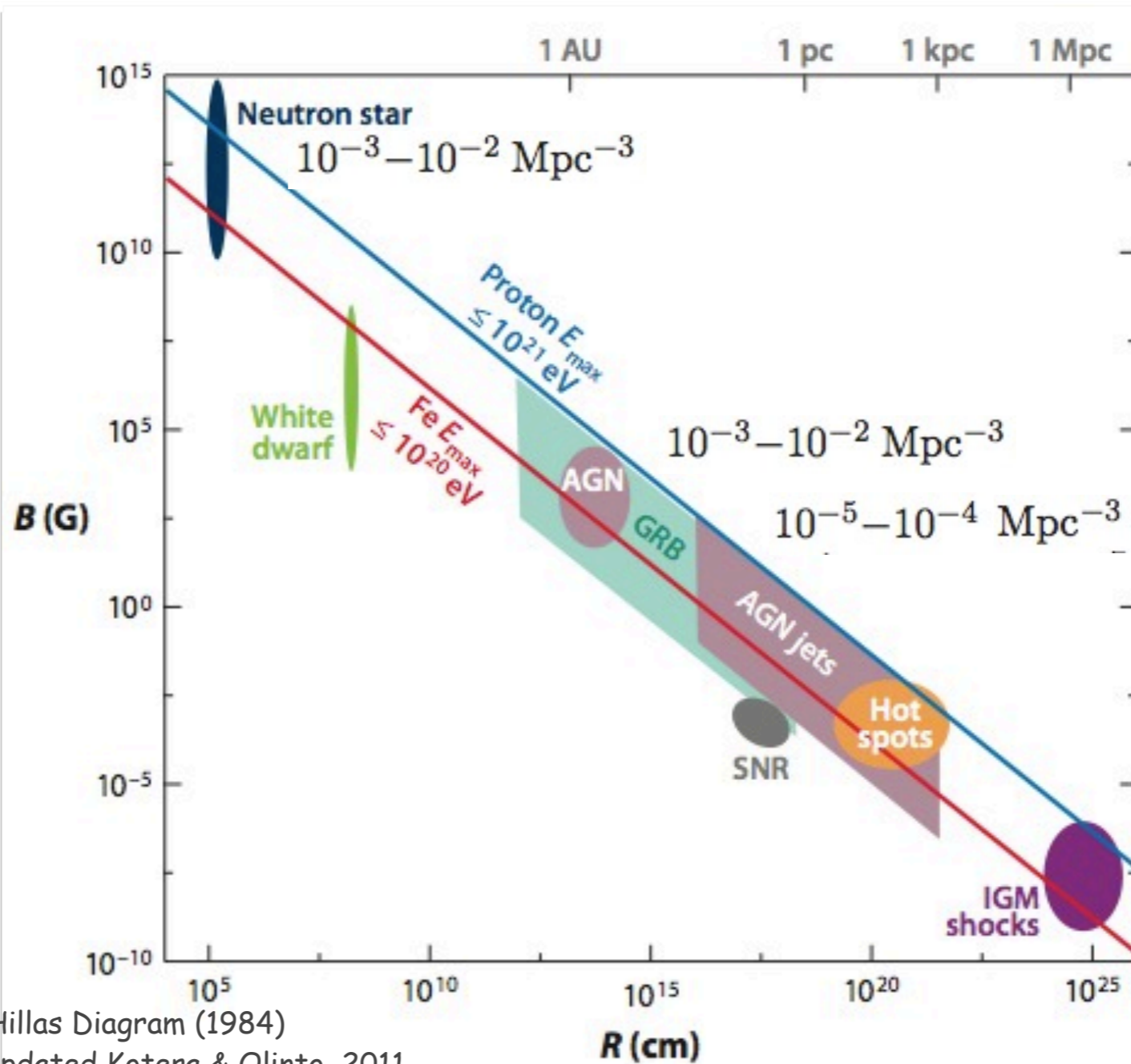
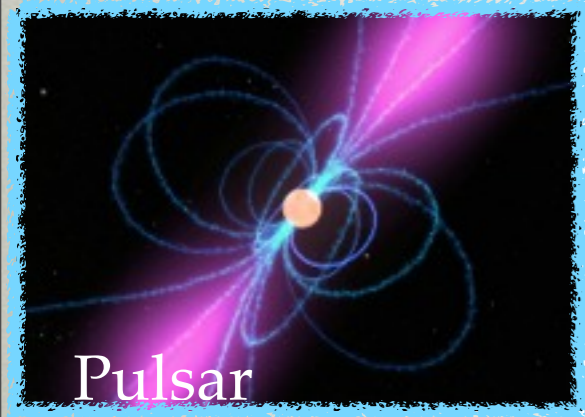
Hillas Diagram (1984)  
updated Kotera & Olinto, 2011



# Possible Candidates of UHECR Sources



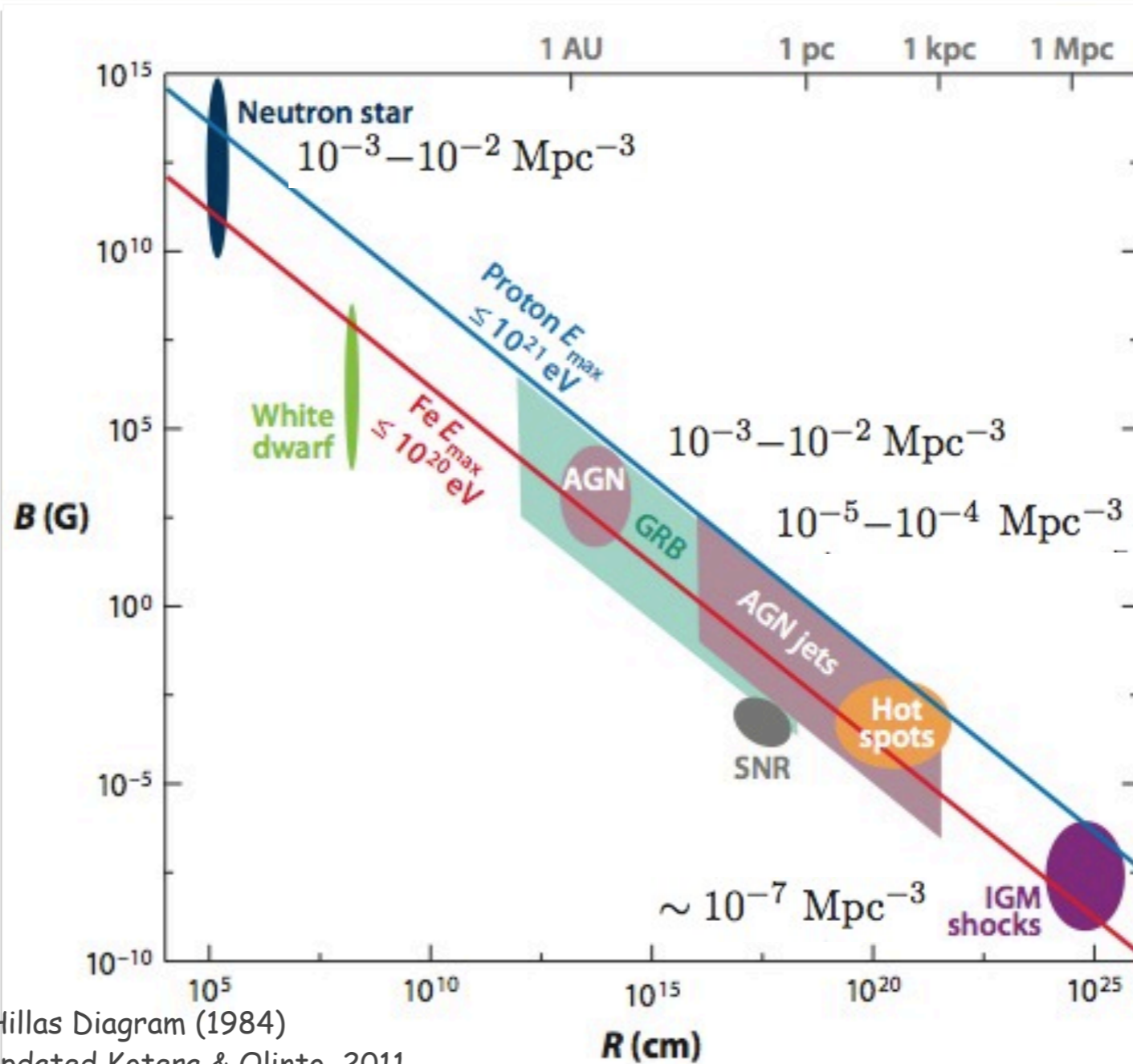
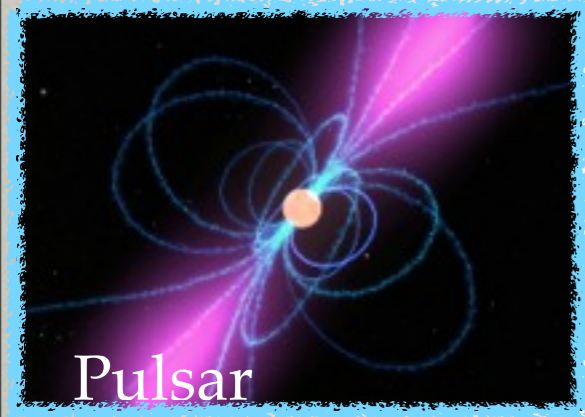
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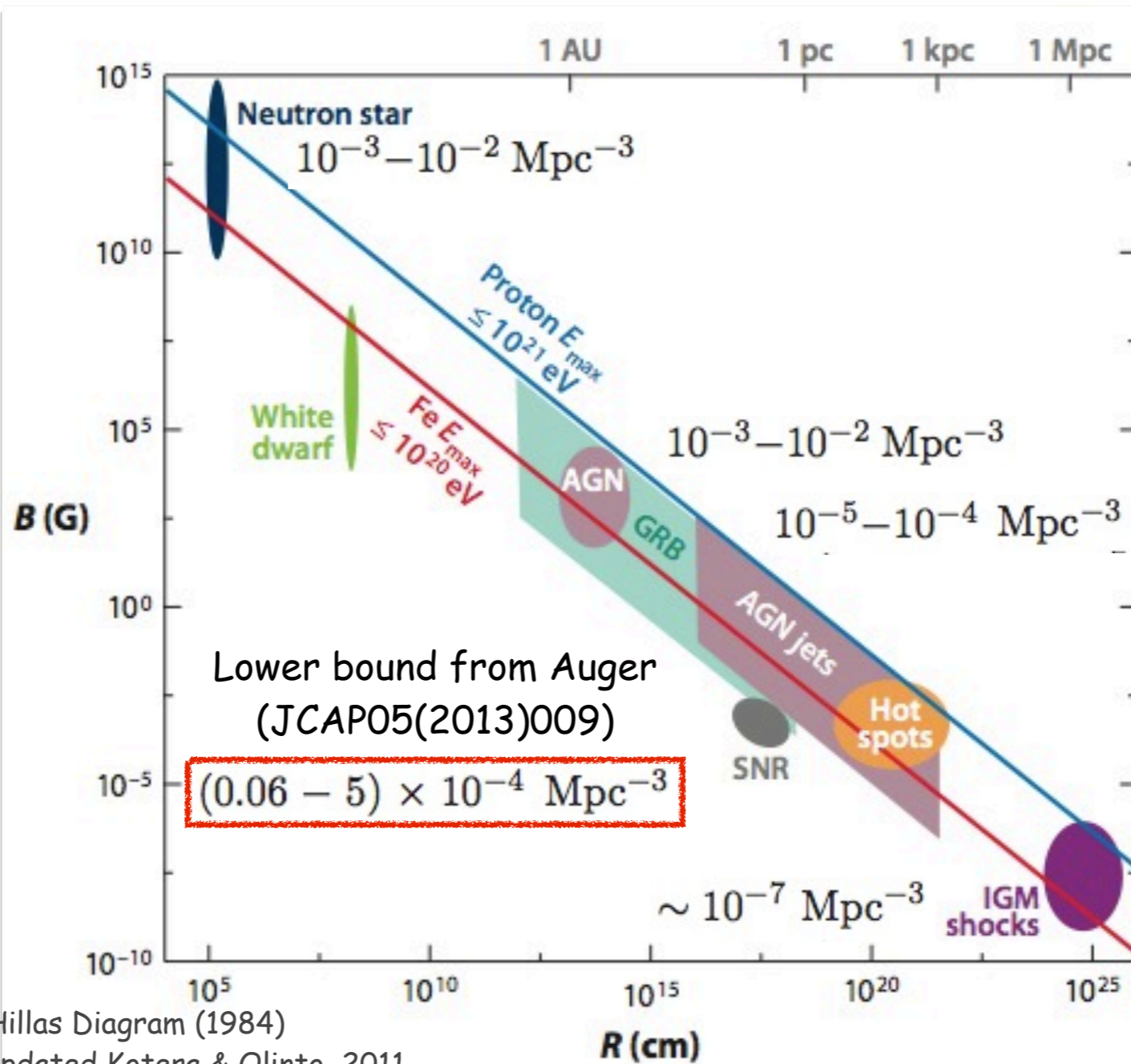
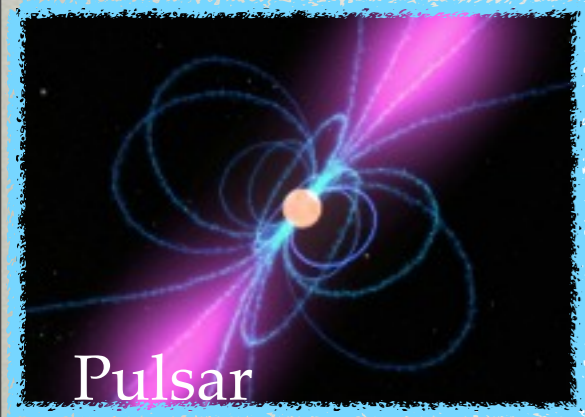
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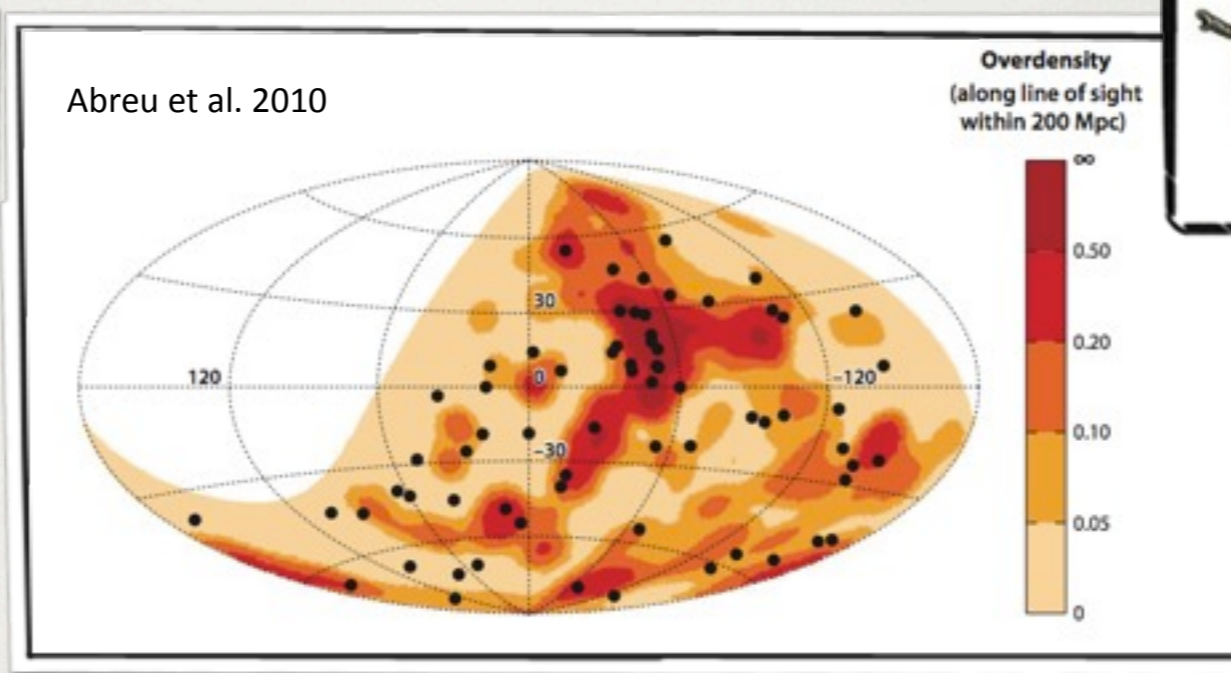
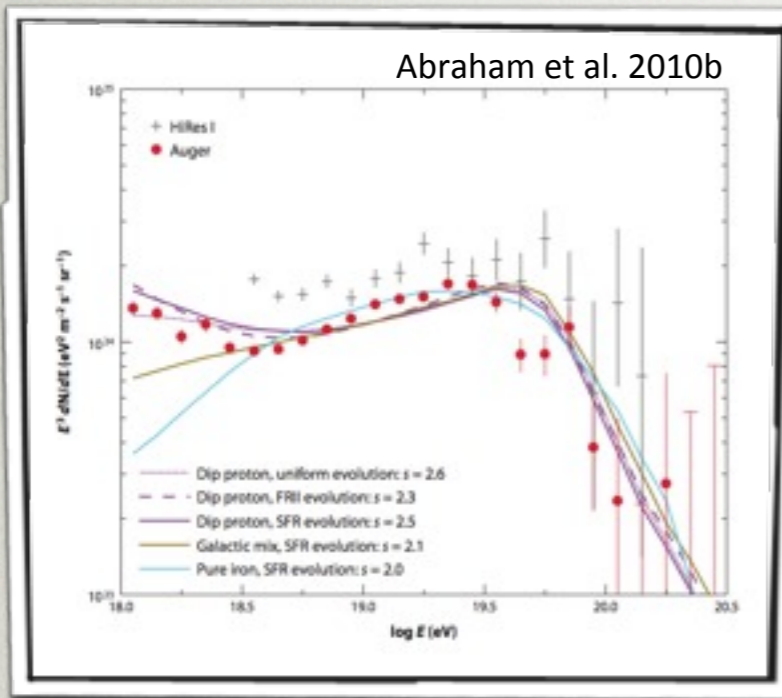
# Possible Candidates of UHECR Sources



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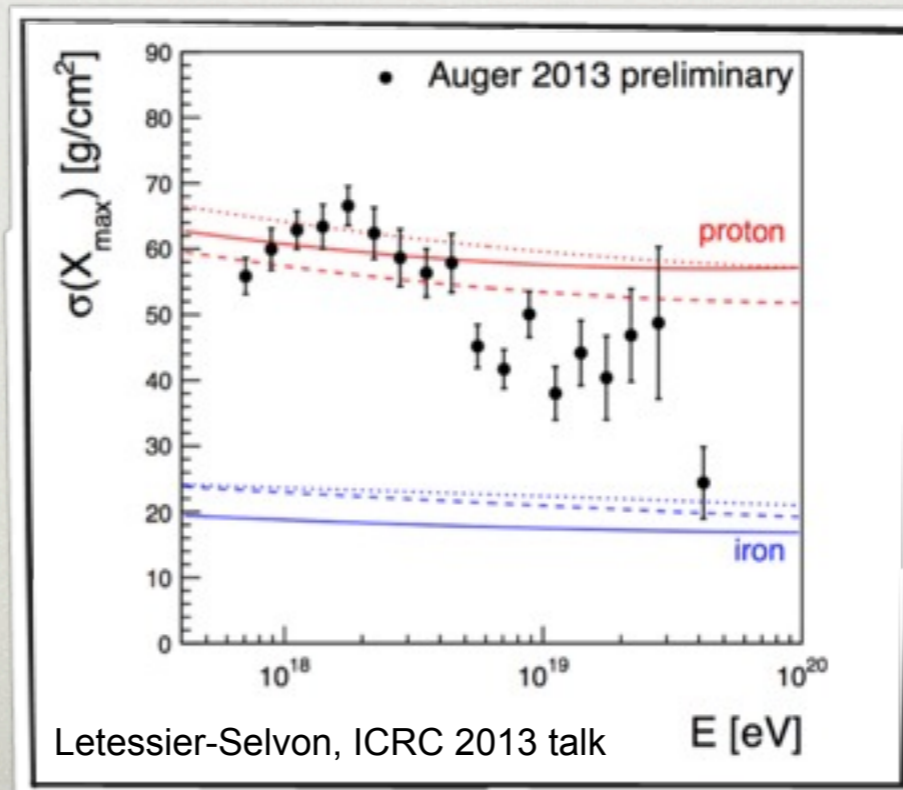
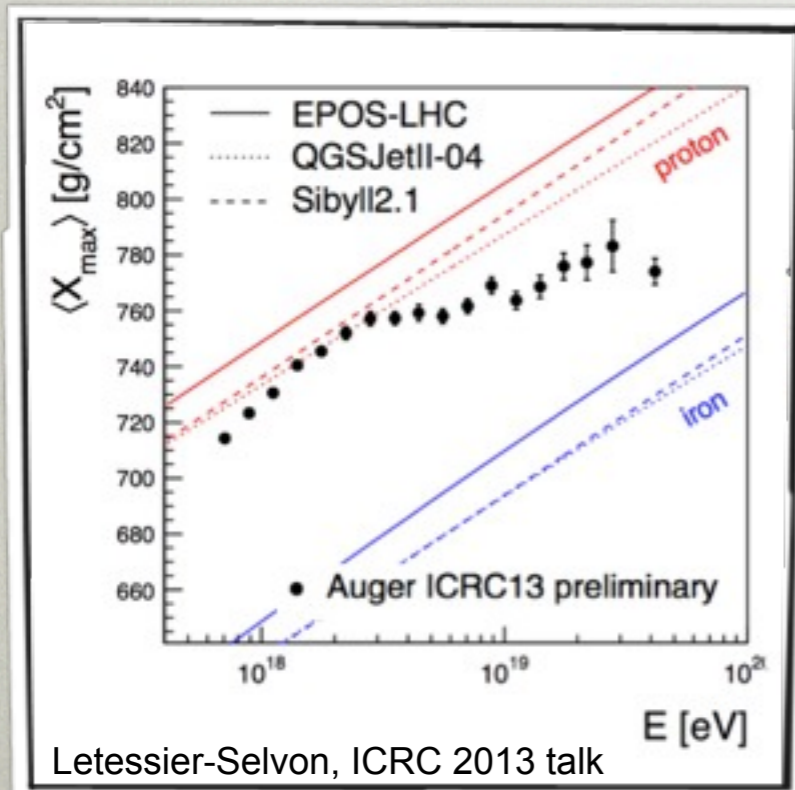


# The Source has to explain..



Intrinsic spectrum index  $\sim 2$

no significant anisotropy



Intermediate and heavy above  $10 \text{ EeV}$   
TA indicates light composition

...  
Multi-messenger signatures



# A tale of newborn pulsars

Blasi, Epstein & Olinto 2000  
Arons 2003  
KF, Kotera, Olinto 2012, 2013

Goldreich-Julian  
charge density at  
the stellar surface

$$\dot{N}_{GJ} = \frac{\Omega^2 \mu}{Zec}$$

Pulsar spins down due  
to electromagnetic  
radiation (neglect GW)

$$\dot{\Omega} = -\frac{\dot{E}_{EM}}{I\Omega} \propto -\mu^2 \Omega^3$$

Particles can be accelerated by the induced  
E-field

$$E = Ze\Phi\eta = 3 \times 10^{20} Z_{26} \eta_1 \Omega_{41}^2 \mu_{30.5} eV$$

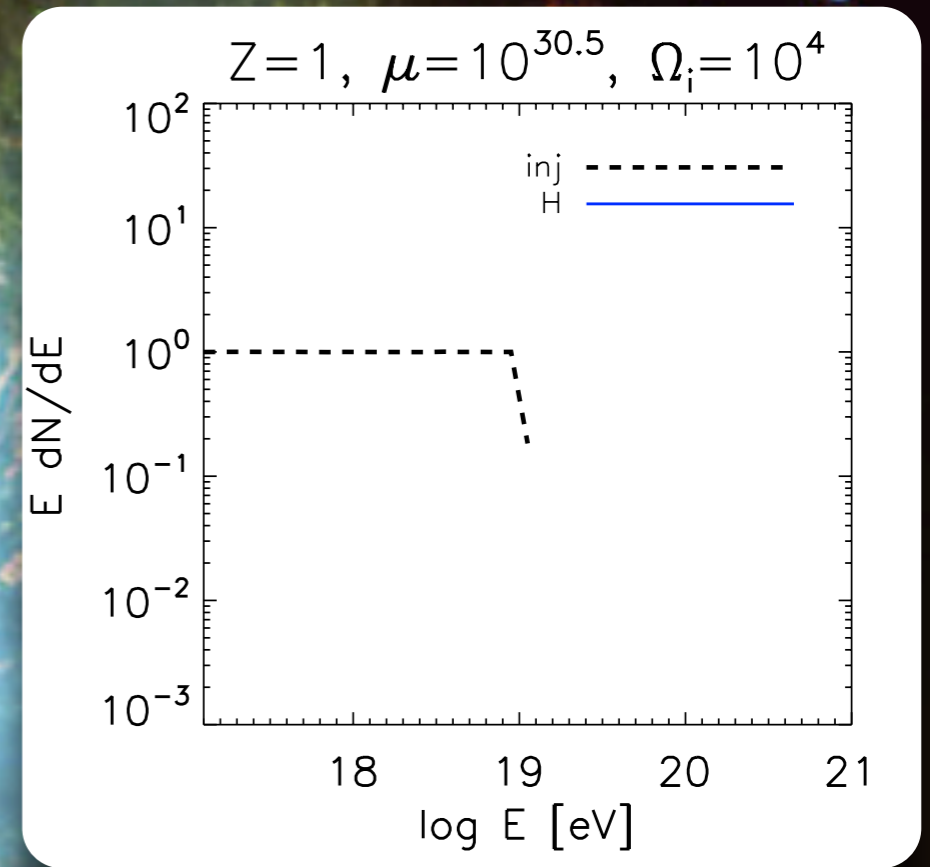
$$t_{spin}(E) = 1 \text{ yr} \left( \frac{3 \times 10^{20} eV}{E} \right) \frac{Z_{26} \eta_1}{\mu_{30.5}}$$

$$\frac{dN_i}{dE} = 5 \times 10^{23} (Z_{26} \mu_{30.5} E_{20})^{-1} eV^{-1}$$

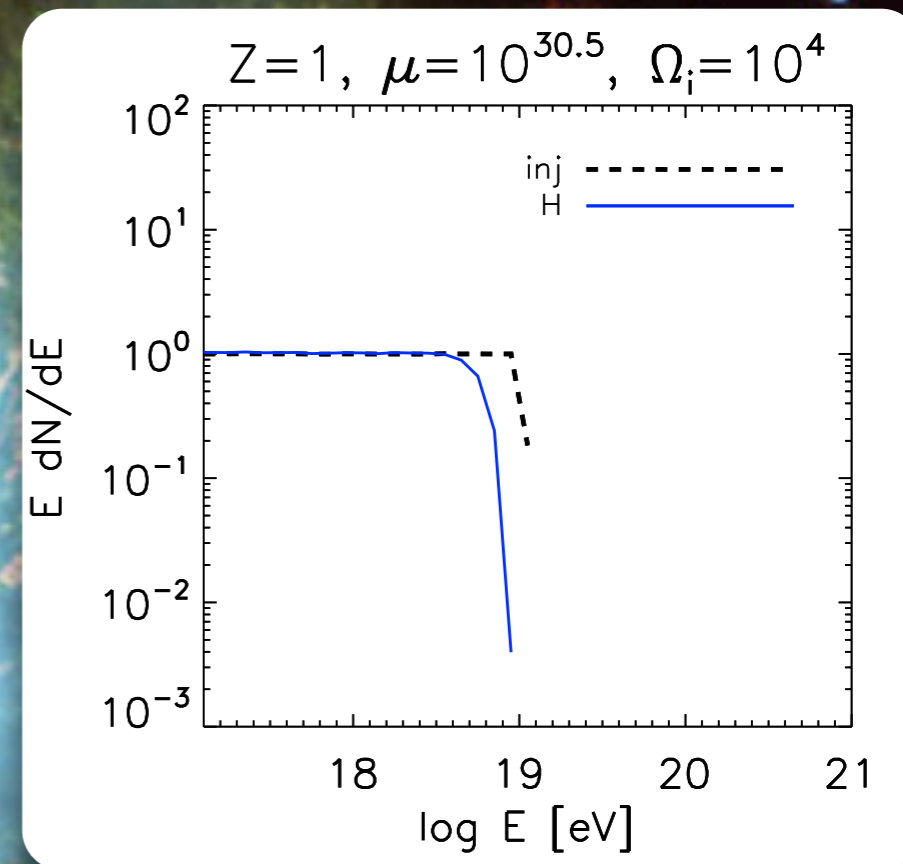
Monte-Carlo propagation  
hadron interactions simulated  
with EPOS + CONEX



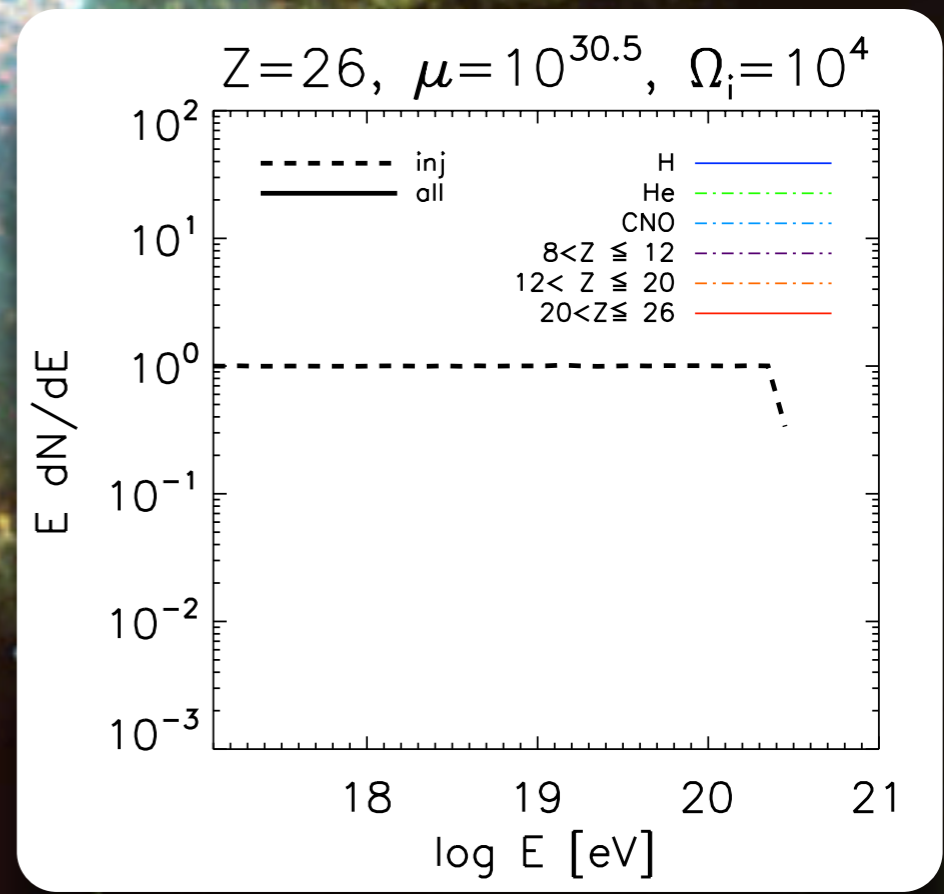
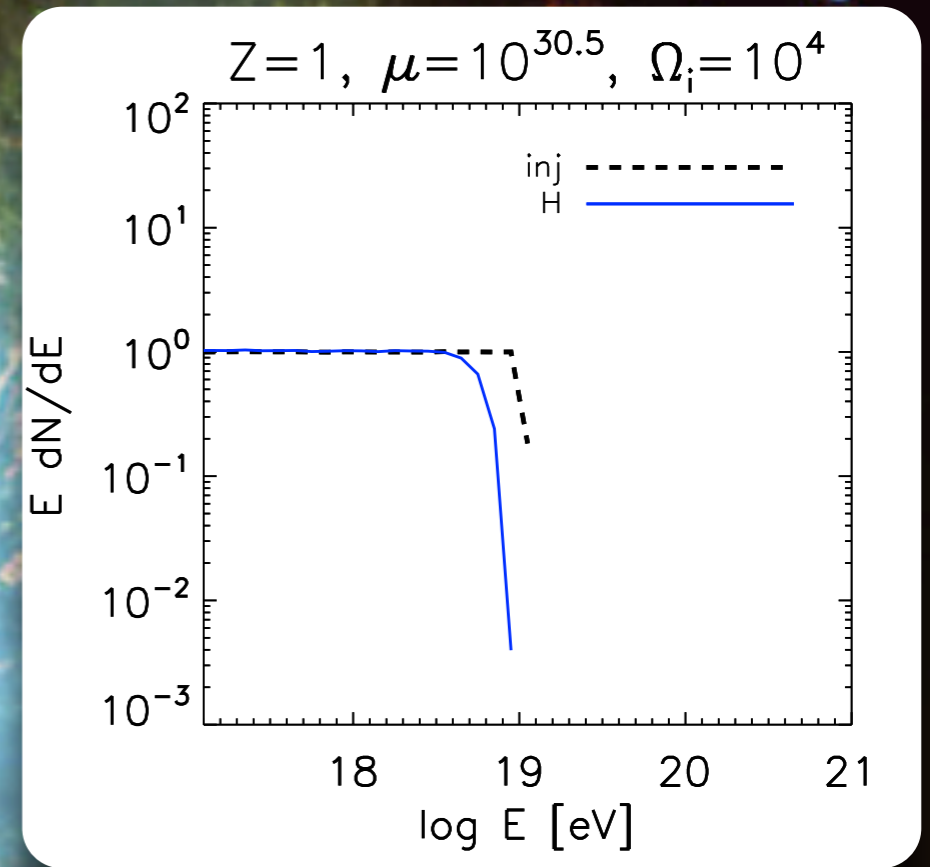
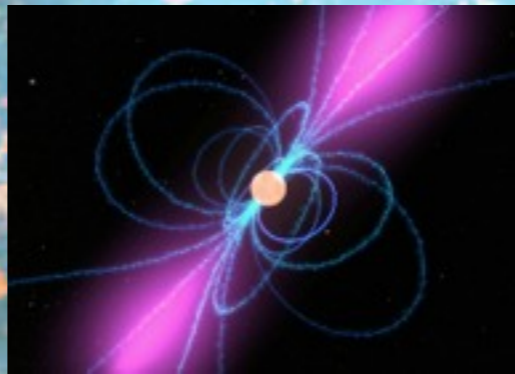
Monte-Carlo propagation  
hadron interactions simulated  
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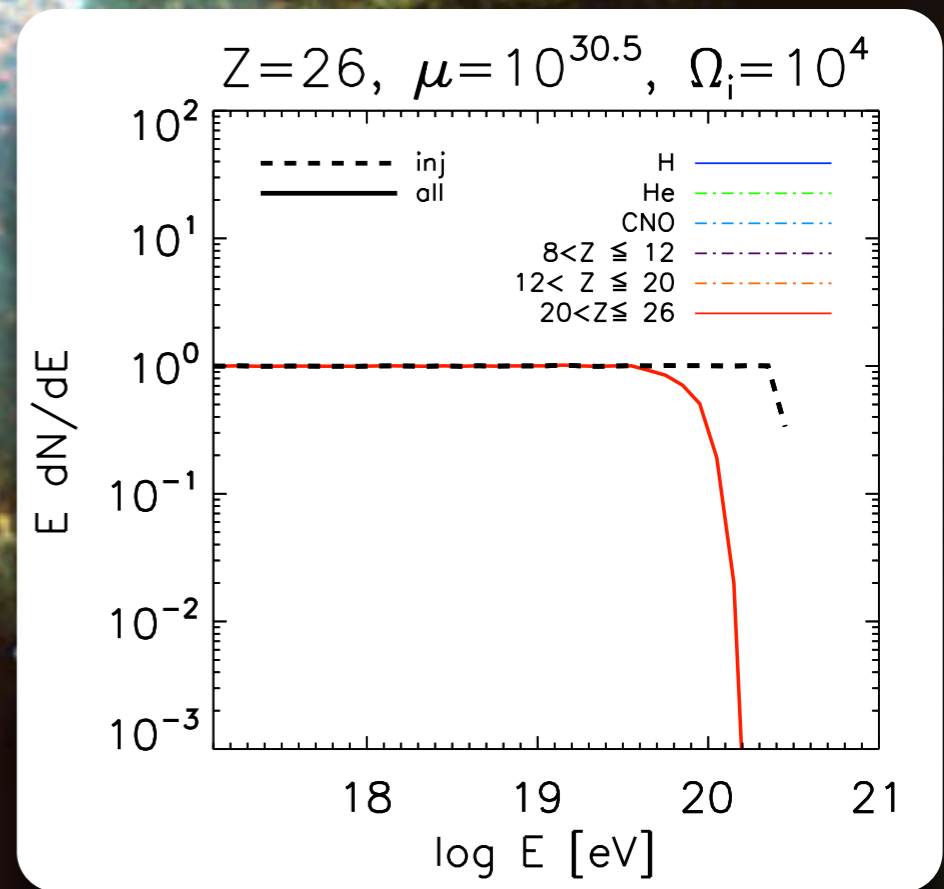
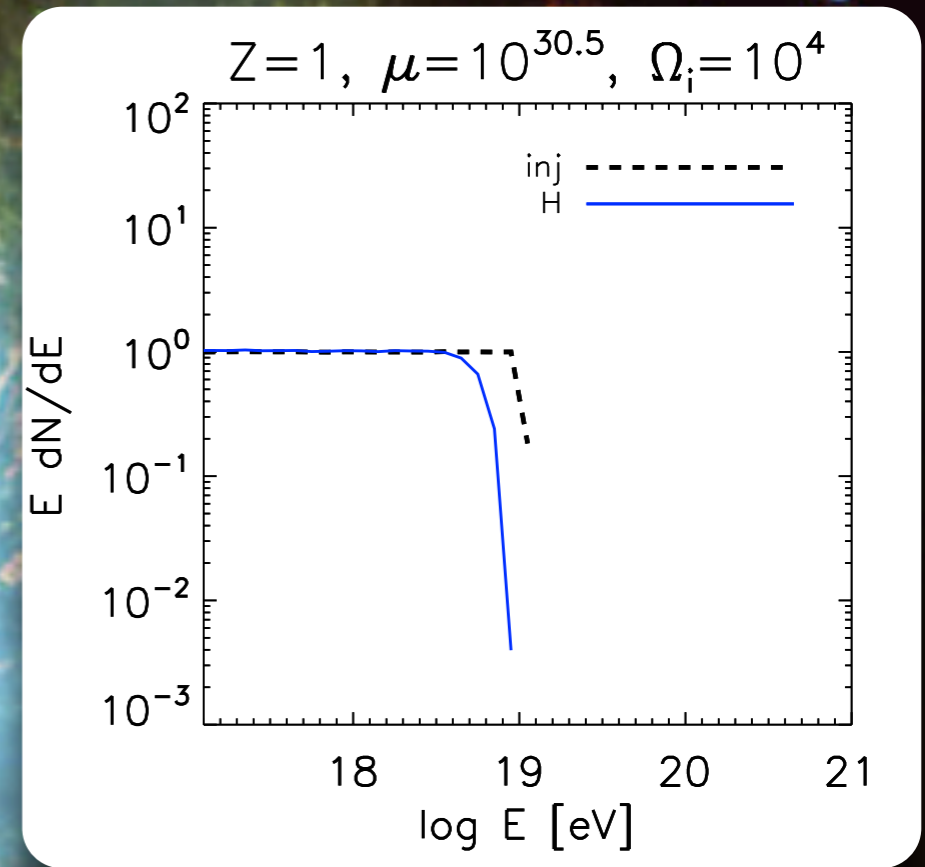
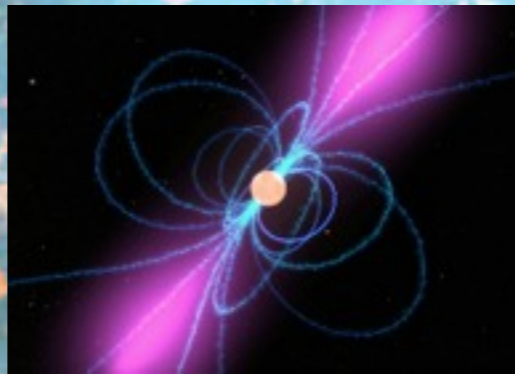
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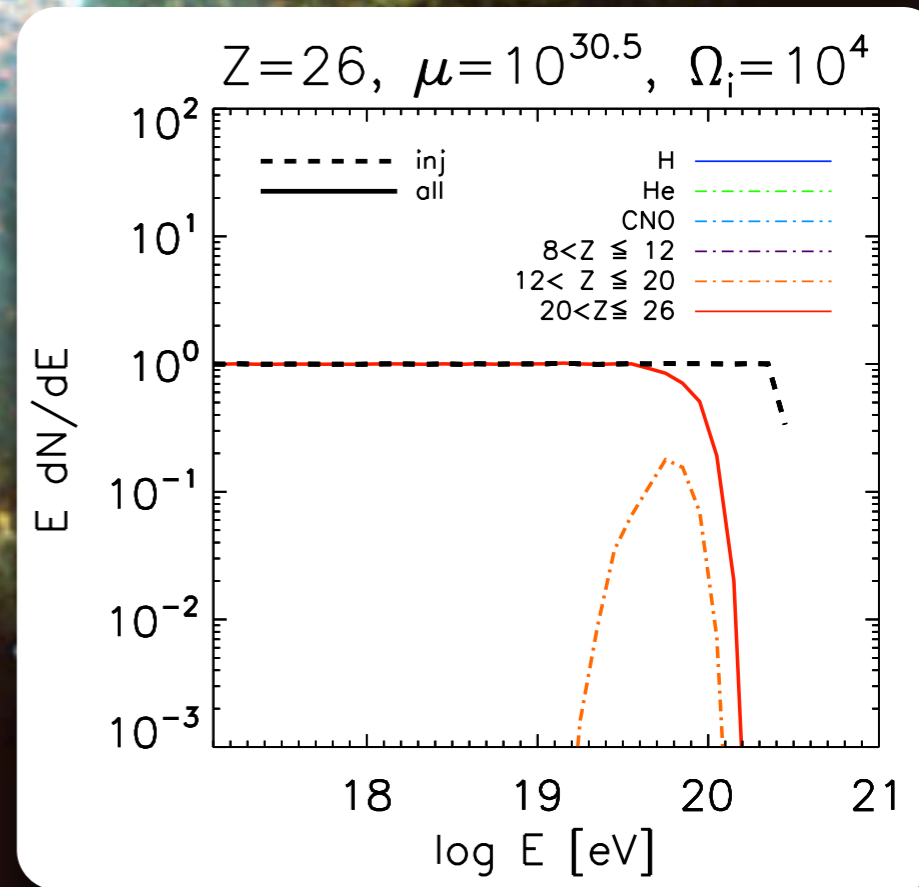
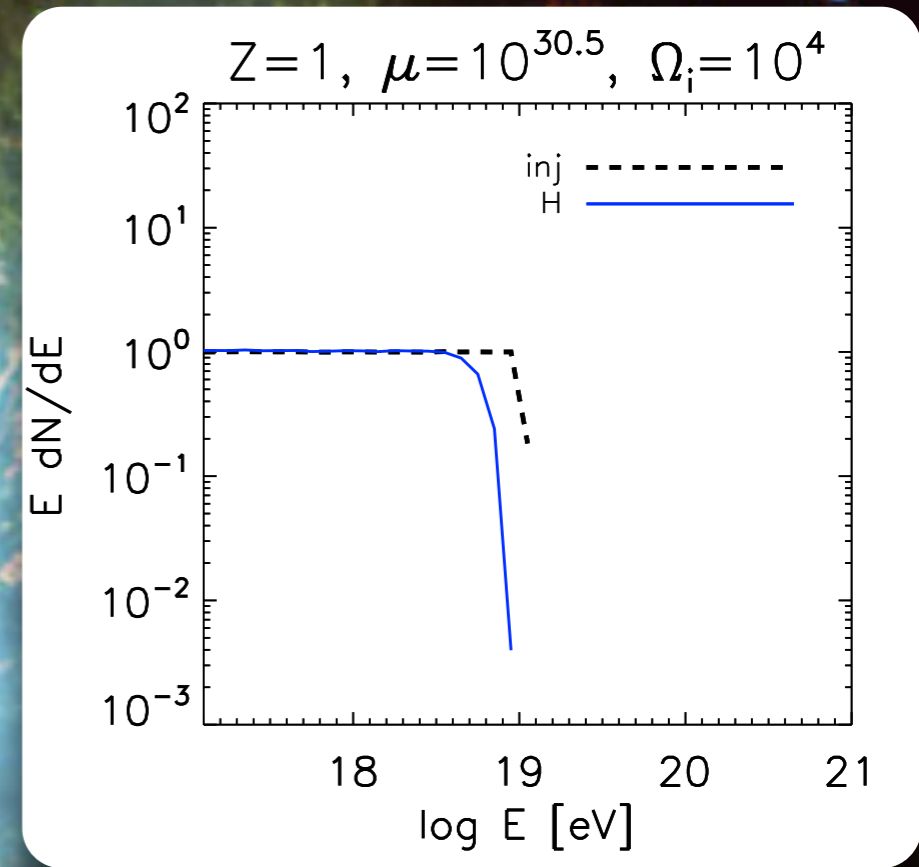
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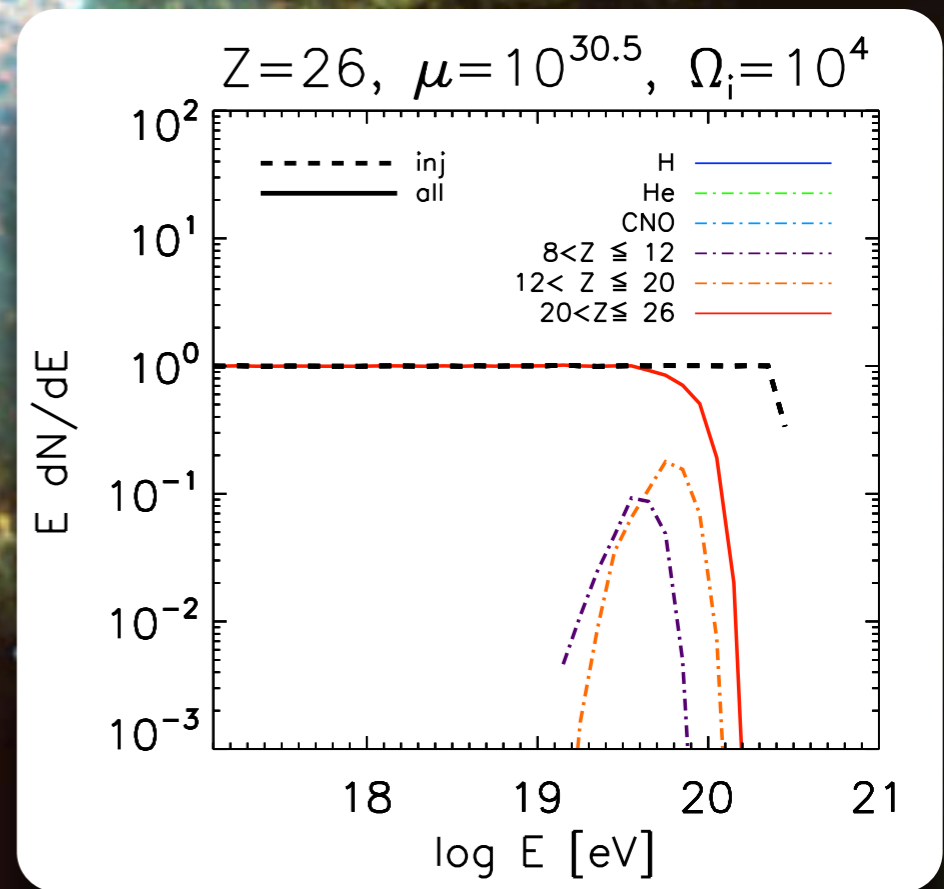
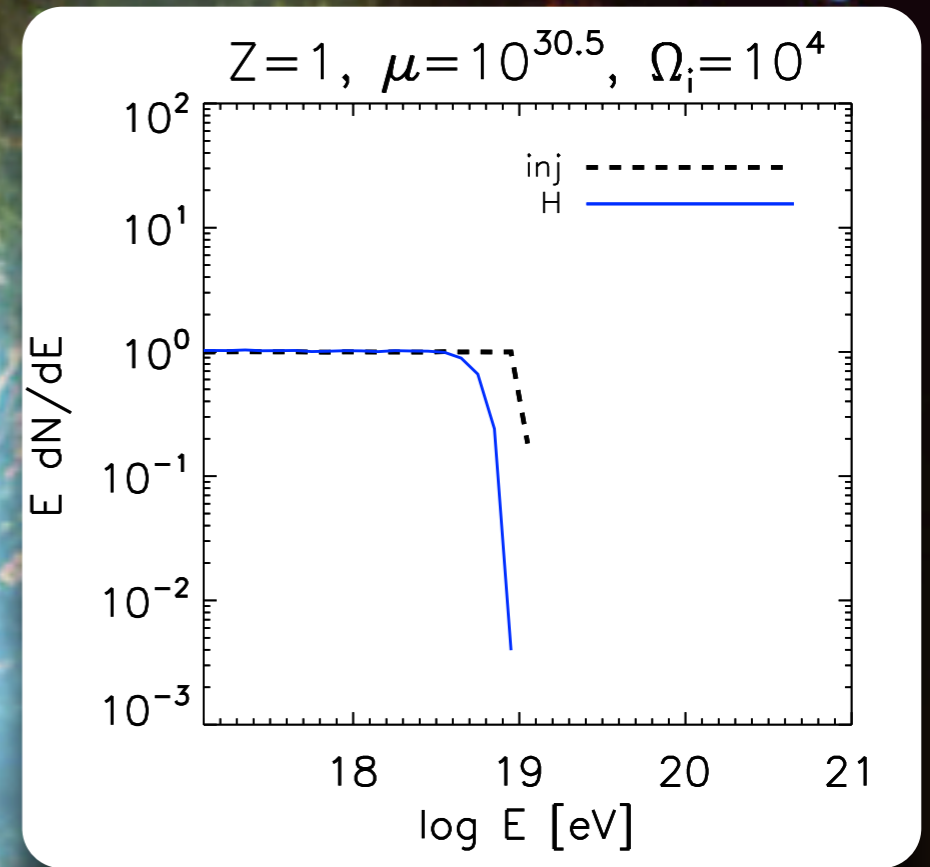
Monte-Carlo propagation  
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Monte-Carlo propagation  
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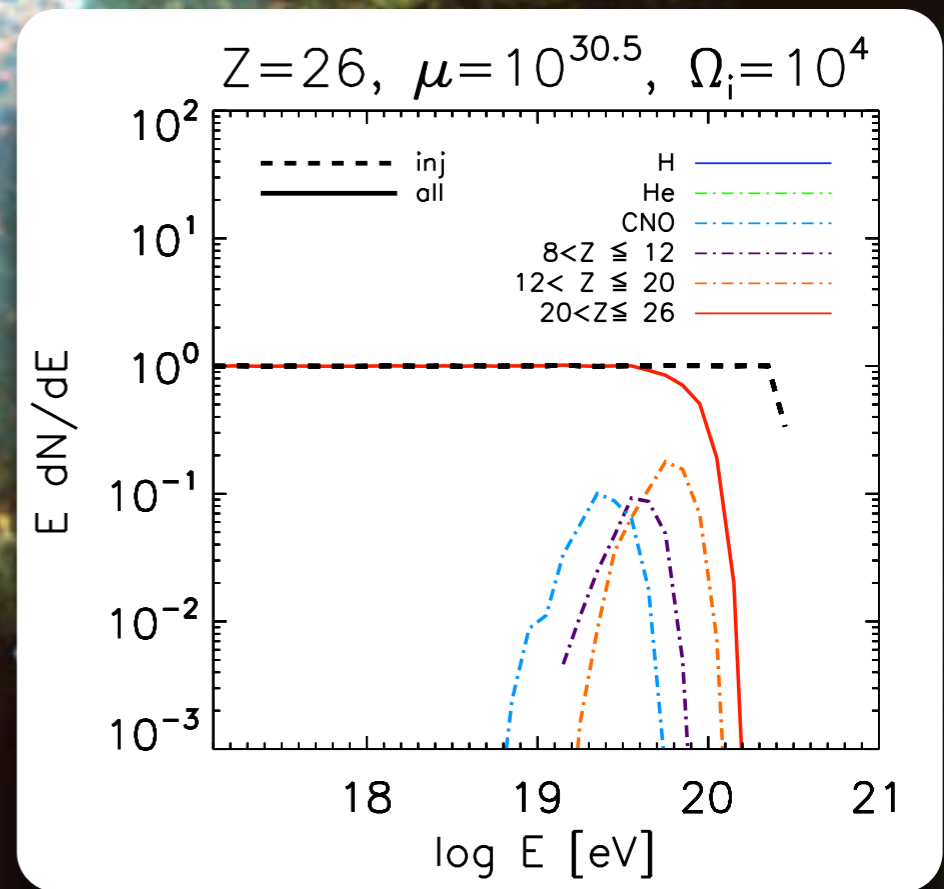
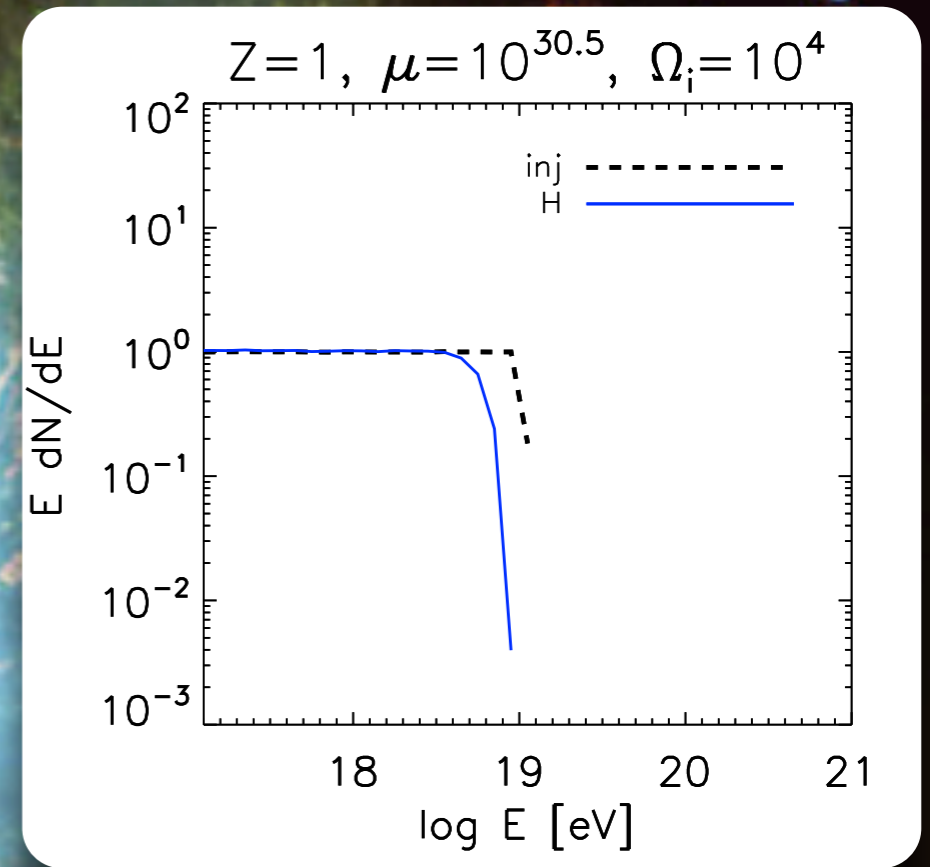
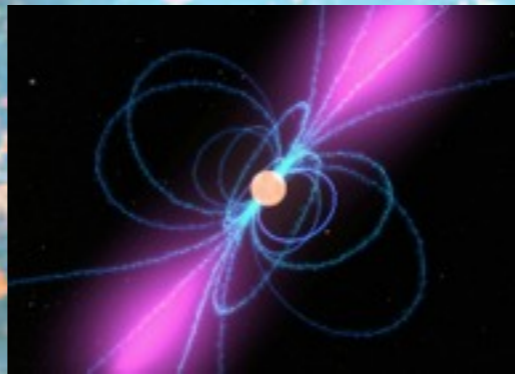


Monte-Carlo propagation  
hadron interactions simulated  
with EPOS + CONEX

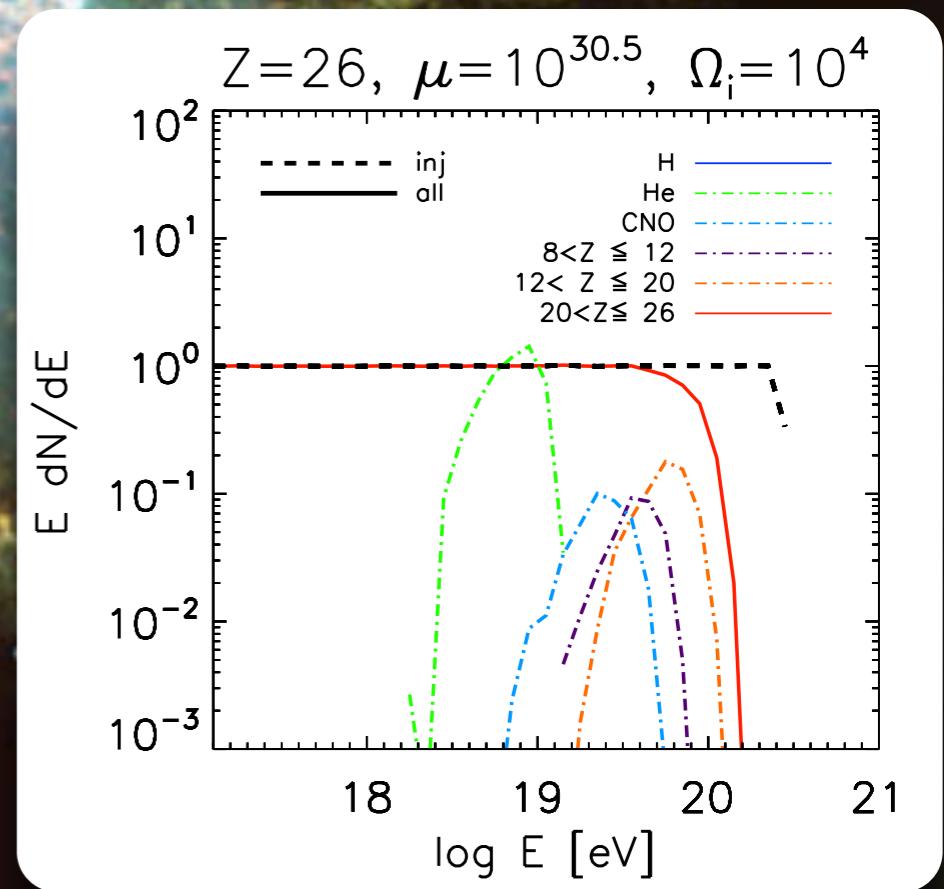
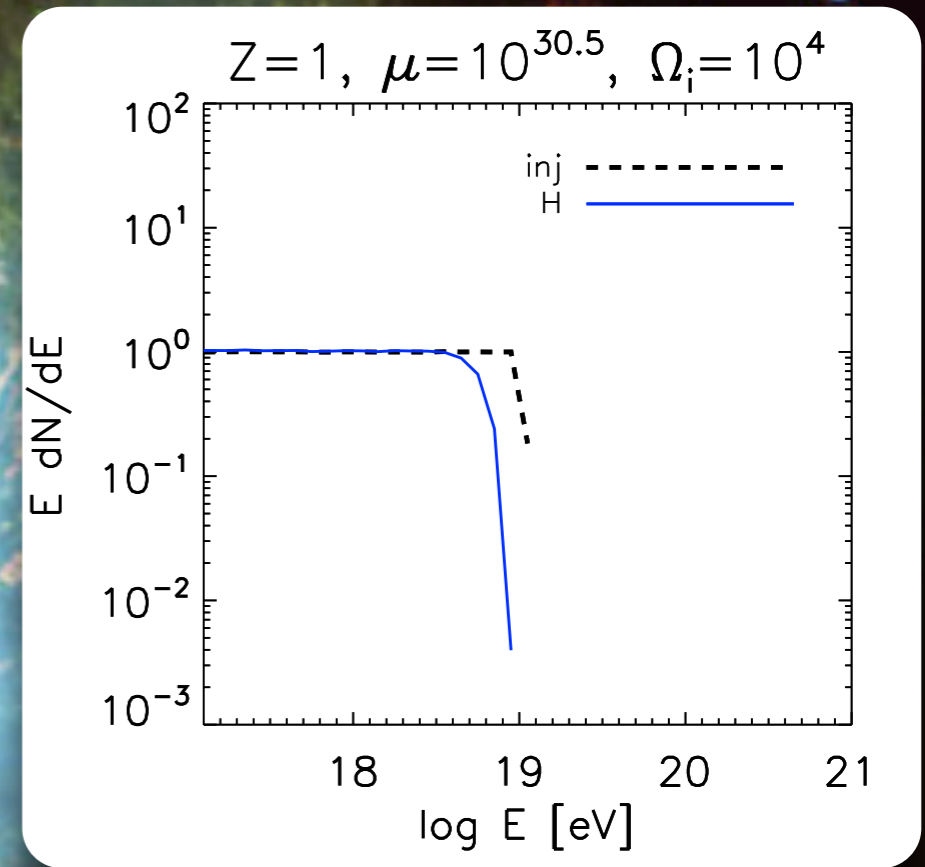




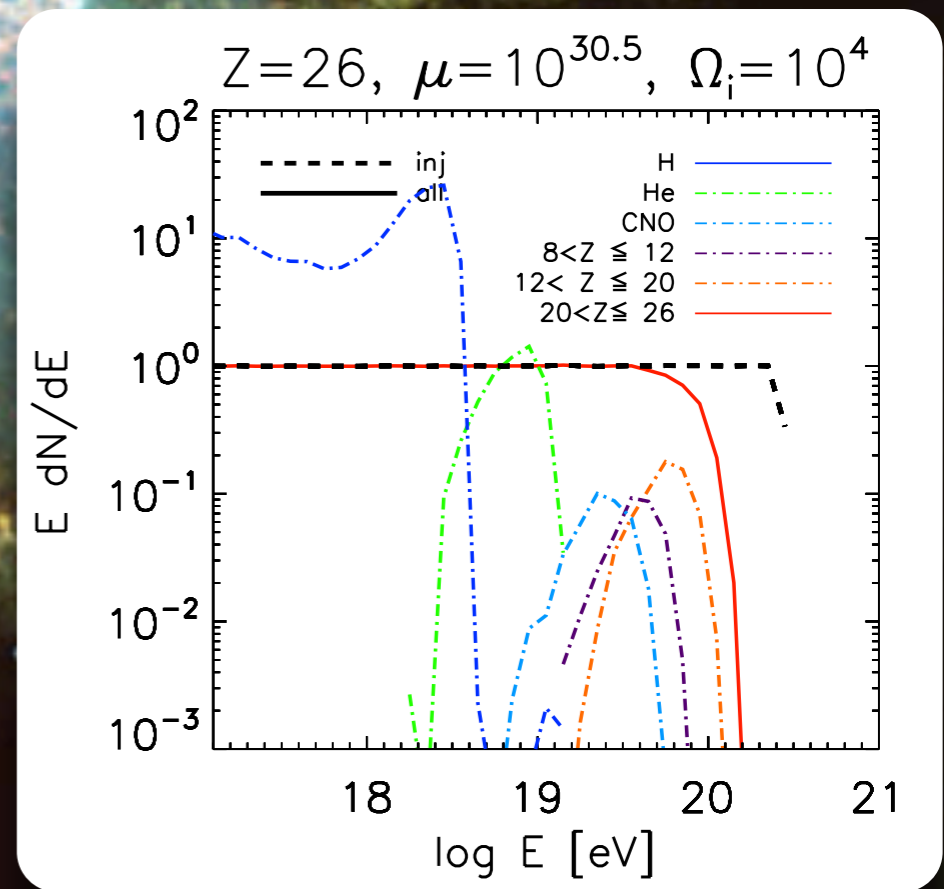
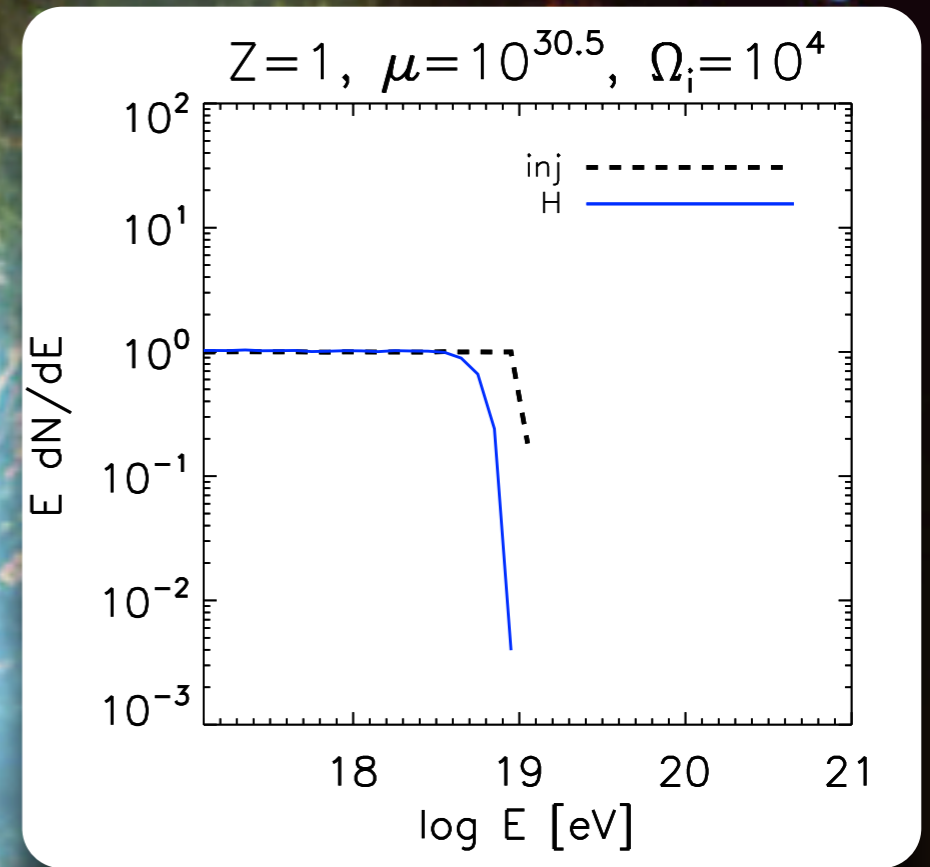
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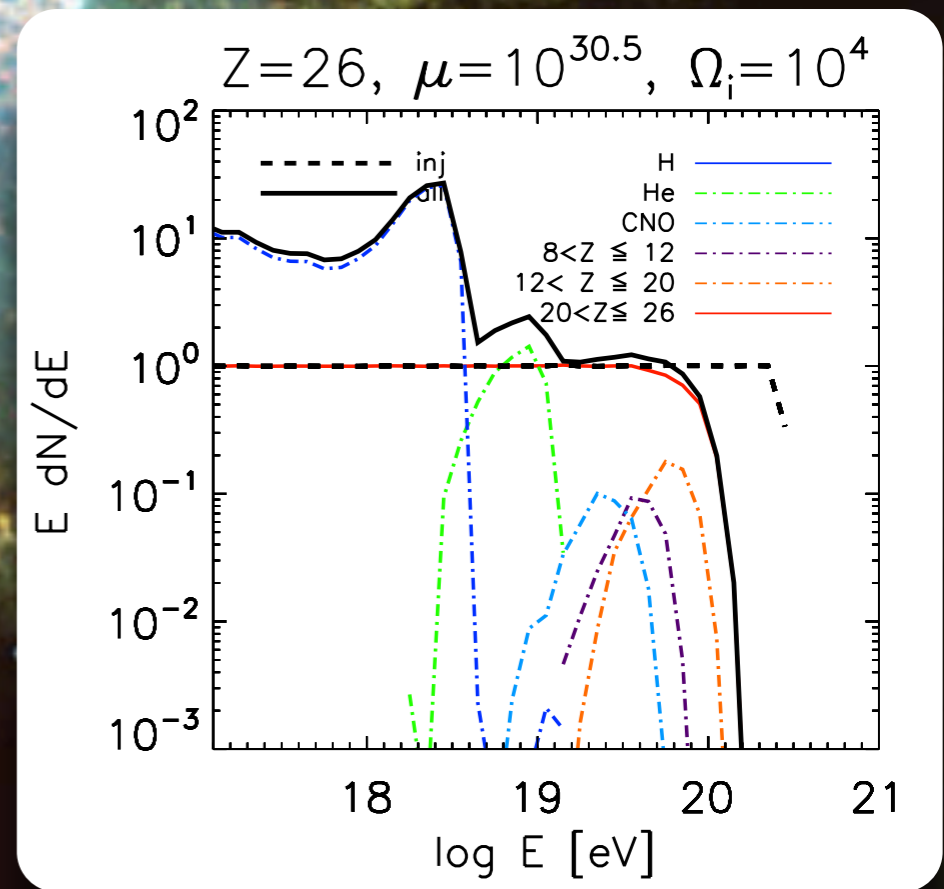
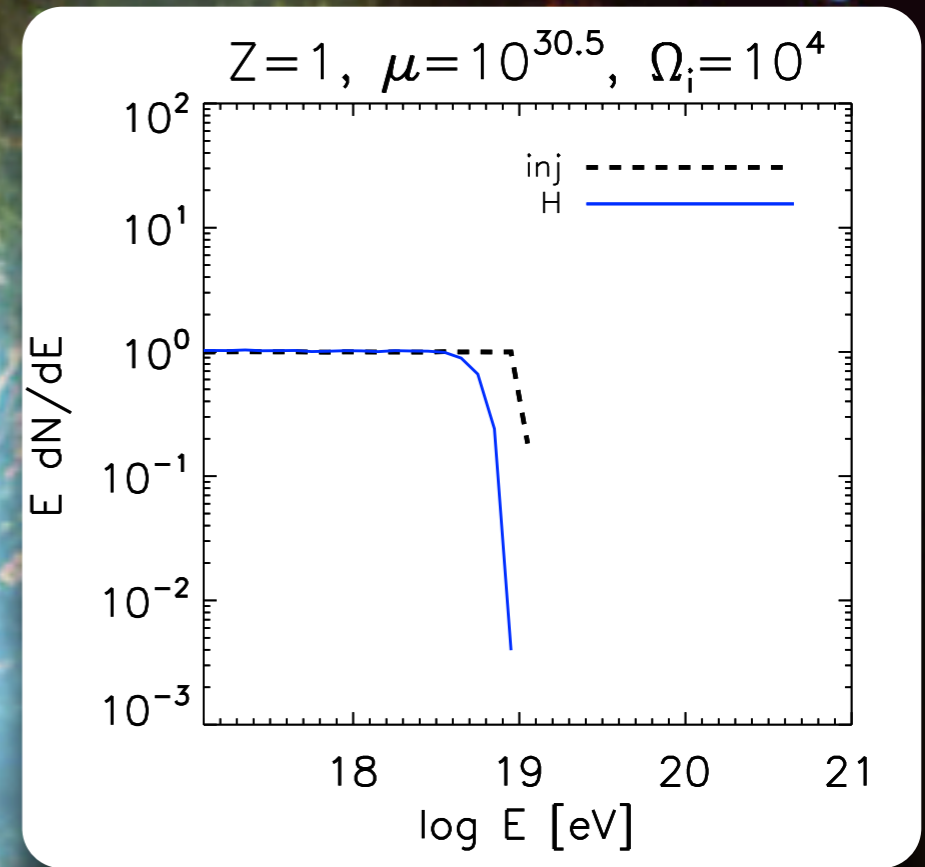
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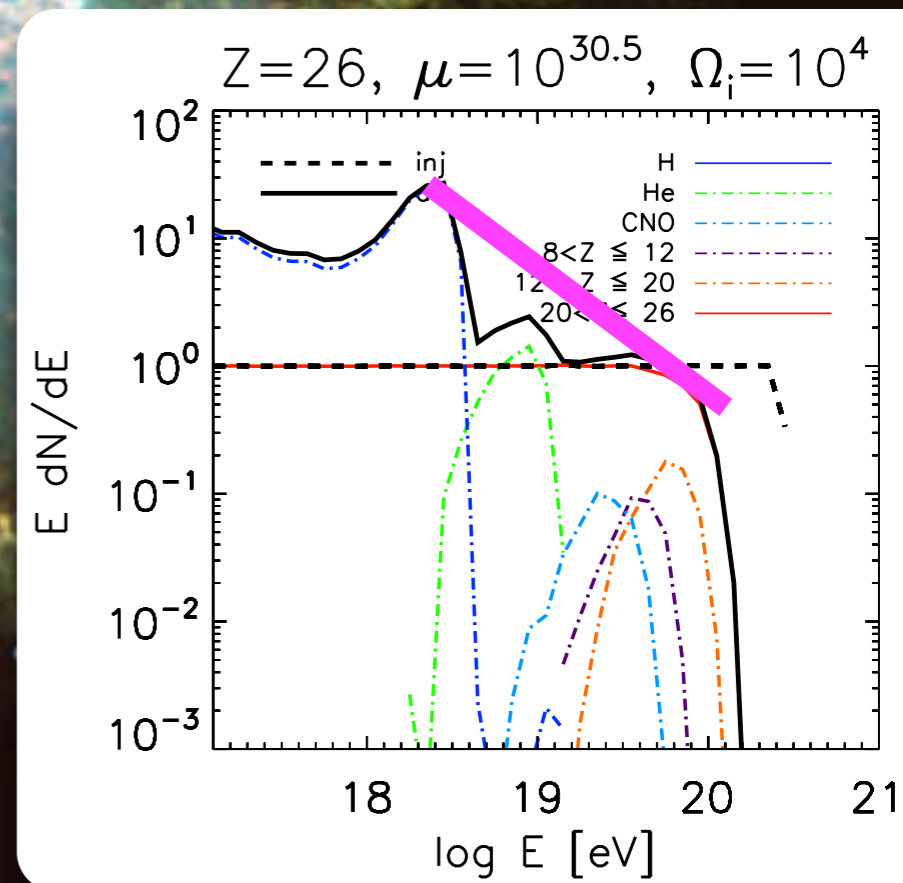
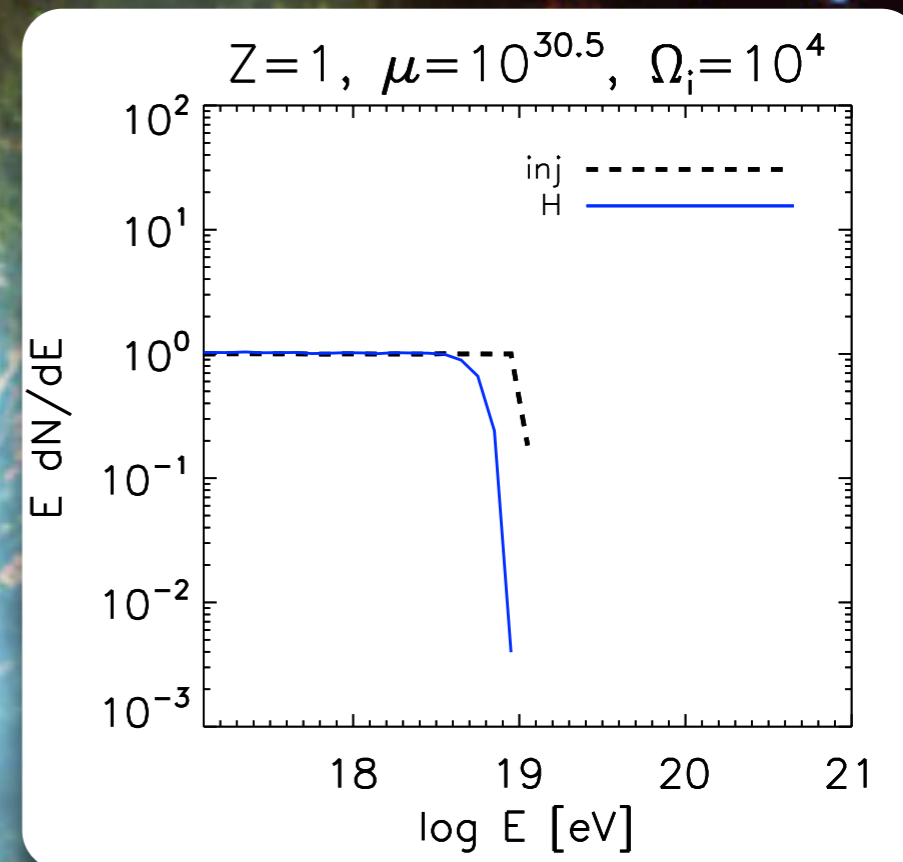
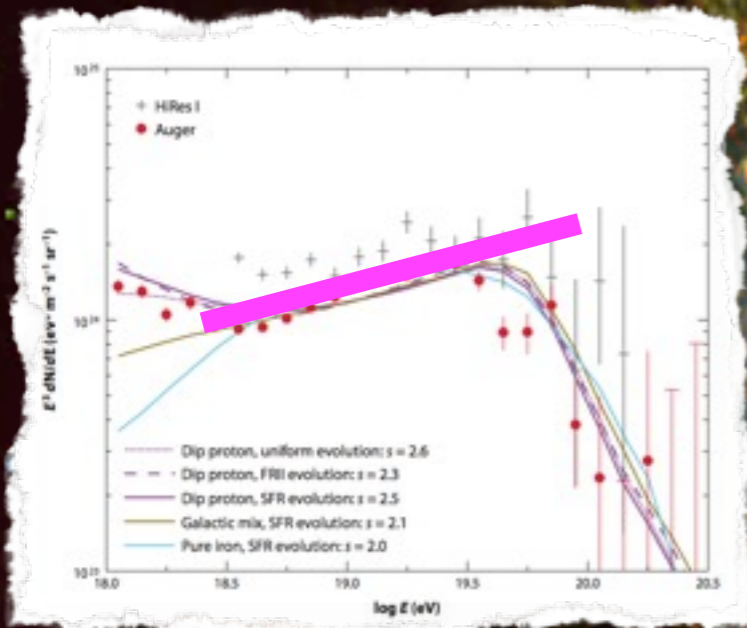
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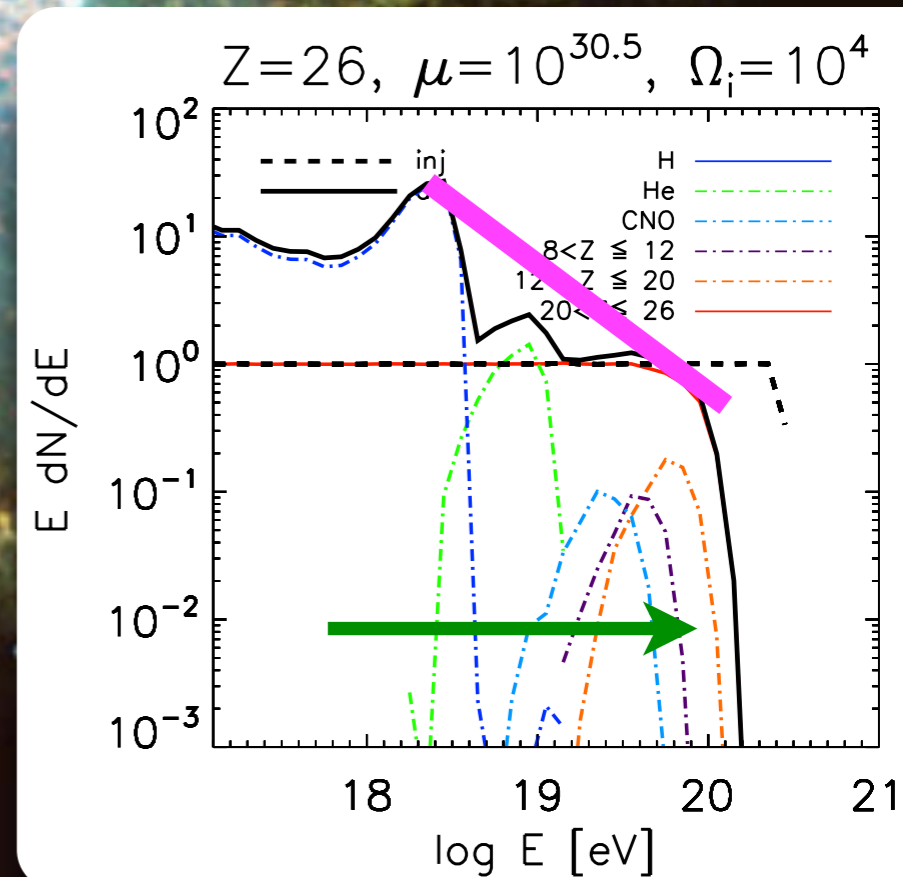
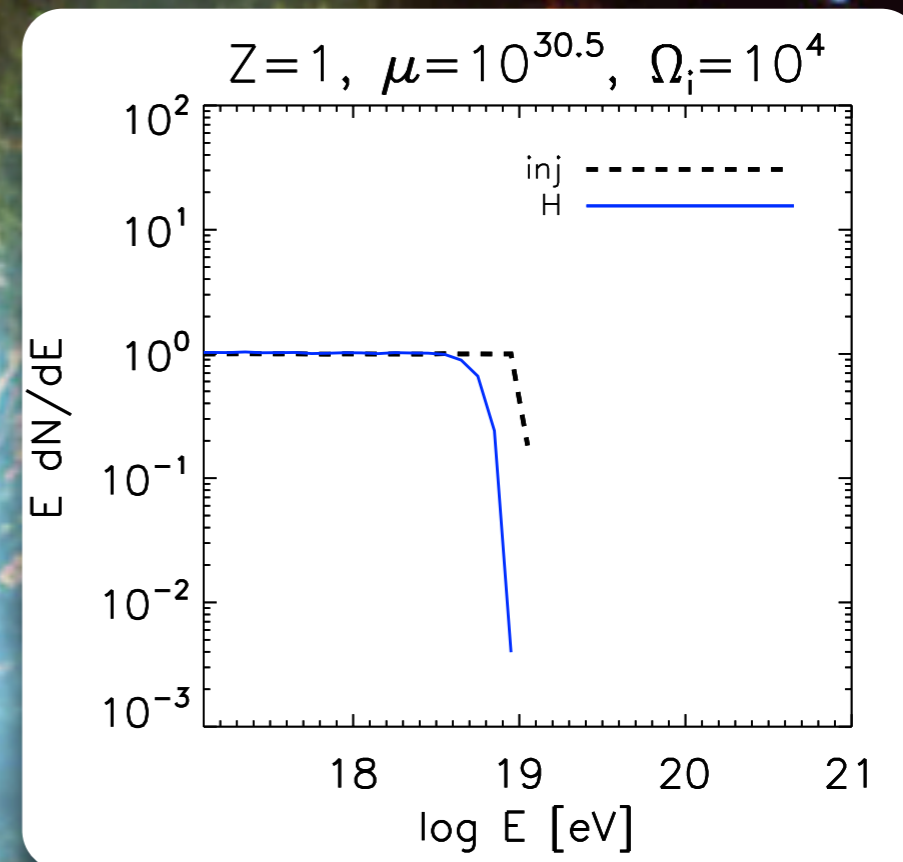
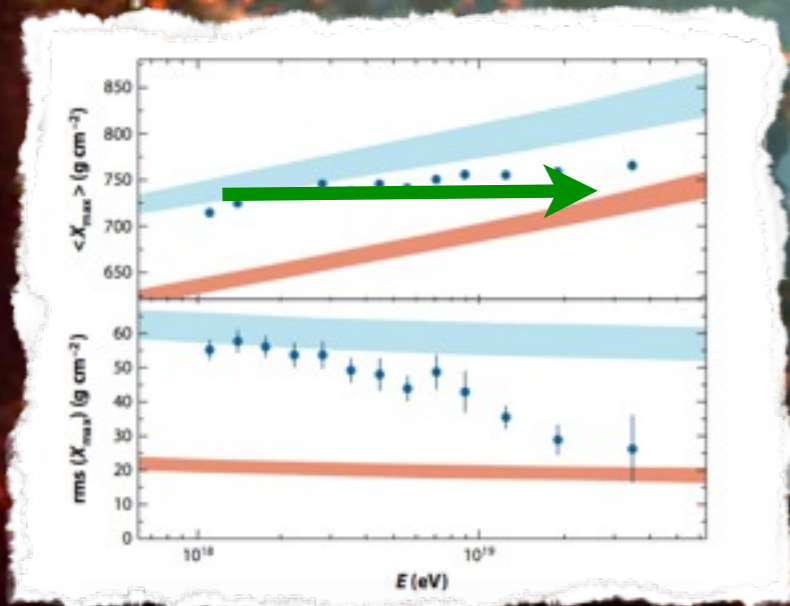
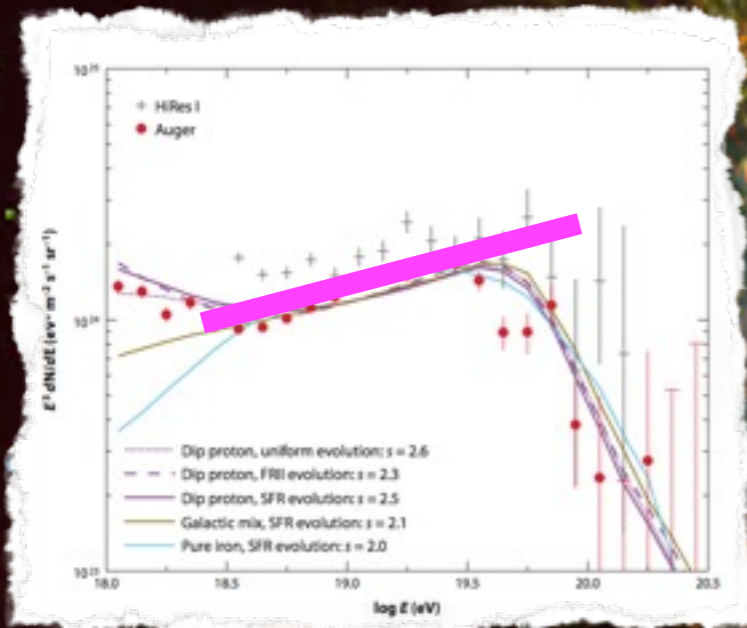
Monte-Carlo propagation  
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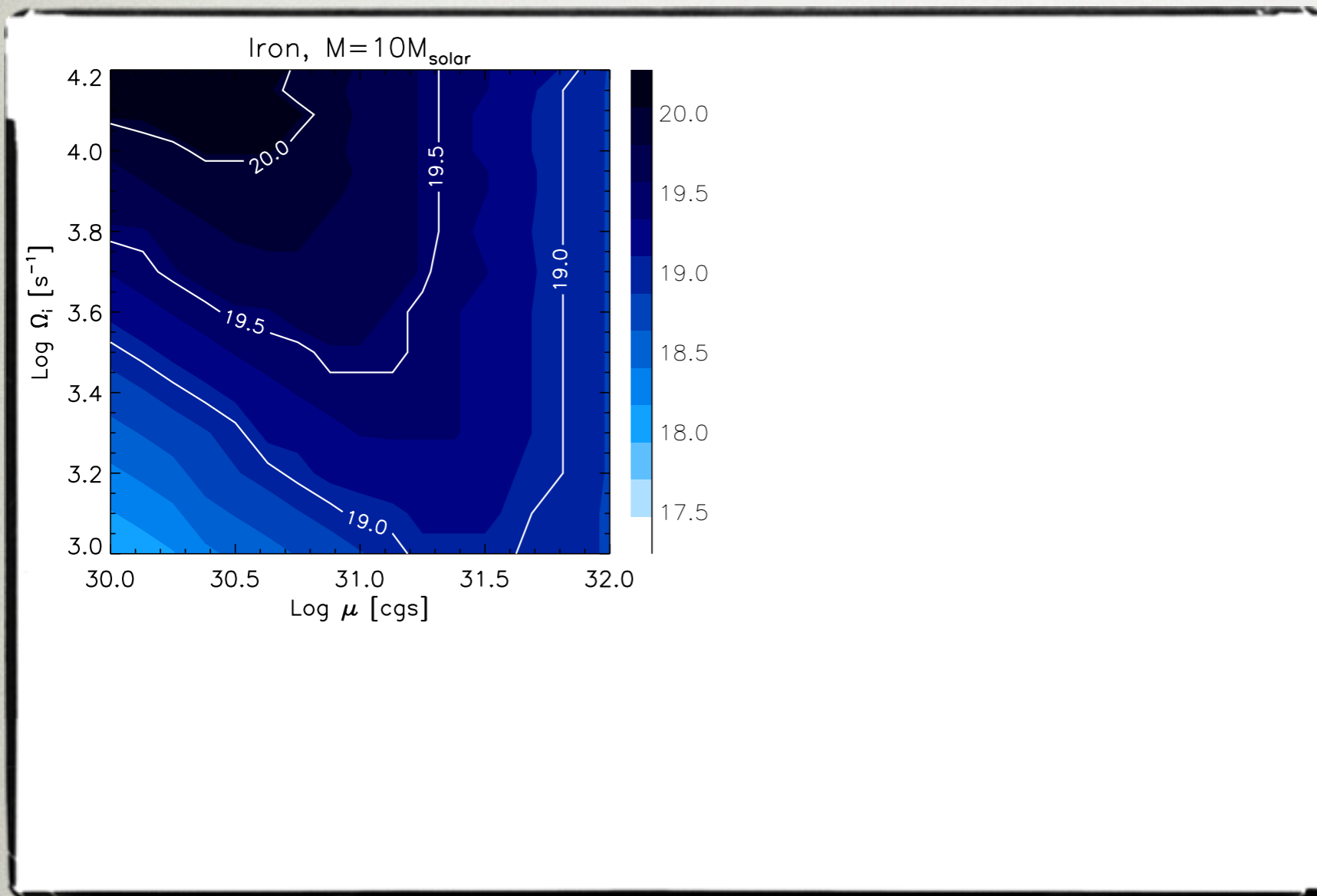
# Monte-Carlo propagation hadron interactions simulated with EPOS + CONEX



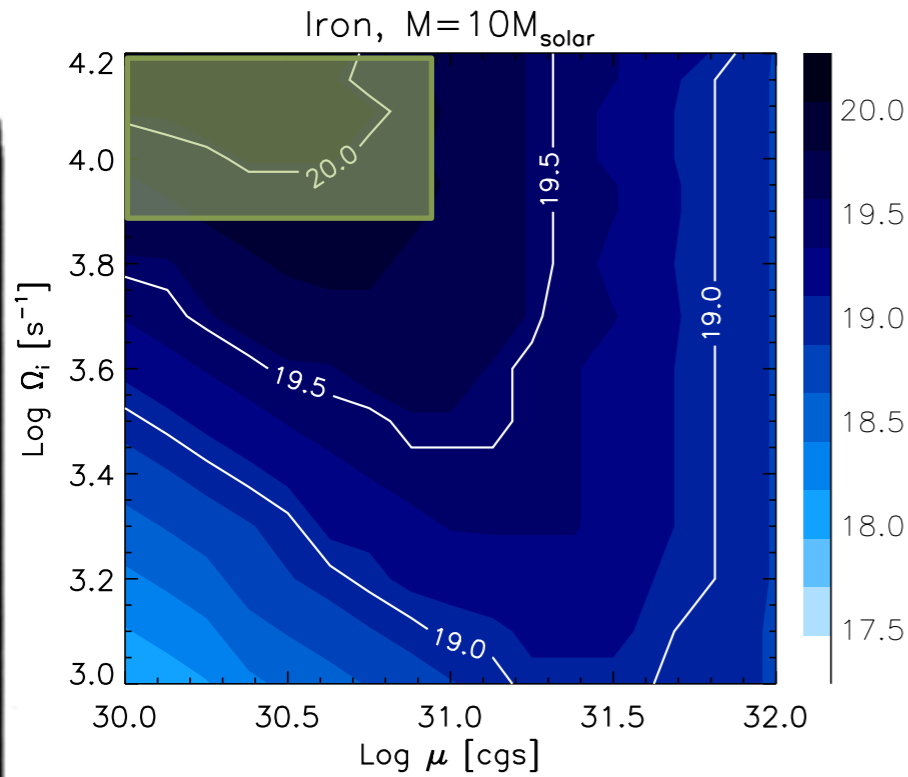
# Monte-Carlo propagation hadron interactions simulated with EPOS + CONEX



# Integrated Extragalactic Pulsars

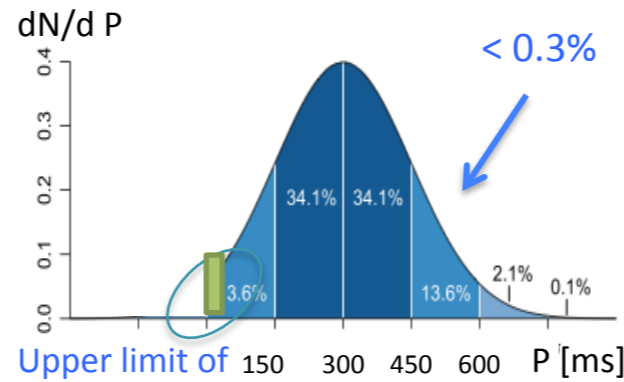
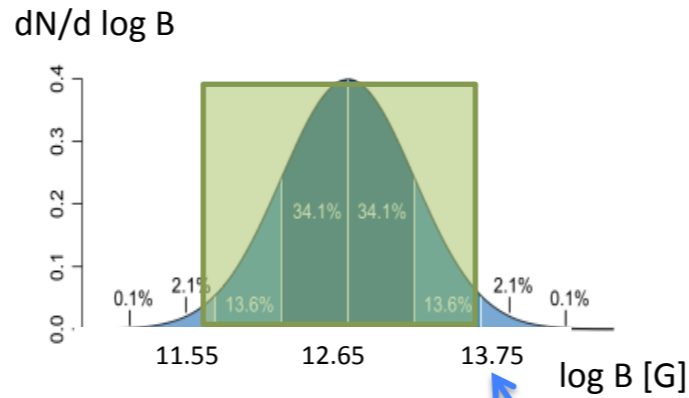


# Integrated Extragalactic Pulsars



## Pulsar distribution in the galaxy

*Faucher-Giguère & Kaspi 06*



Upper limit of rotational speed

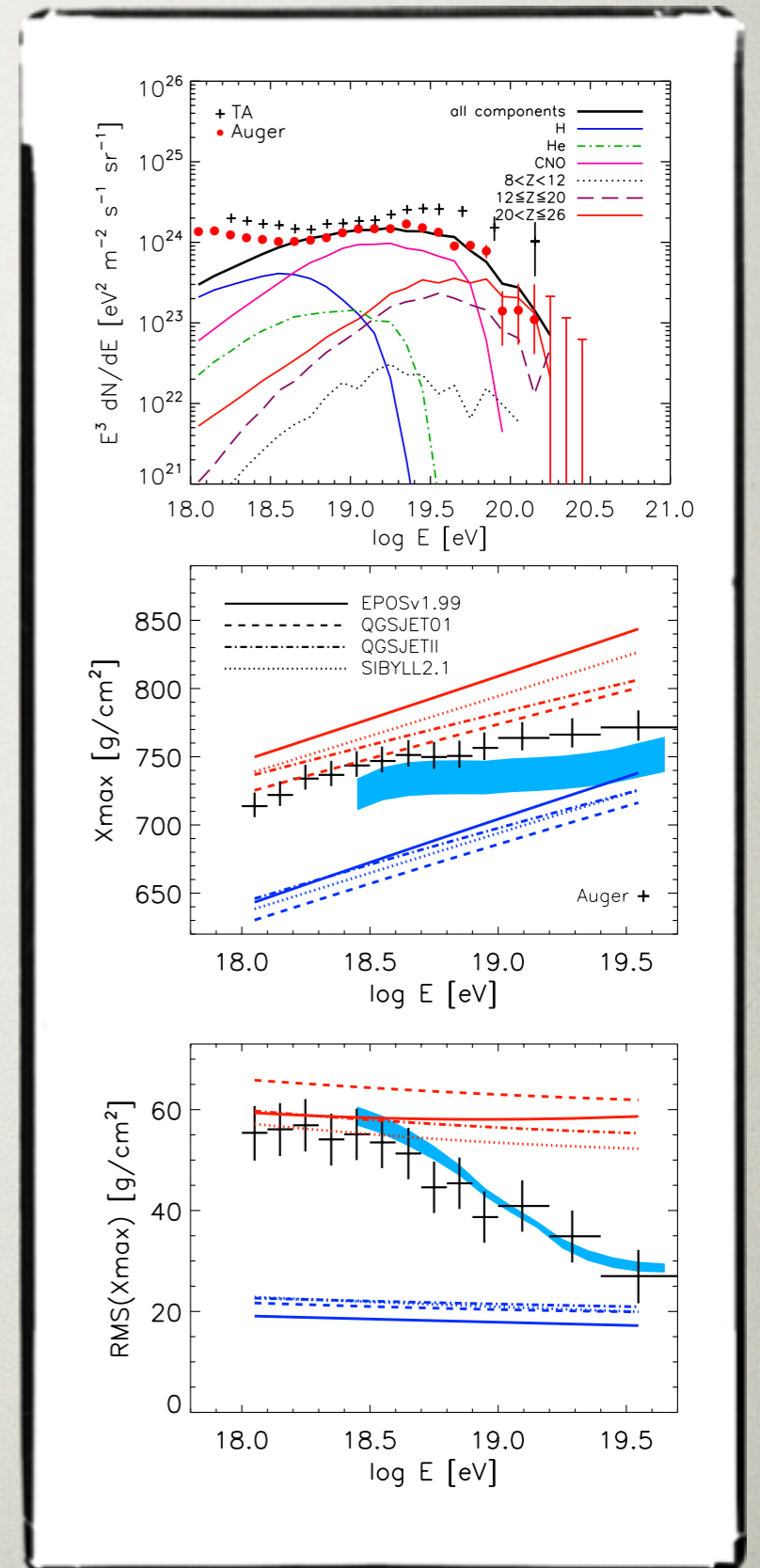
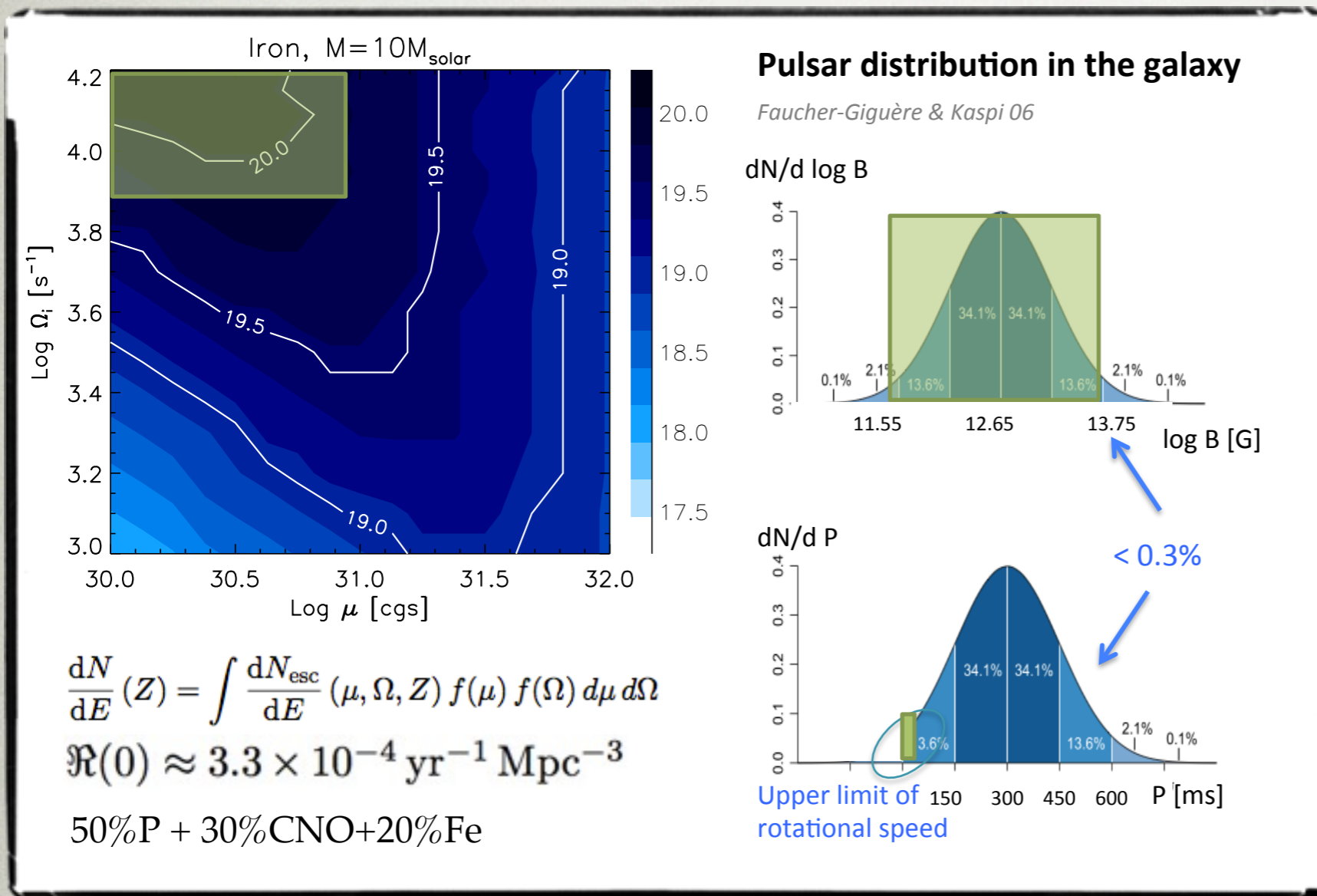
$$\frac{dN}{dE}(Z) = \int \frac{dN_{\text{esc}}}{dE}(\mu, \Omega, Z) f(\mu) f(\Omega) d\mu d\Omega$$

$$\mathcal{R}(0) \approx 3.3 \times 10^{-4} \text{ yr}^{-1} \text{ Mpc}^{-3}$$

50%P + 30%CNO+20%Fe



# Integrated Extragalactic Pulsars



## Conclusion I

Newborn pulsars can be successful  
UHE sources!

# Anisotropy Check

$$r_L = 10 \text{ Mpc} \frac{1}{Z} \frac{E}{10^{20} \text{ eV}} \left( \frac{B}{10^{-8} \text{ G}} \right)^{-1}$$

$\lambda \approx 10 - 100 \text{ kpc} \ll r_L \Rightarrow$  small deflections

$$\delta\theta^2 \approx \frac{r_{\text{structure}}}{r_L^2 / l_c}$$

$$\delta\theta_i \simeq 1.7^\circ \left( \frac{\bar{r}_i}{2 \text{ Mpc}} \right)^{1/2} \left( \frac{B_i}{10^{-8} \text{ G}} \right) \times \left( \frac{\lambda_i}{0.1 \text{ Mpc}} \right)^{1/2} \left( \frac{E}{10^{20} \text{ eV}} \right)^{-1}$$

Kotera et al 2009

Time delay after the deflections

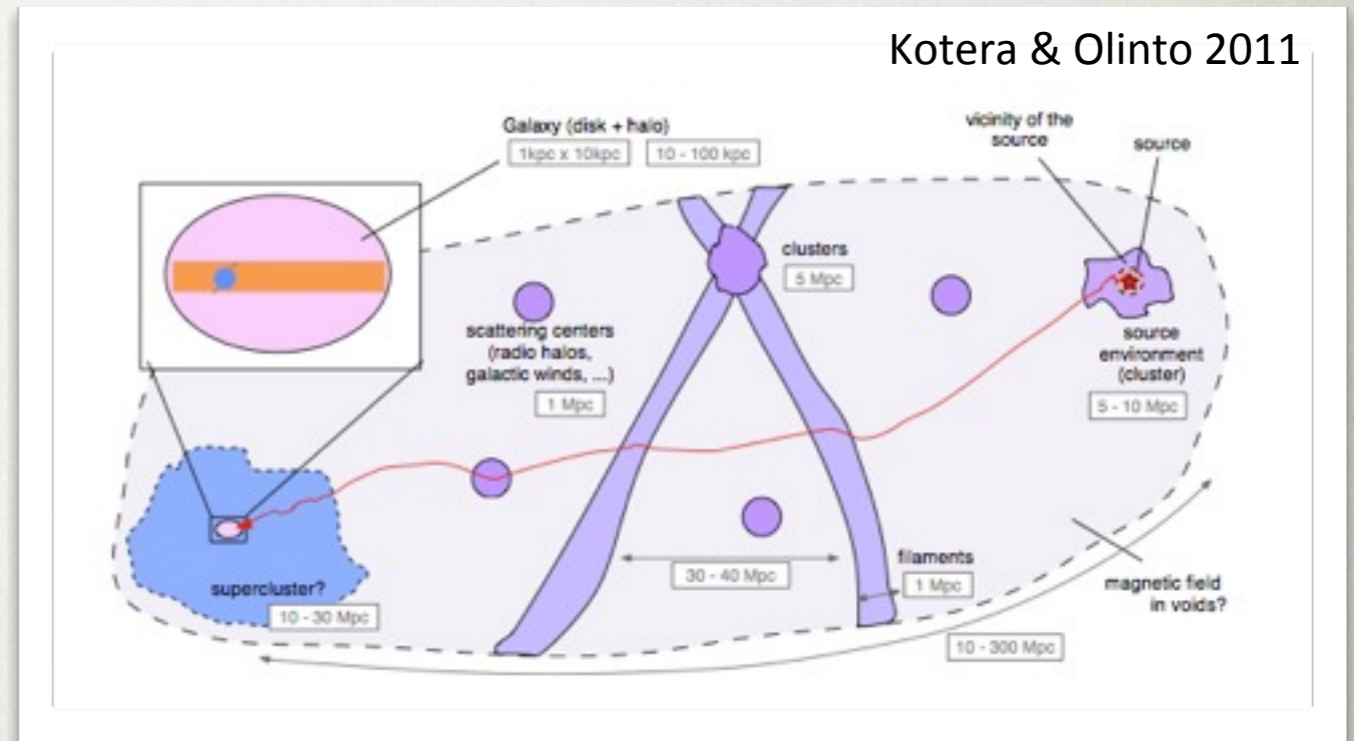
$$\delta t_i \simeq 0.93 \times 10^3 \text{ yr} \left( \frac{\bar{r}_i}{2 \text{ Mpc}} \right)^2 \left( \frac{B_i}{10^{-8} \text{ G}} \right)^2 \times \left( \frac{\lambda_i}{0.1 \text{ Mpc}} \right) \left( \frac{E}{10^{20} \text{ eV}} \right)^{-2}$$

Kotera et al 2009

$\gg$

$$t_{\text{spin}} = 3 \text{ yr} \left( \frac{10^{20} \text{ eV}}{E} \right) \frac{Z_{26} \eta_1}{\mu_{30.5}}$$

$\Rightarrow$



Time the source was lighted

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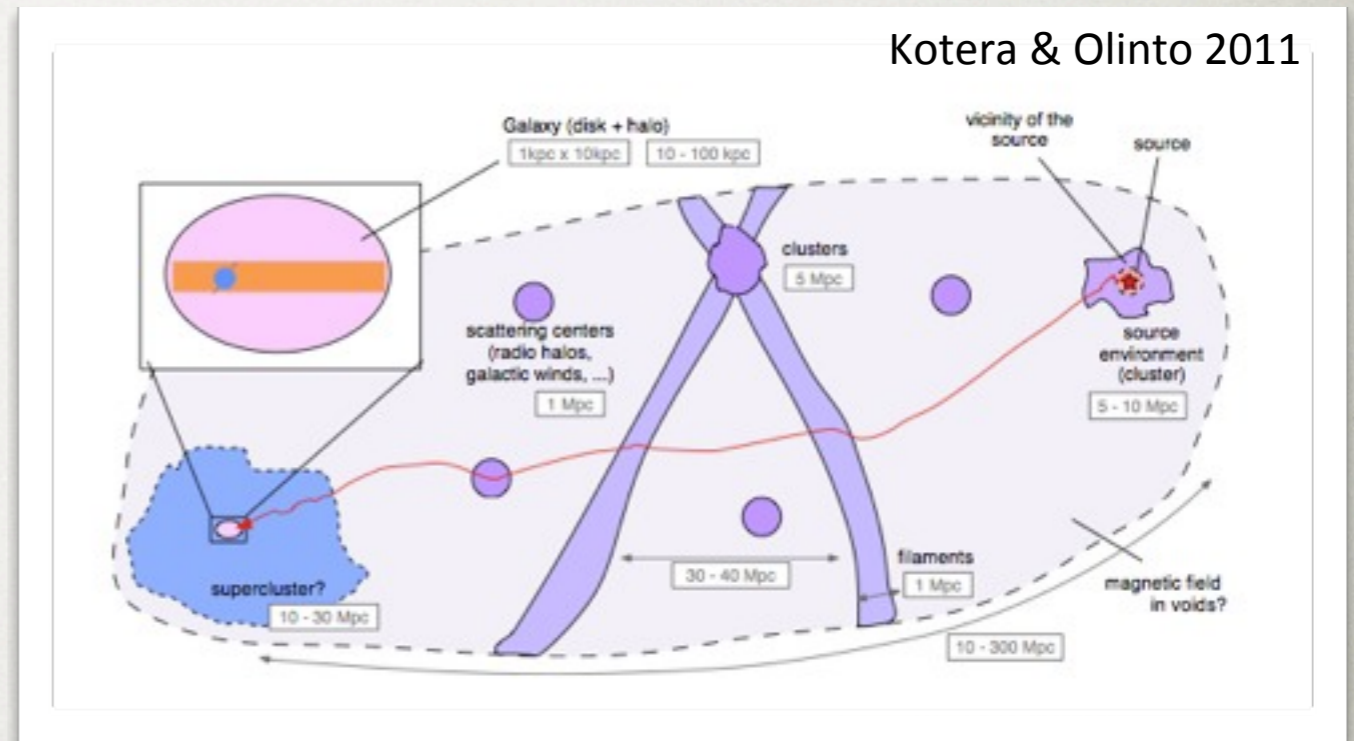
Kotera et al 2009

$\gg$

$$t_{\text{spin}} = 3 \text{ yr} \left( \frac{10^{20} \text{ eV}}{E} \right) \frac{Z_{26} \eta_1}{\mu_{30.5}}$$

$\Rightarrow$

Transients, no source-arrival direction correlation

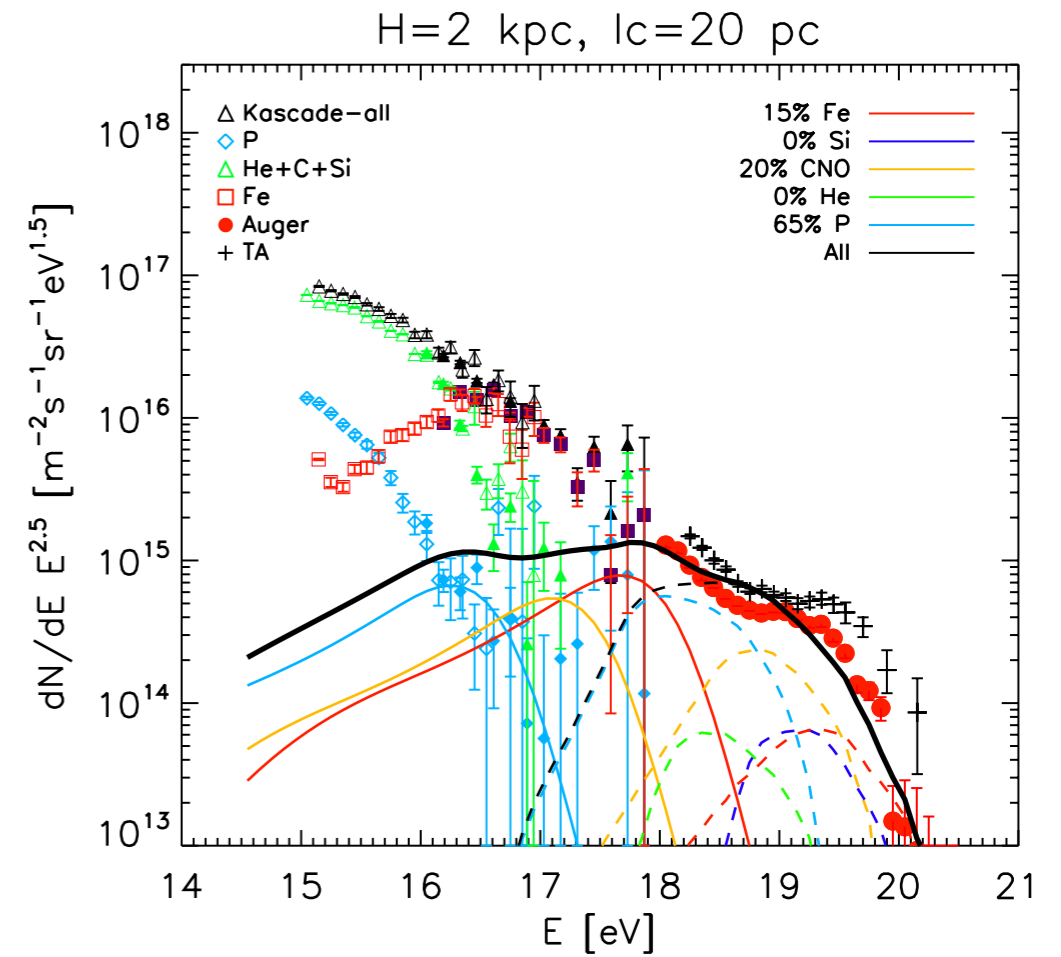
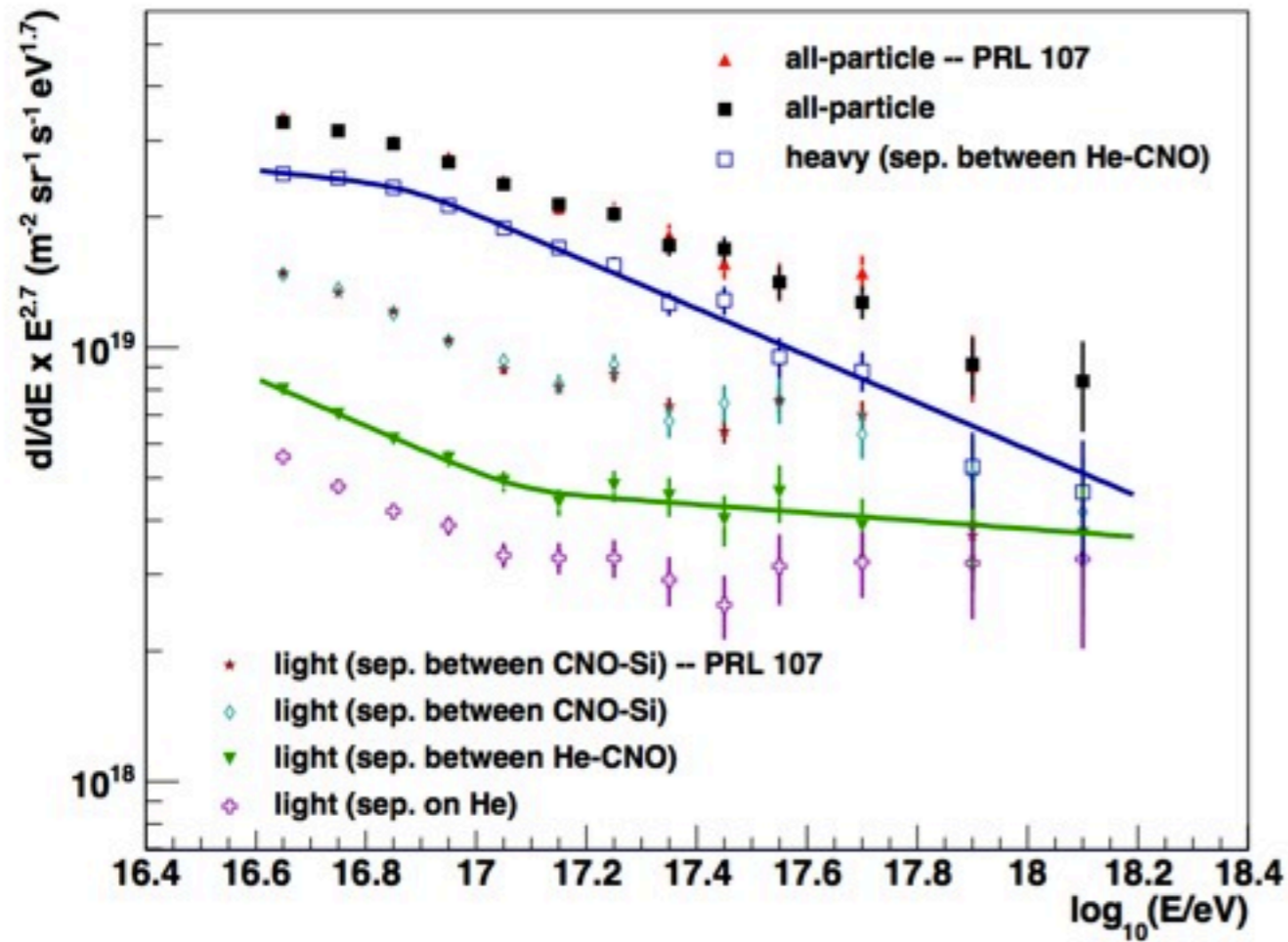


Time the source was lighted

**What about their Galactic Counterparts?**

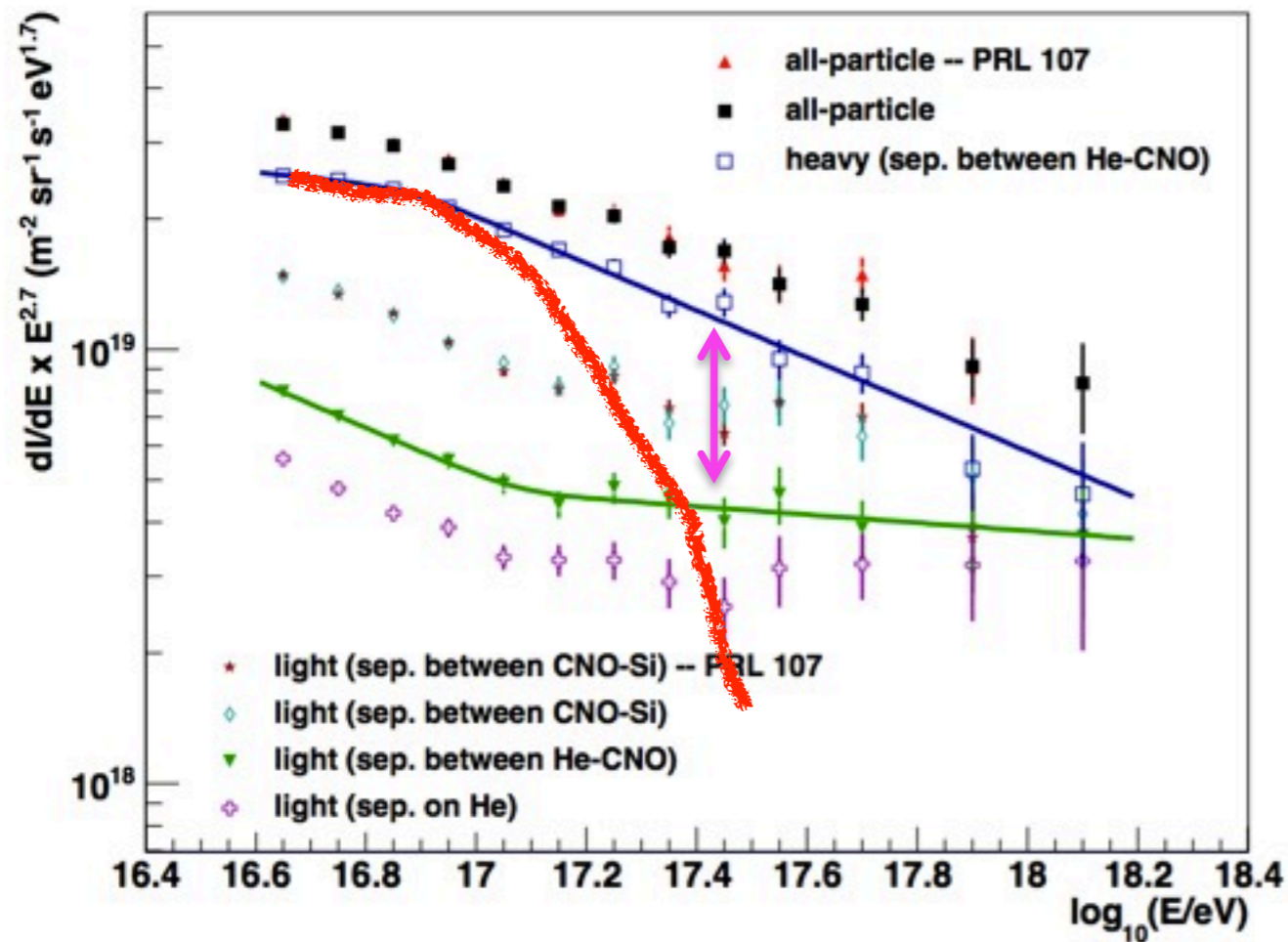
# Contribution from Galactic pulsars

Apel et al 2013



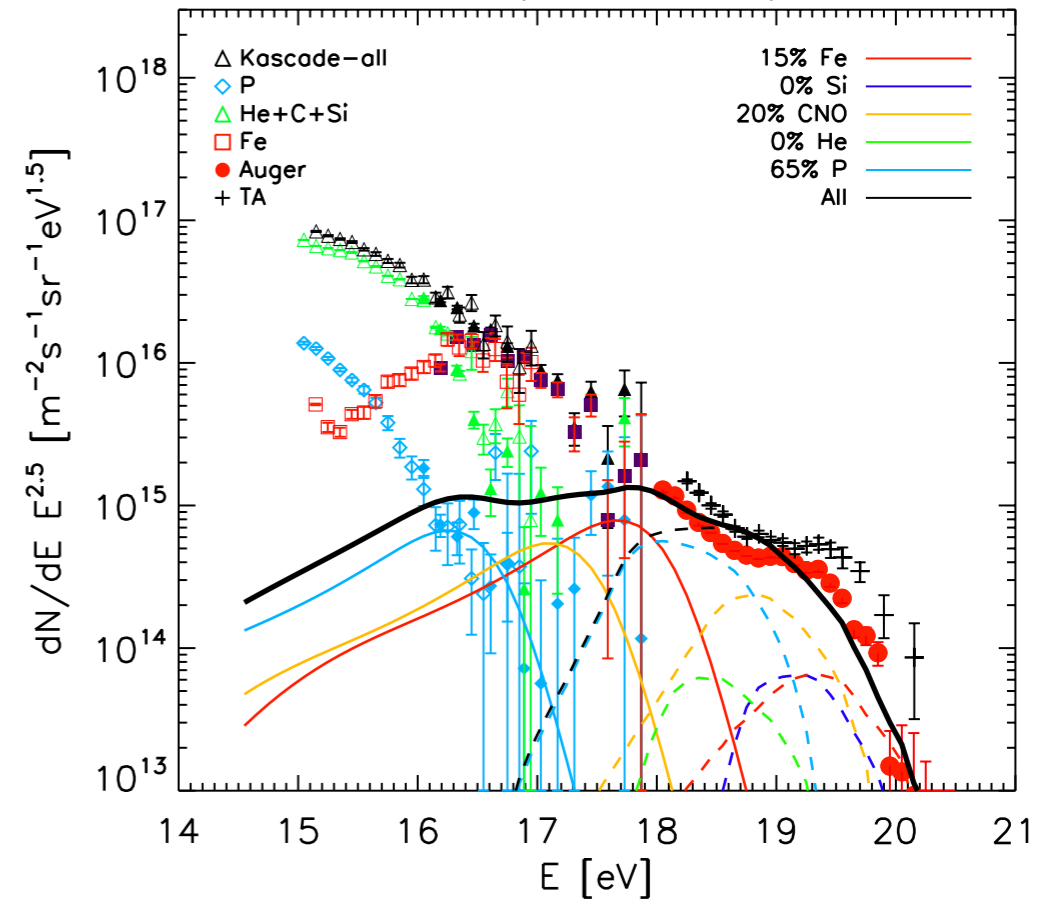
# Contribution from Galactic pulsars

Apel et al 2013



No Cutoff, Mind the Gap!

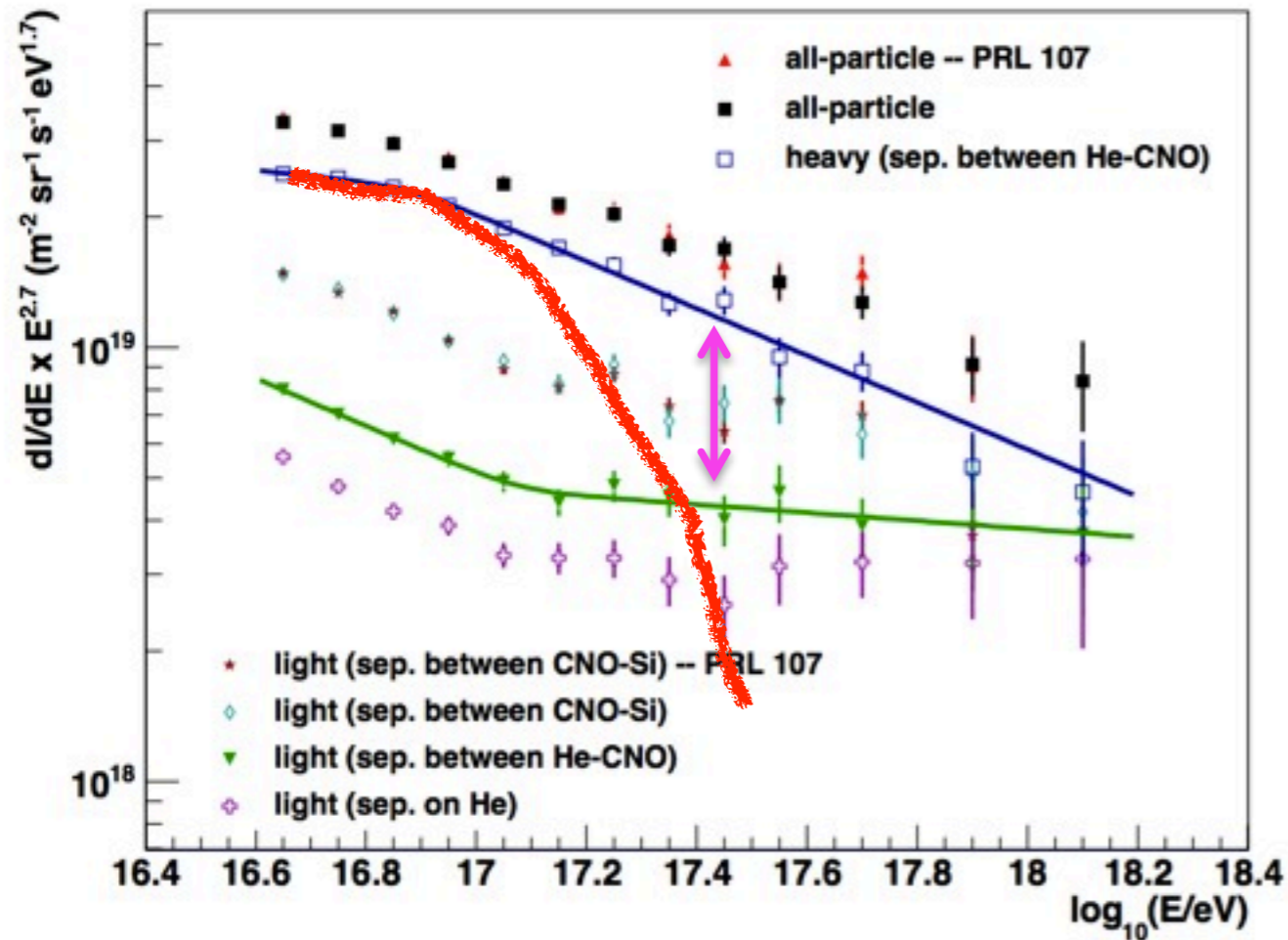
H=2 kpc, lc=20 pc



Galactic pulsars can fill the gap

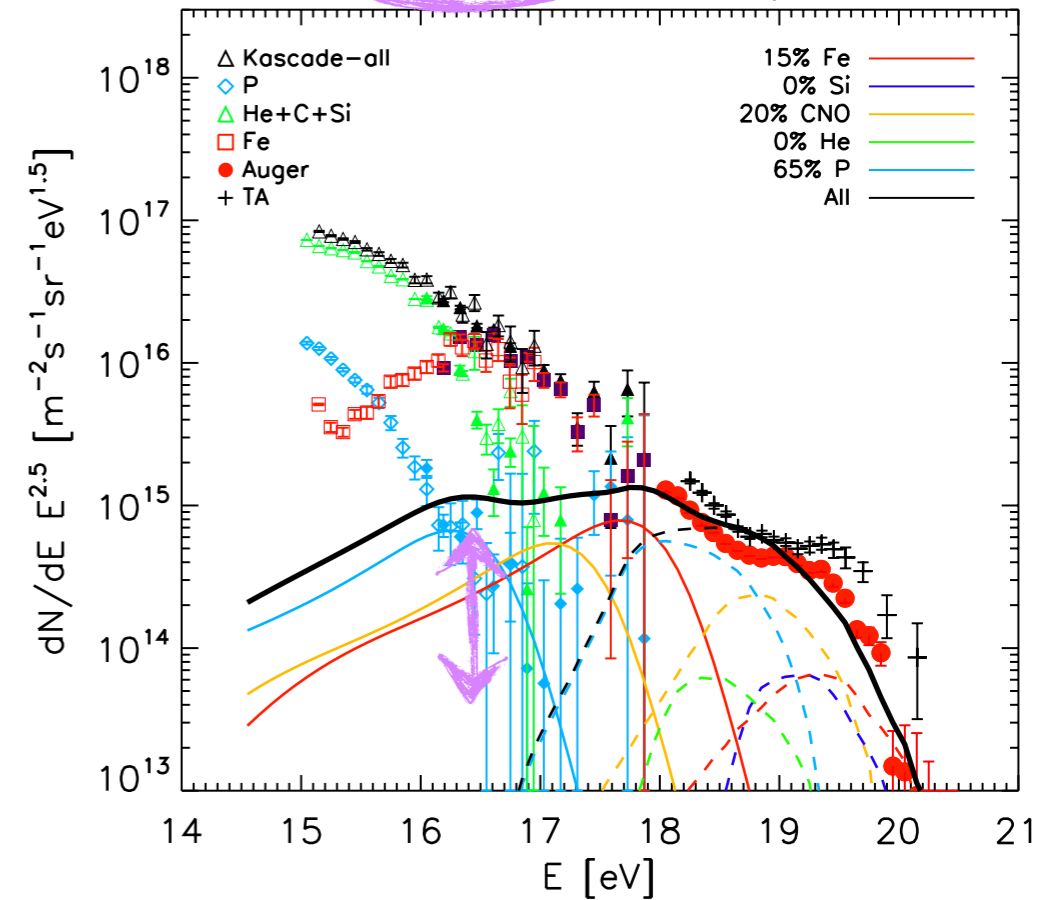
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Apel et al 2013



No Cutoff, Mind the Gap!

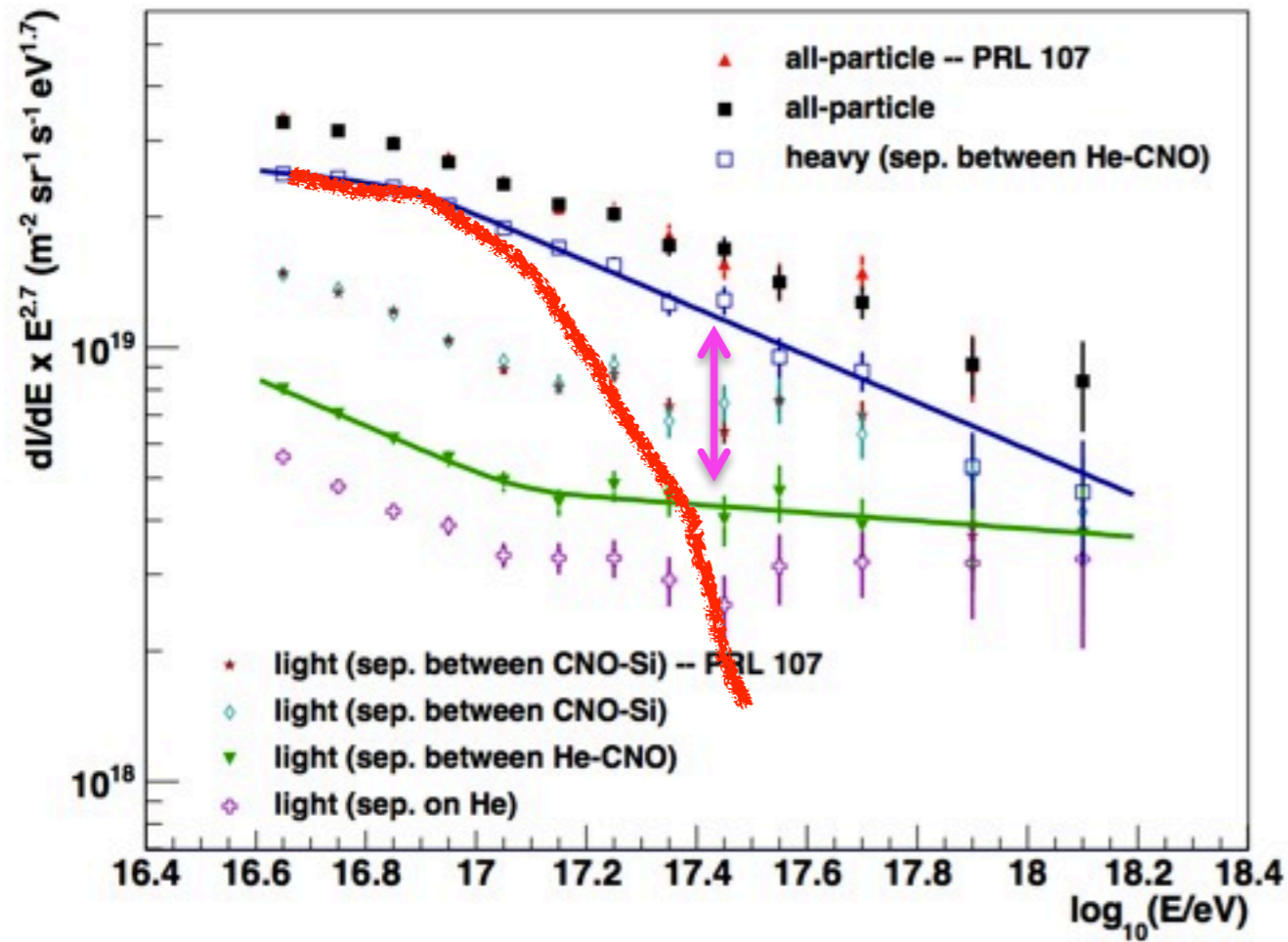
$H=2 \text{ kpc}$ ,  $l_c=20 \text{ pc}$



Galactic pulsars can fill the gap

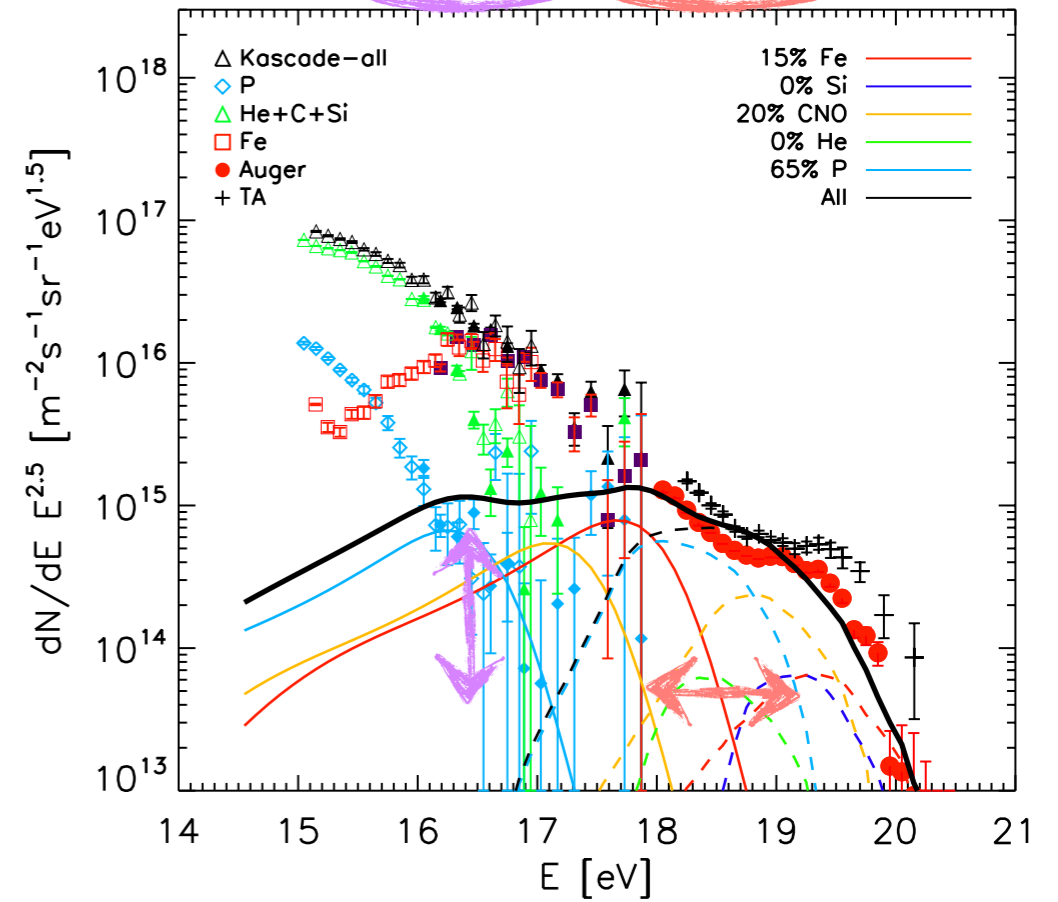
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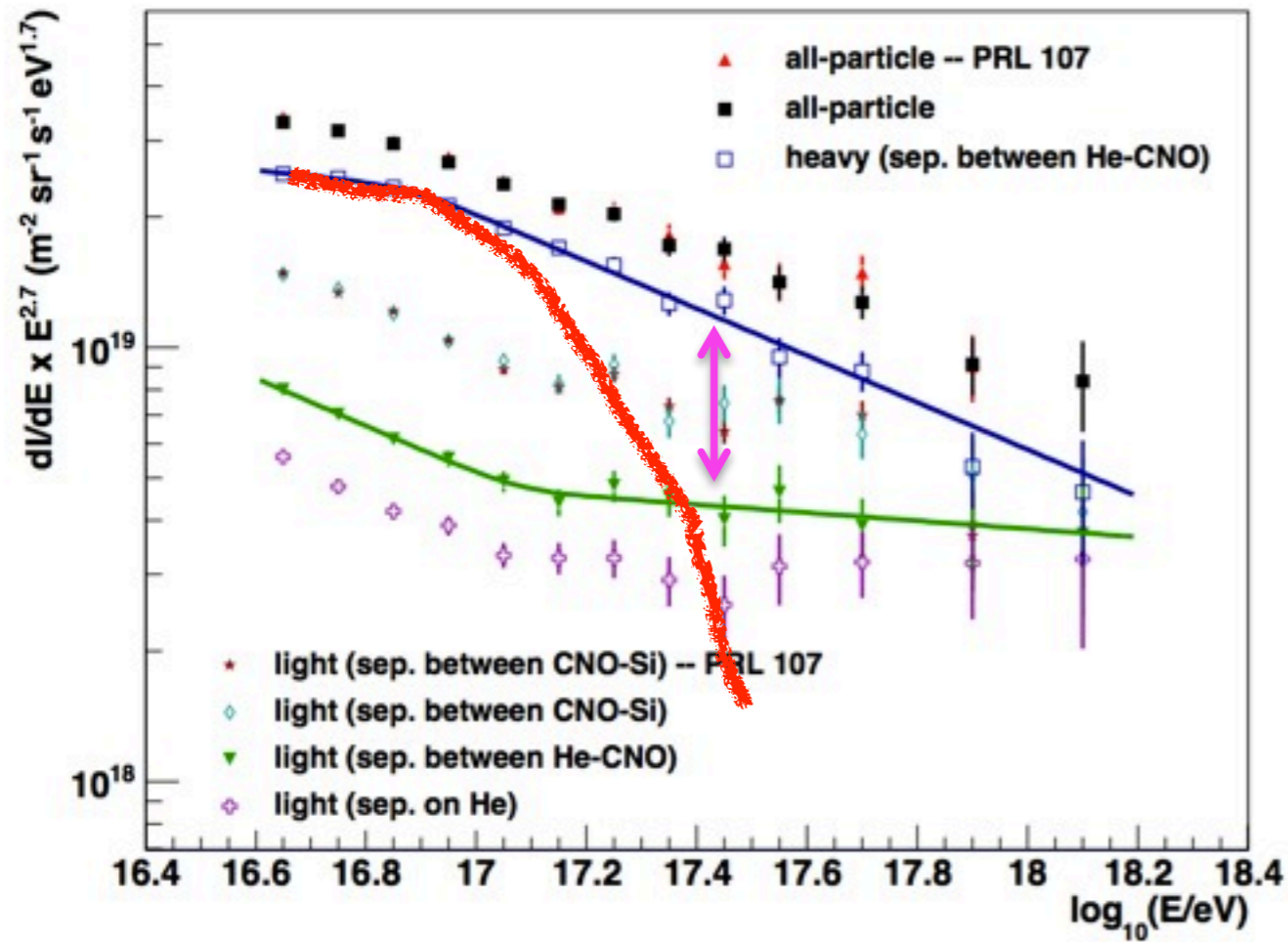


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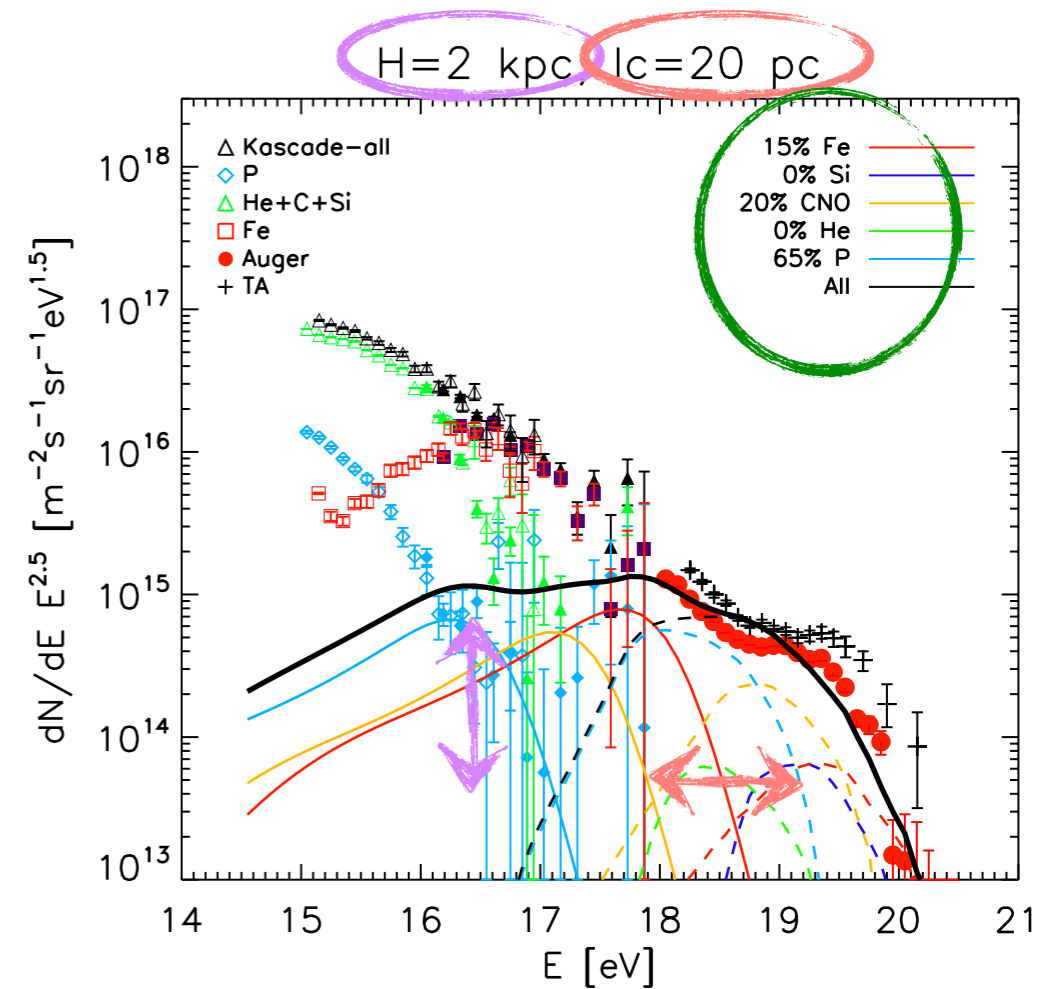


# Contribution from Galactic pulsars

Apel et al 2013

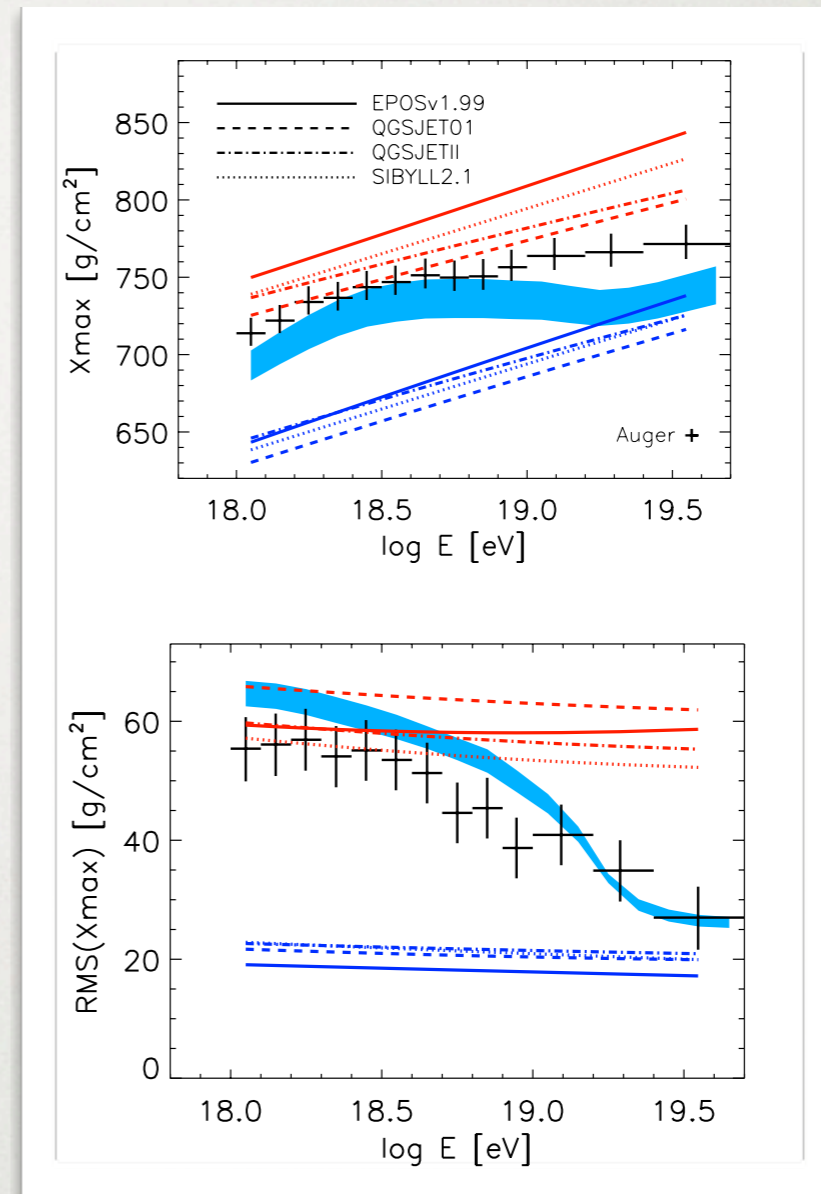
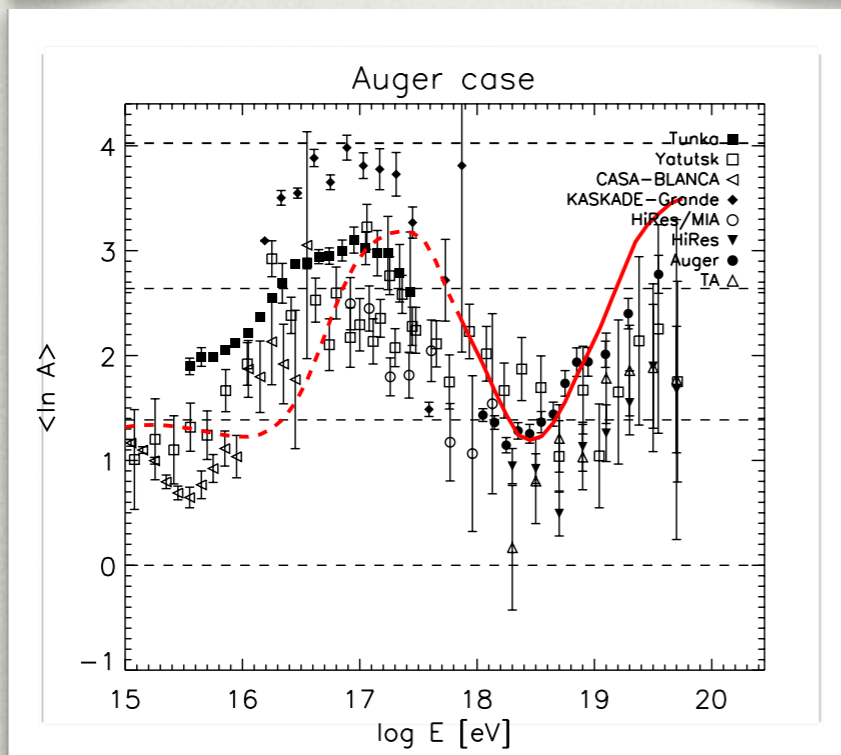
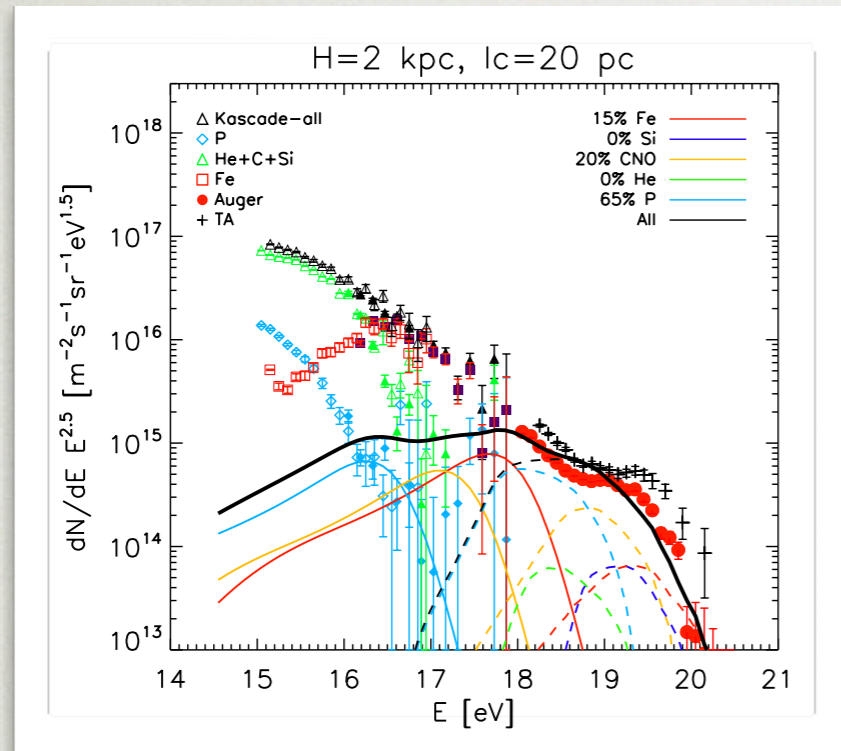


No Cutoff, Mind the Gap!

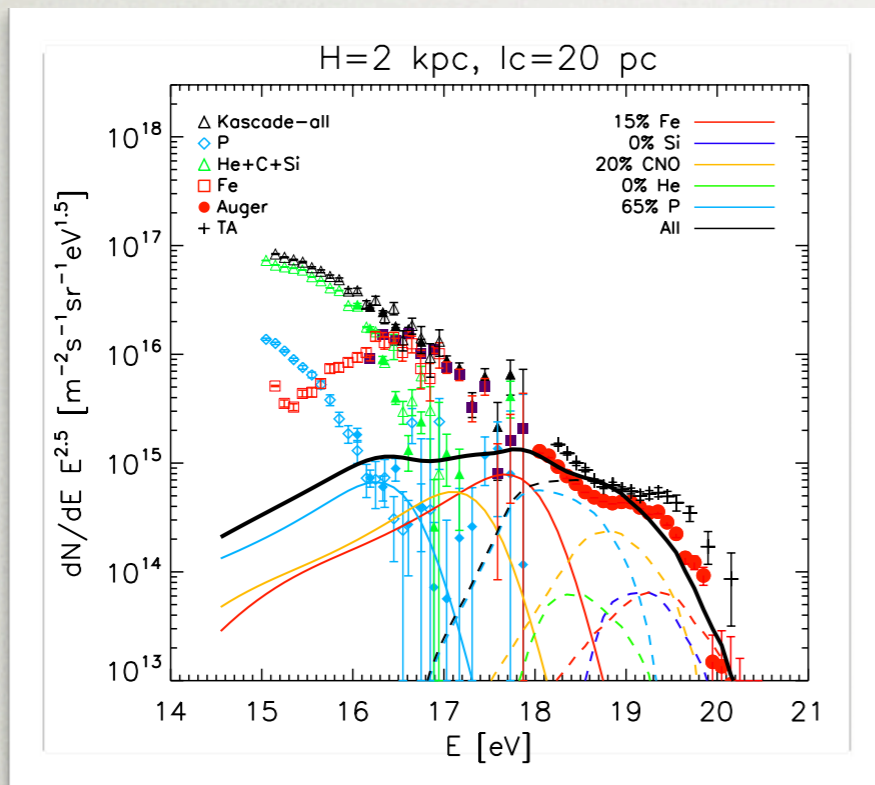


Galactic pulsars can fill the gap

# Composition

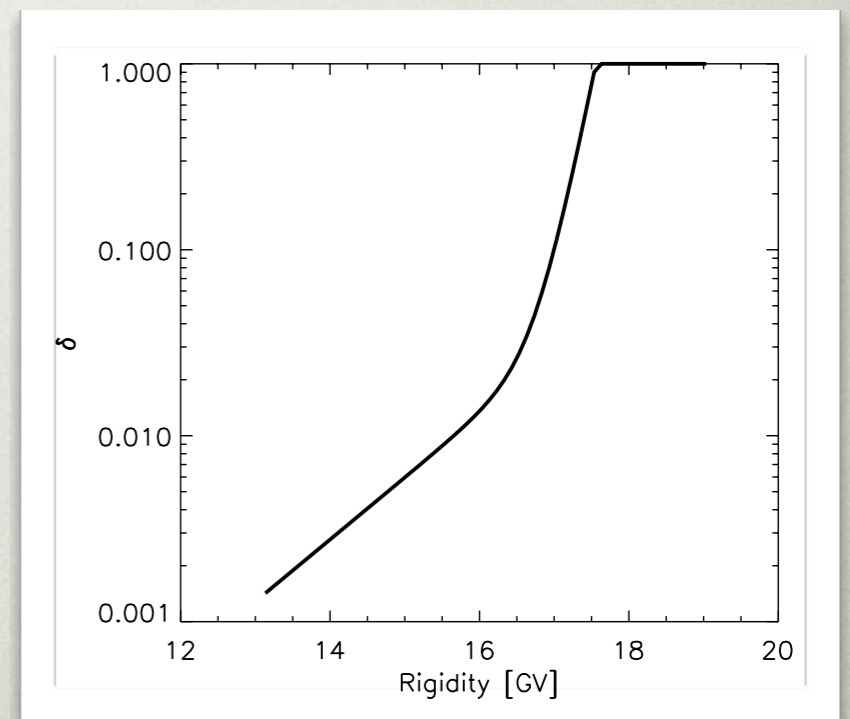


# Estimation on Anisotropy

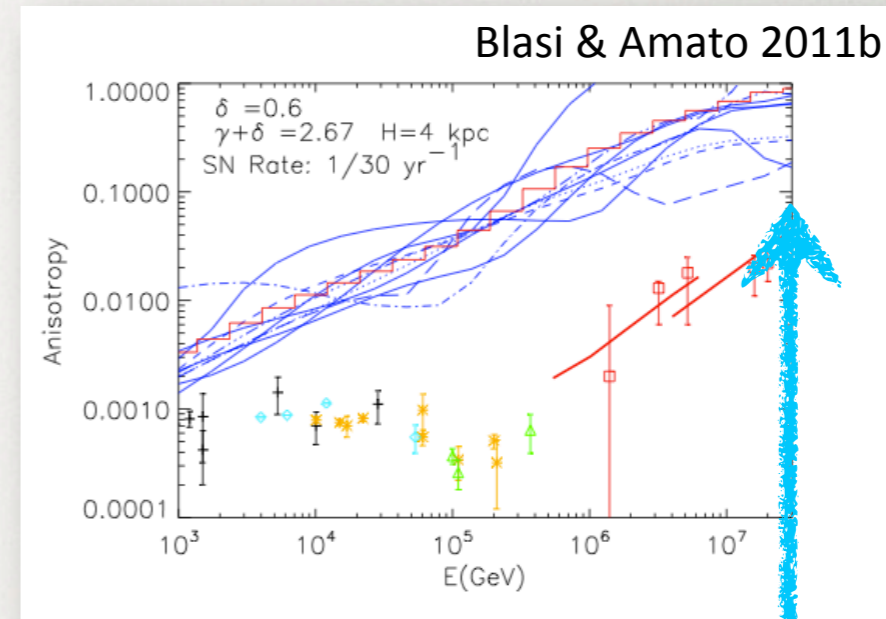
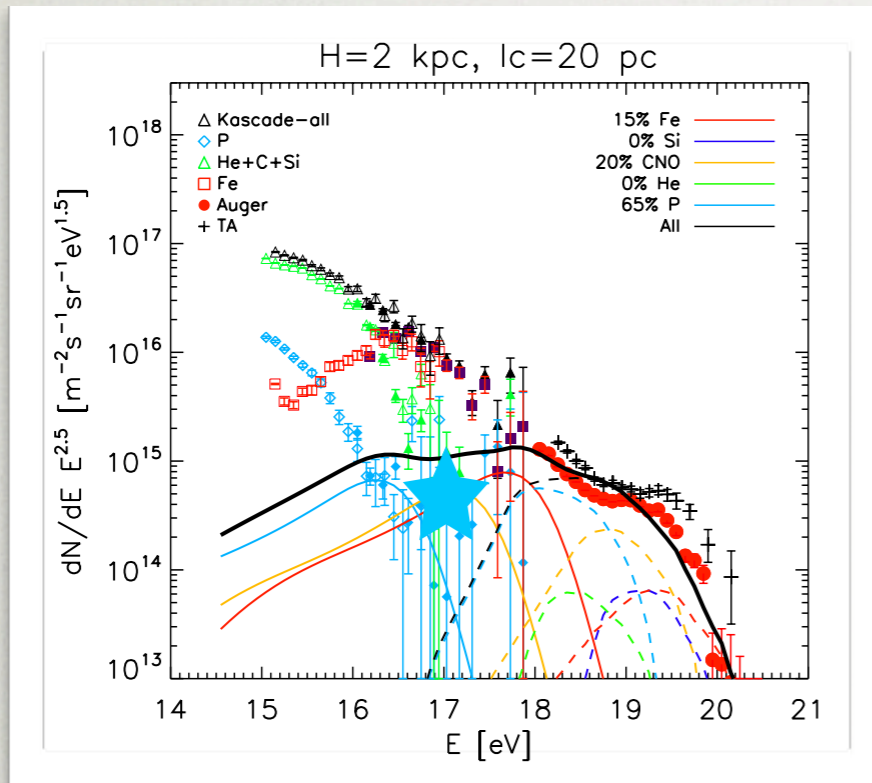


Assume sources homogeneously distributed in the disc, small scale anisotropy can be estimated as (Blasi & Amato 2011b)

$$\delta = \frac{3}{2^{3/2} \pi^{1/2}} \frac{D(E)}{Hc}$$

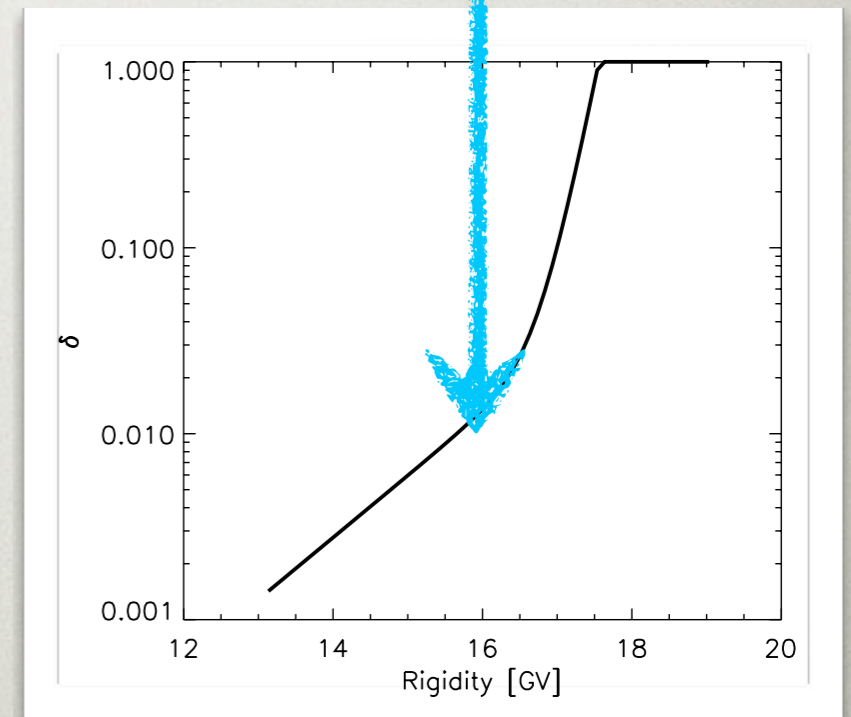


# Estimation on Anisotropy

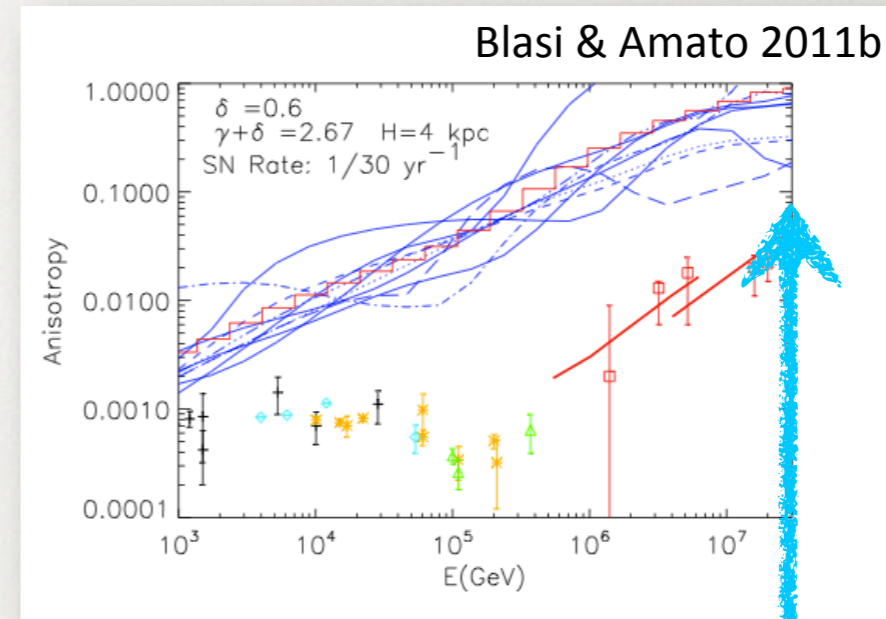
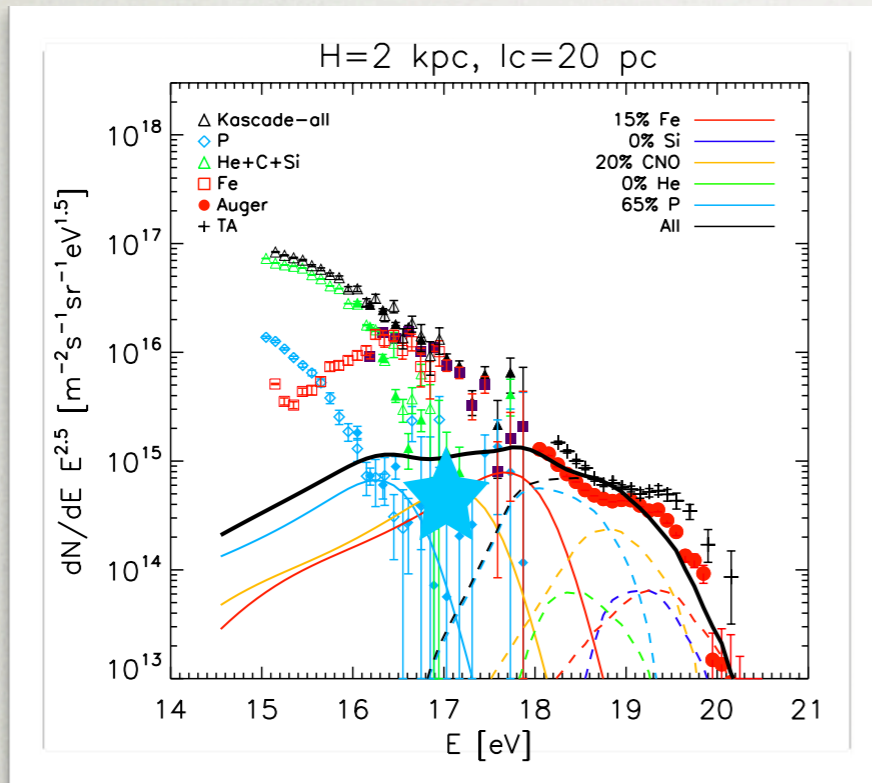


Assume sources homogeneously distributed in the disc, small scale anisotropy can be estimated as (Blasi & Amato 2011b)

$$\delta = \frac{3}{2^{3/2} \pi^{1/2}} \frac{D(E)}{Hc}$$



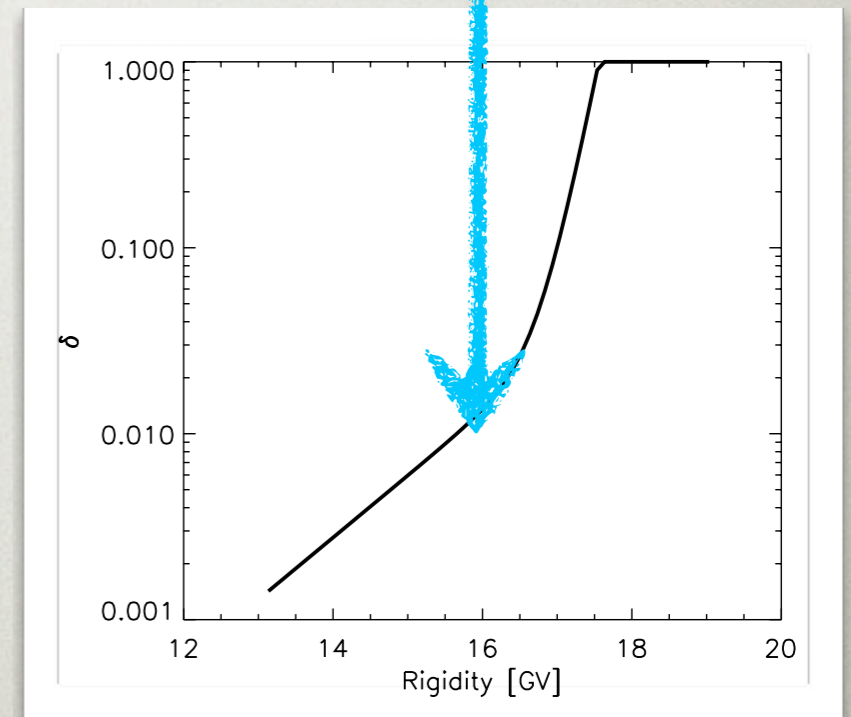
# Estimation on Anisotropy



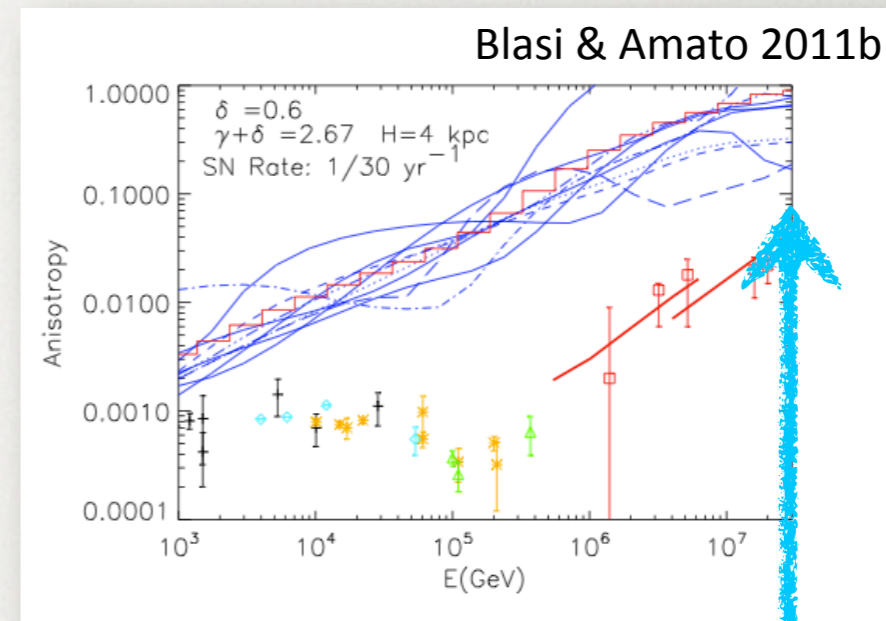
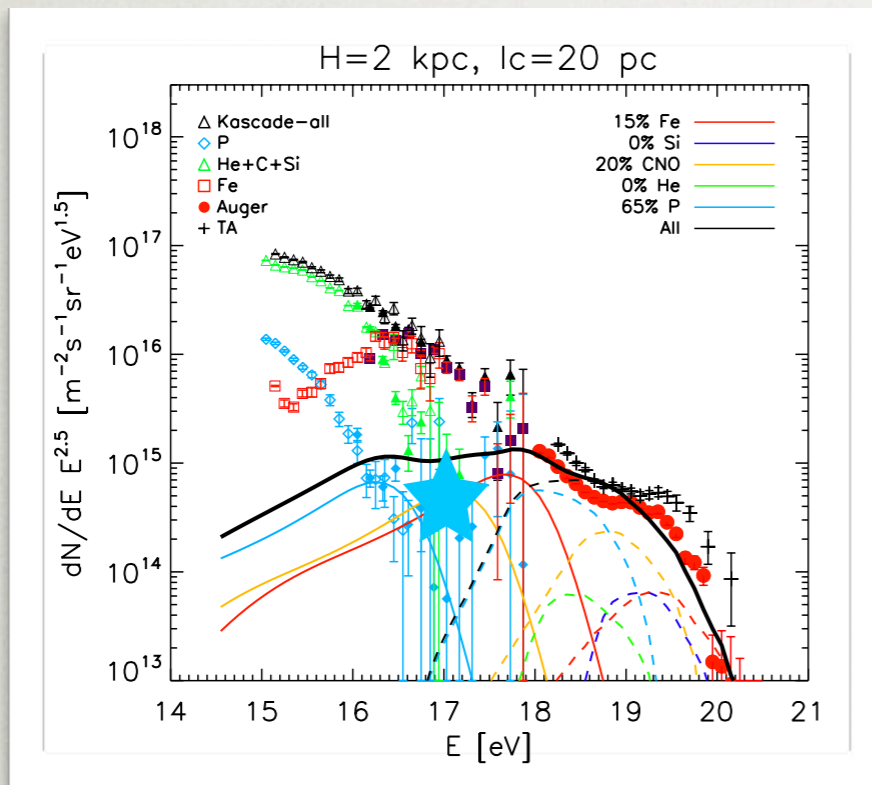
Heavy composition reduces anisotropy levels

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# Estimation on Anisotropy



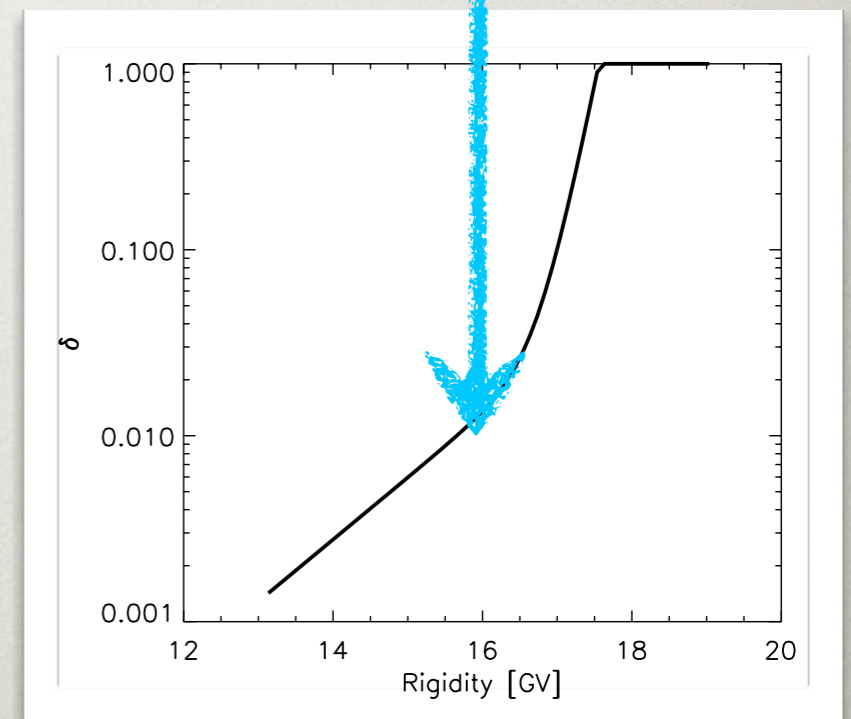
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Assume sources homogeneously distributed in the disc, small scale anisotropy can be estimated as (Blasi & Amato 2011b)

$$\delta = \frac{3}{2^{3/2} \pi^{1/2}} \frac{D(E)}{Hc}$$

Conclusion II

Galactic pulsars can contribute between the knee and the ankle!



# Testable Scenario?

# Neutrinos from Integrated Pulsar Sources

$$\Phi_\nu = \frac{f_s}{4\pi} \int_0^{z_H} \int_0^{t_\nu} \frac{dN_\nu}{dt' dE_\nu 4\pi D^2} dt' \Re(z) 4\pi D^2 \frac{dD}{dz} dz$$



# Neutrinos from Integrated Pulsar Sources

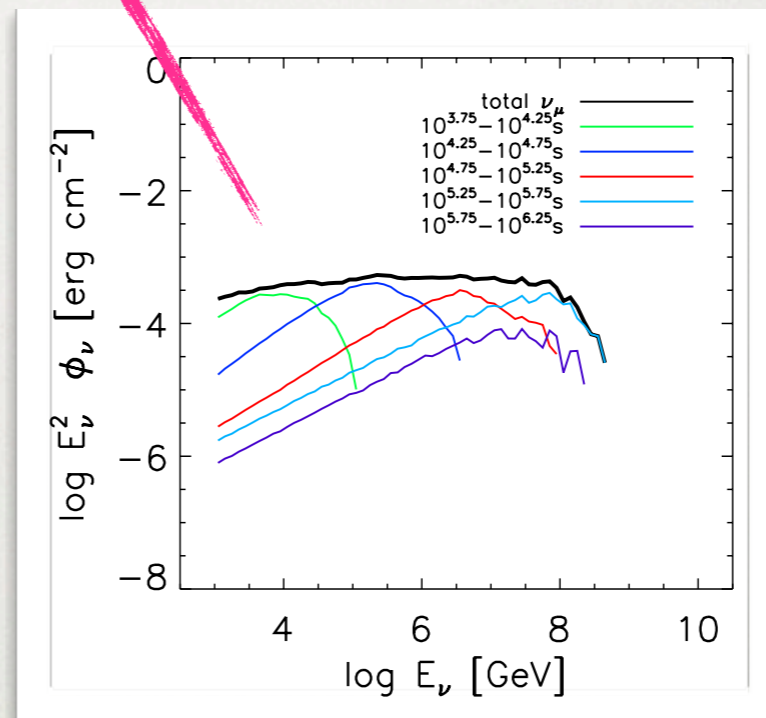
$$\Phi_\nu = \frac{f_s}{4\pi} \int_0^{z_H} \int_0^{t_\nu} \frac{dN_\nu}{dt' dE_\nu} dt' \mathcal{R}(z) 4\pi D^2 \frac{dD}{dz} dz$$

$$\mathcal{R}(0) \approx 3.3 \times 10^{-4} \text{ yr}^{-1} \text{ Mpc}^{-3}$$

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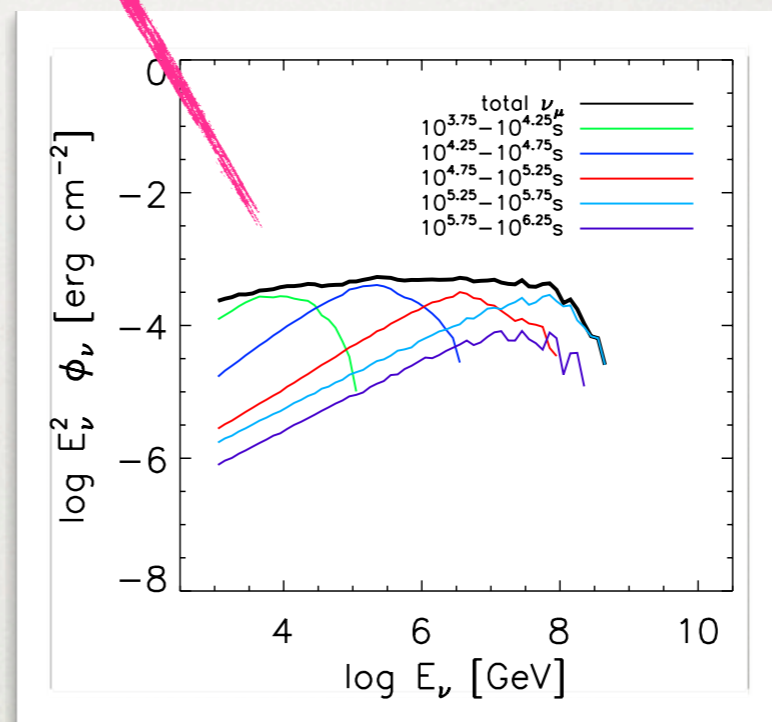
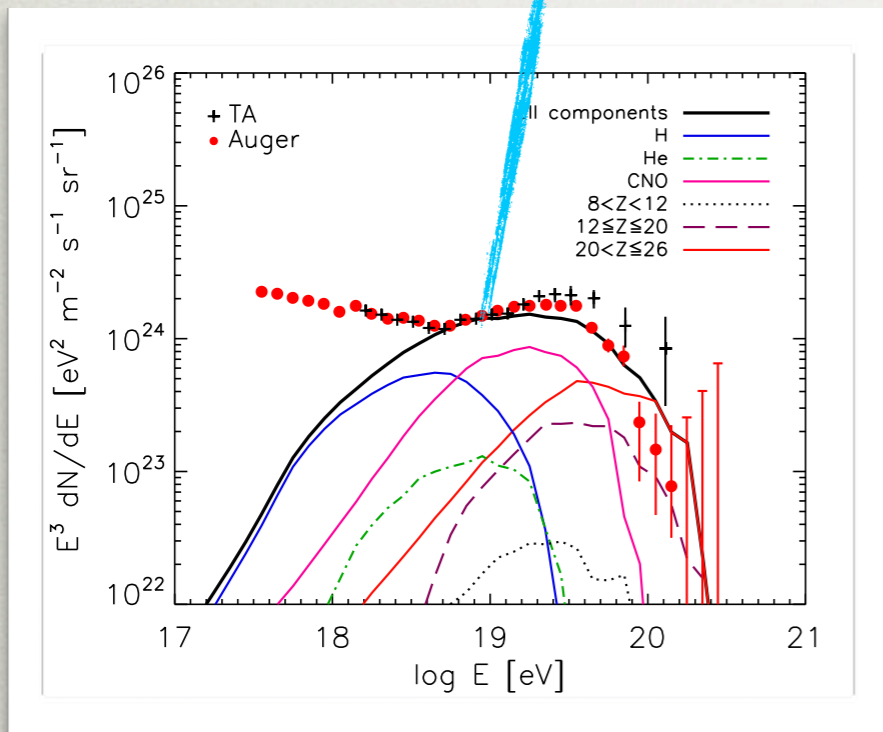


Insensible to  
injection composition:  
Fe~56P/26  
CNO~28P/14

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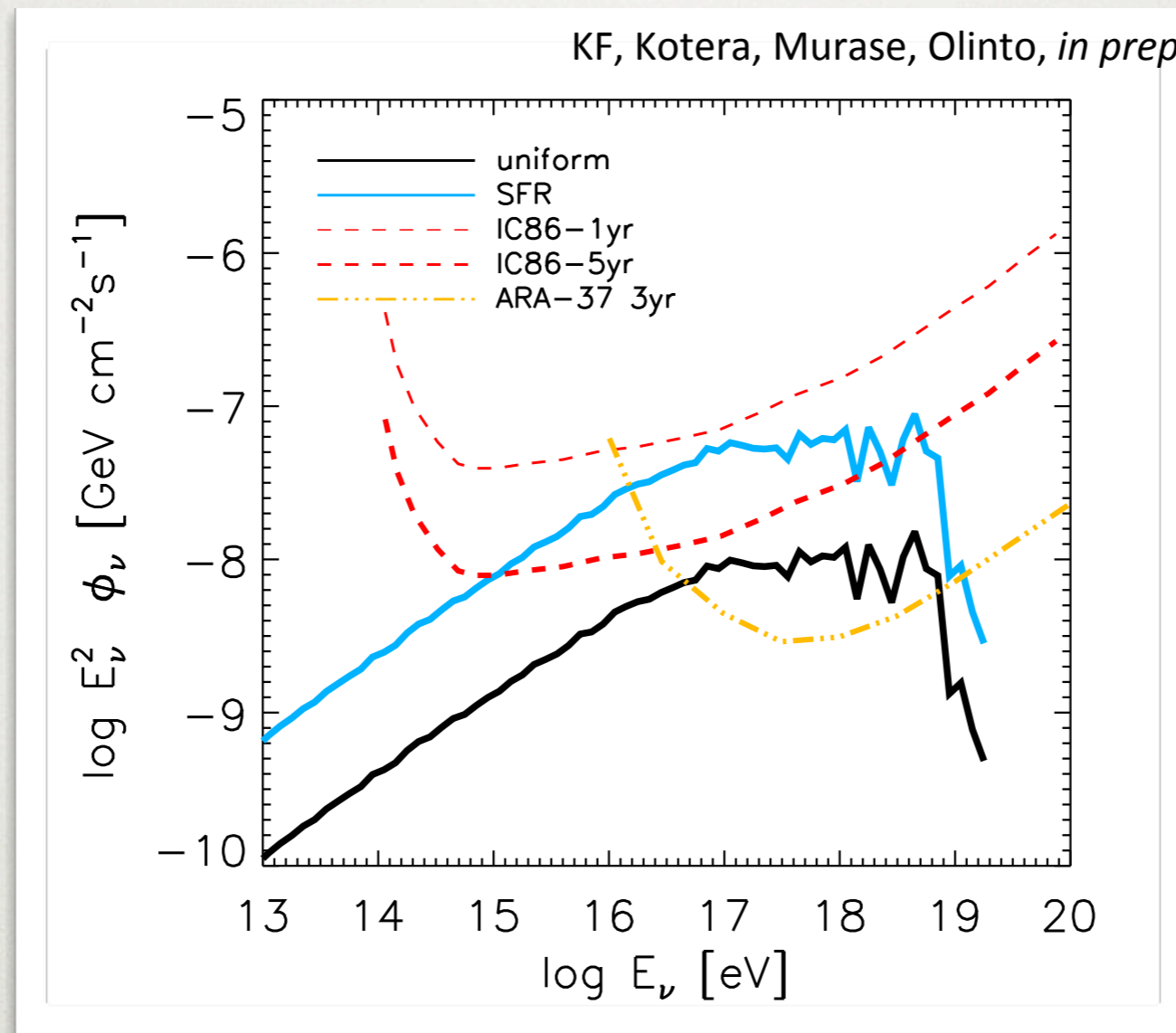


Invisible to injection composition:  
 $\text{Fe} \sim 56\text{P}/26$   
 $\text{CNO} \sim 28\text{P}/14$

$$\dot{N} = c \left( \frac{\Omega B_{*, \text{dipole}}}{2\pi c} \right) (2\pi A_{\text{cap}}) \times \text{Neutrino-loud lifetime} \rightarrow f_s \approx 0.05$$

Measured UHECR flux

# Neutrinos from Integrated Pulsar Sources



## Conclusion III

Consistent with current detection upper limits;  
Robustly tested with IC86-5 year and projected ARA-37 3 year operations.